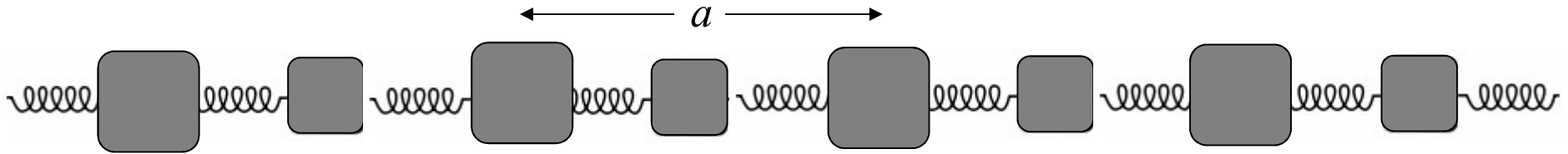


Phonons

Linear chain M_1 and M_2



$$-\omega^2 M_1 u_k = C v_k (1 + \exp(-ika)) - 2C u_k$$

$$-\omega^2 M_2 v_k = C u_k (1 + \exp(ika)) - 2C v_k$$

$$\begin{bmatrix} \omega^2 M_1 - 2C & C(1 + \exp(-ika)) \\ C(1 + \exp(ika)) & \omega^2 M_2 - 2C \end{bmatrix} \begin{bmatrix} u_k \\ v_k \end{bmatrix} = 0$$

$$M_1 M_2 \omega^4 - 2C(M_1 + M_2) \omega^2 + 2C^2 (1 - \cos(ka)) = 0$$

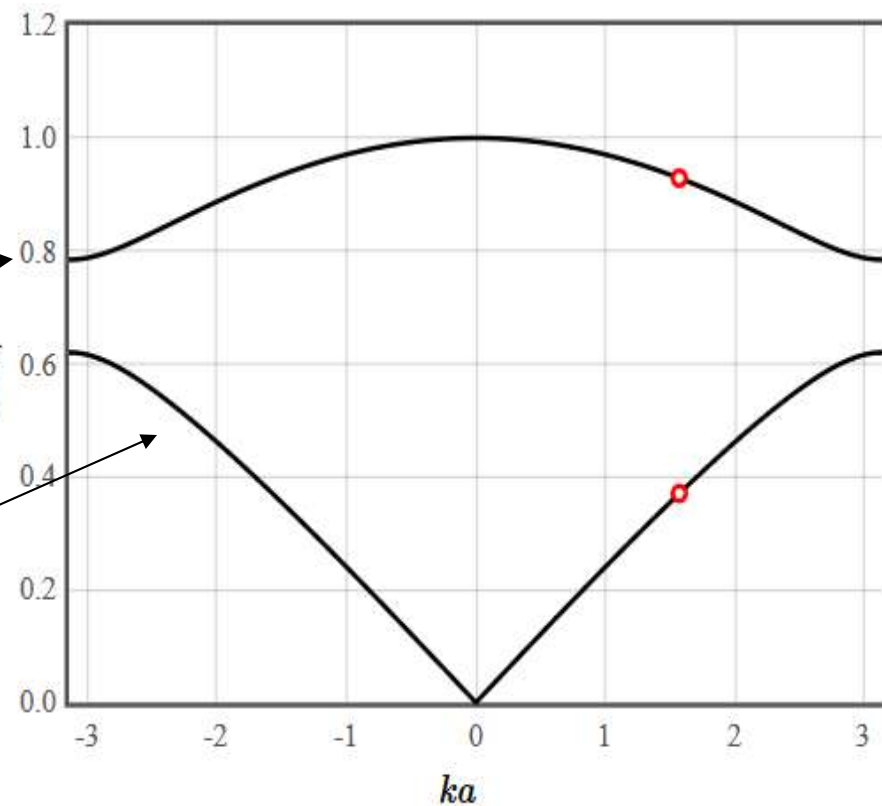
dispersion relation

$$\omega^2 = C \left(\frac{1}{M_1} + \frac{1}{M_2} \right) \pm C \sqrt{\left(\frac{1}{M_1} + \frac{1}{M_2} \right)^2 - \frac{4 \sin^2 \left(\frac{ka}{2} \right)}{M_1 M_2}}$$

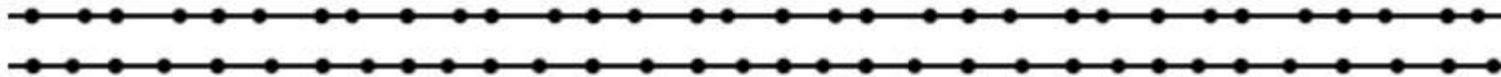
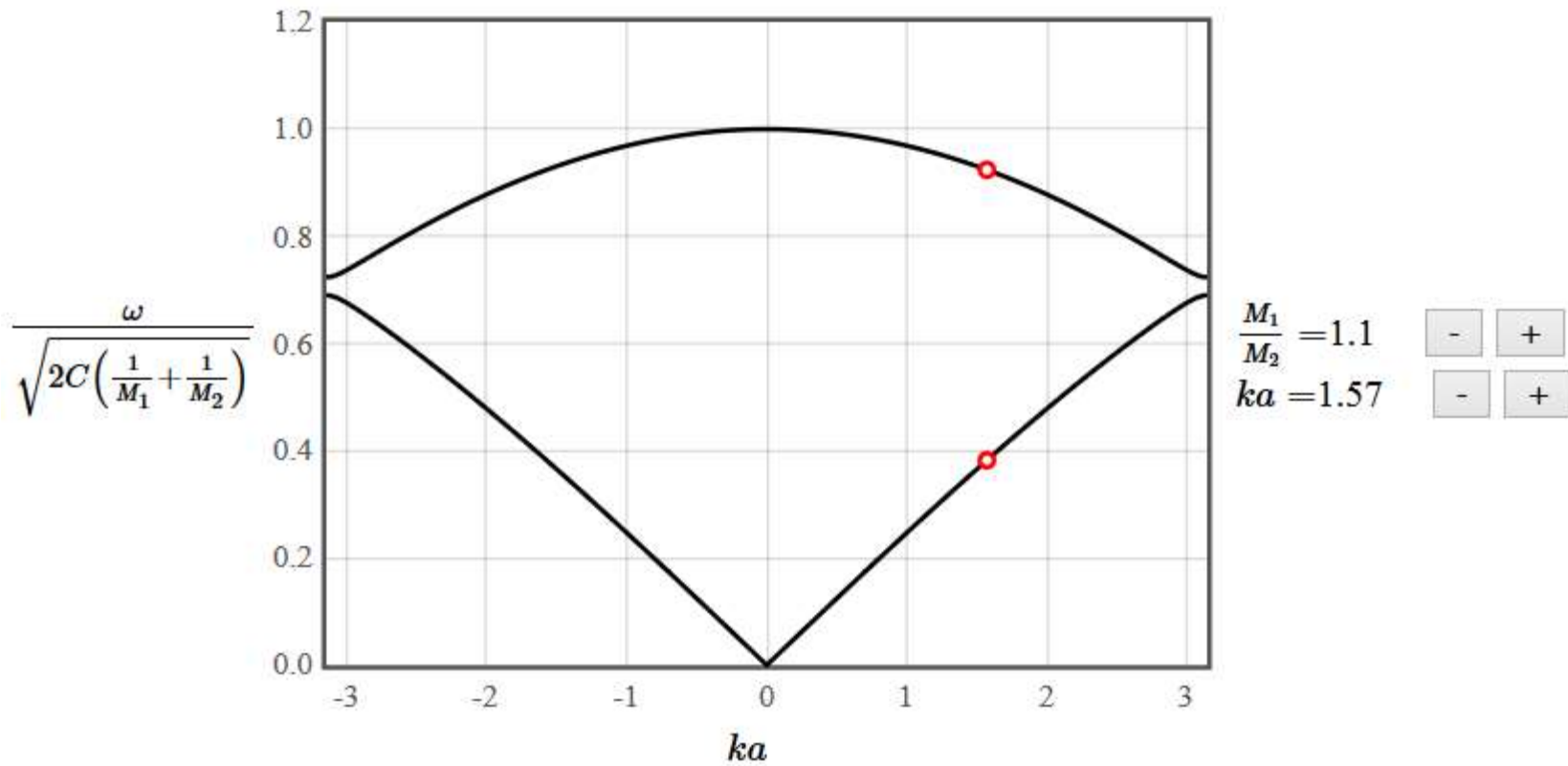
Optical phonon branch

$$\frac{\omega}{\sqrt{2C \left(\frac{1}{M_1} + \frac{1}{M_2} \right)}}$$

Acoustic phonon branch

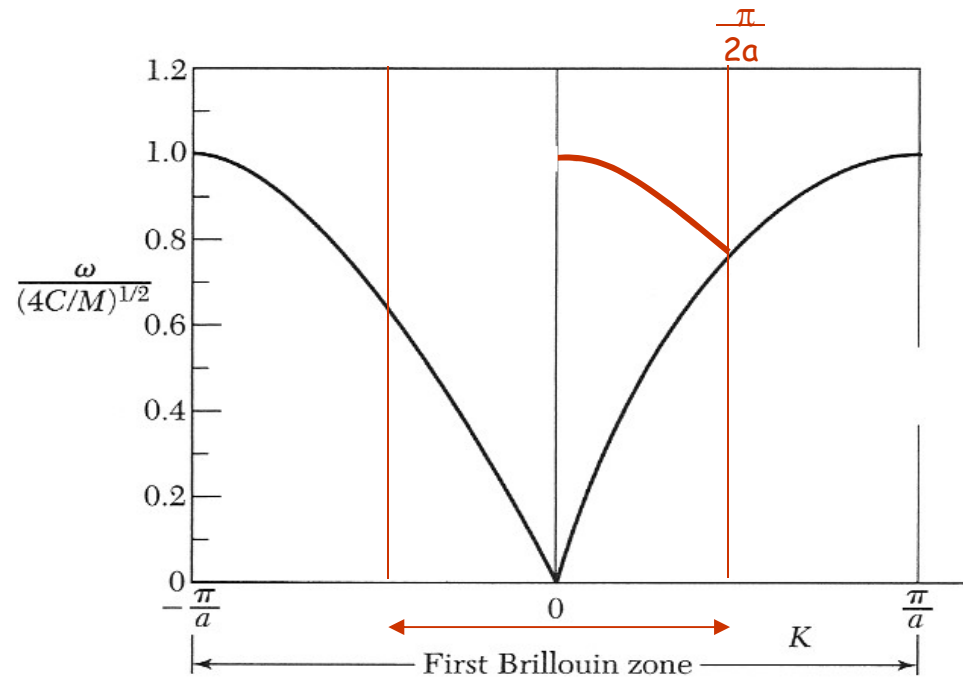


normal modes



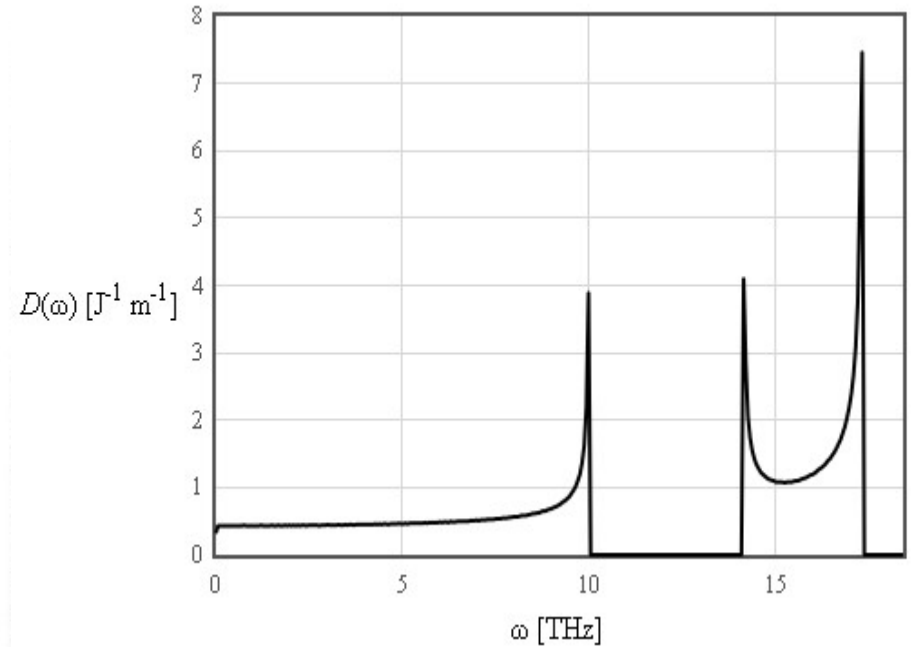
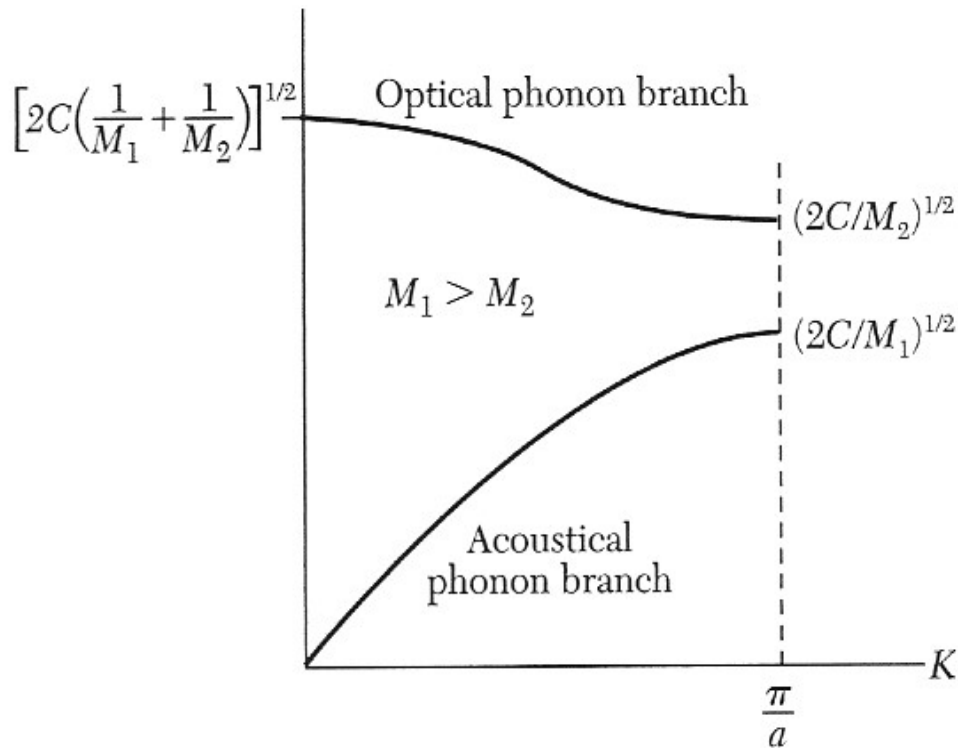
<http://lampx.tugraz.at/~hadley/ss1/phonons/1d/1d2m.php>

Linear chain M_1 and M_2

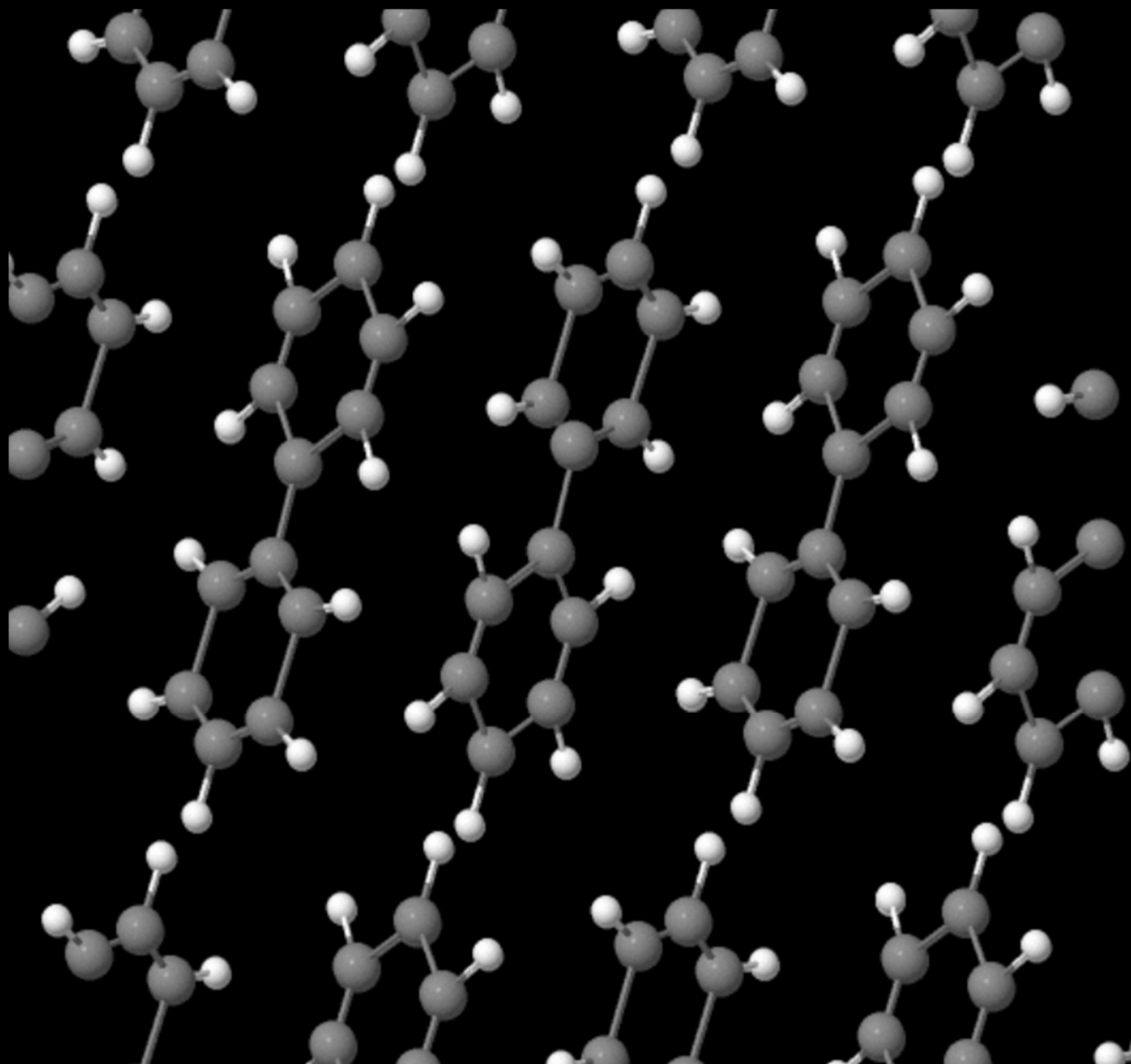


The branches of the dispersion curves can be translated by a reciprocal lattice vector \vec{G} .

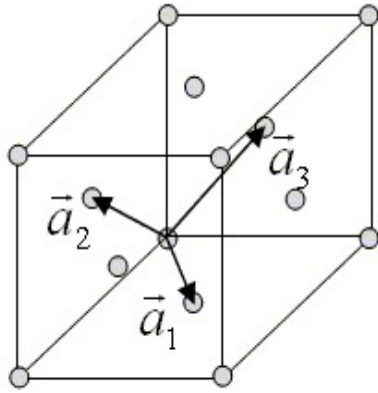
density of states



$$\omega^2 = C \left(\frac{1}{M_1} + \frac{1}{M_2} \right) \pm C \sqrt{\left(\frac{1}{M_1} + \frac{1}{M_2} \right)^2 - \frac{4 \sin^2 ka}{M_1 M_2}}$$



fcc



$$\vec{a}_1 = \frac{a}{2} \hat{x} + \frac{a}{2} \hat{y}$$

$$\vec{a}_2 = \frac{a}{2} \hat{x} + \frac{a}{2} \hat{z}$$

$$\vec{a}_3 = \frac{a}{2} \hat{y} + \frac{a}{2} \hat{z}$$

$$\vec{b}_1 = \frac{2\pi}{a} (\hat{k}_x + \hat{k}_y - \hat{k}_z)$$

$$\vec{b}_2 = \frac{2\pi}{a} (\hat{k}_x - \hat{k}_y + \hat{k}_z)$$

$$\vec{b}_3 = \frac{2\pi}{a} (-\hat{k}_x + \hat{k}_y + \hat{k}_z)$$

$$\begin{aligned} m \frac{d^2 u_{lmn}^x}{dt^2} = & \frac{C}{2} \left[(u_{l+1mn}^x - u_{lmn}^x) + (u_{l-1mn}^x - u_{lmn}^x) + (u_{lm+1n}^x - u_{lmn}^x) + (u_{lm-1n}^x - u_{lmn}^x) \right. \\ & + (u_{l+1mn-1}^x - u_{lmn}^x) + (u_{l-1mn+1}^x - u_{lmn}^x) + (u_{lm+1n-1}^x - u_{lmn}^x) + (u_{lm-1n+1}^x - u_{lmn}^x) \\ & + (u_{l+1mn}^y - u_{lmn}^y) + (u_{l-1mn}^y - u_{lmn}^y) - (u_{lm+1n-1}^y - u_{lmn}^y) - (u_{lm-1n+1}^y - u_{lmn}^y) \\ & \left. + (u_{lm+1n}^z - u_{lmn}^z) + (u_{lm-1n}^z - u_{lmn}^z) - (u_{l+1mn-1}^z - u_{lmn}^z) - (u_{l-1mn+1}^z - u_{lmn}^z) \right] \end{aligned}$$

and similar expressions for the y and z motion

Normal modes are eigenfunctions of T

$$u_{lmn}^x = u_{\vec{k}}^x \exp\left(i\left(l\vec{k} \cdot \vec{a}_1 + m\vec{k} \cdot \vec{a}_2 + n\vec{k} \cdot \vec{a}_3 - \omega t\right)\right)$$

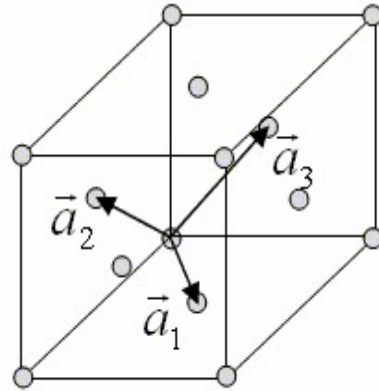
$$u_{lmn}^y = u_{\vec{k}}^y \exp\left(i\left(l\vec{k} \cdot \vec{a}_1 + m\vec{k} \cdot \vec{a}_2 + n\vec{k} \cdot \vec{a}_3 - \omega t\right)\right)$$

$$u_{lmn}^z = u_{\vec{k}}^z \exp\left(i\left(l\vec{k} \cdot \vec{a}_1 + m\vec{k} \cdot \vec{a}_2 + n\vec{k} \cdot \vec{a}_3 - \omega t\right)\right)$$

These are eigenfunctions of T.

$$\begin{aligned} T_{pqr} u_{lmn}^x &= u_{\vec{k}}^x \exp\left(i\left(l\vec{k} \cdot (\vec{a}_1 + p\vec{a}_1) + m\vec{k} \cdot (\vec{a}_2 + q\vec{a}_2) + n\vec{k} \cdot (\vec{a}_3 + r\vec{a}_3) - \omega t\right)\right) \\ &= \exp\left(i\left(lp\vec{k} \cdot \vec{a}_1 + qm\vec{k} \cdot \vec{a}_2 + rn\vec{k} \cdot \vec{a}_3\right)\right) u_{\vec{k}}^x \exp\left(i\left(l\vec{k} \cdot \vec{a}_1 + m\vec{k} \cdot \vec{a}_2 + n\vec{k} \cdot \vec{a}_3 - \omega t\right)\right) \\ &= \exp\left(i\left(lp\vec{k} \cdot \vec{a}_1 + qm\vec{k} \cdot \vec{a}_2 + rn\vec{k} \cdot \vec{a}_3\right)\right) u_{lmn}^x \end{aligned}$$

fcc



$$\vec{a}_1 = \frac{a}{2} \hat{x} + \frac{a}{2} \hat{y}$$

$$\vec{a}_2 = \frac{a}{2} \hat{x} + \frac{a}{2} \hat{z}$$

$$\vec{a}_3 = \frac{a}{2} \hat{y} + \frac{a}{2} \hat{z}$$

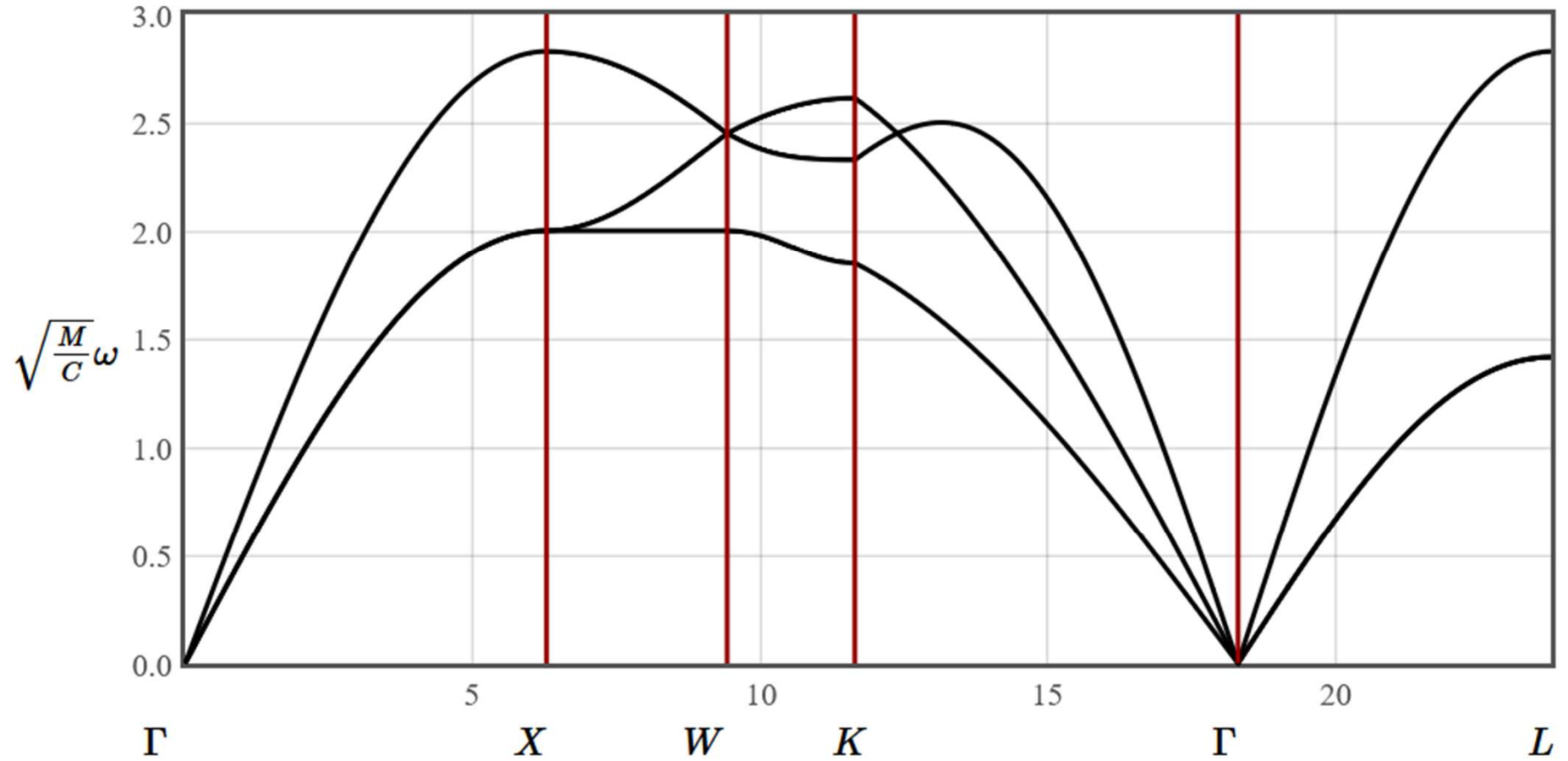
Substitute the eigenfunctions of T into Newton's laws.

$$u_{lmn}^x = u_k^x \exp\left(i\left(l\vec{k} \cdot \vec{a}_1 + m\vec{k} \cdot \vec{a}_2 + n\vec{k} \cdot \vec{a}_3\right)\right) = u_k^x \exp\left(i\left(\frac{(l+m)k_x a}{2} + \frac{(l+n)k_y a}{2} + \frac{(m+n)k_z a}{2}\right)\right).$$

$$\begin{vmatrix} 4 - \cos\left(\frac{k_x a}{2} + \frac{k_y a}{2}\right) - \cos\left(\frac{k_x a}{2} + \frac{k_z a}{2}\right) - \cos\left(\frac{k_x a}{2} - \frac{k_y a}{2}\right) - \cos\left(\frac{k_x a}{2} - \frac{k_z a}{2}\right) - \frac{m\omega^2}{\sqrt{2}C} & -\cos\left(\frac{k_x a}{2} + \frac{k_y a}{2}\right) + \cos\left(\frac{k_x a}{2} - \frac{k_y a}{2}\right) & -\cos\left(\frac{k_x a}{2} + \frac{k_z a}{2}\right) + \cos\left(\frac{k_x a}{2} - \frac{k_z a}{2}\right) \\ -\cos\left(\frac{k_x a}{2} + \frac{k_y a}{2}\right) + \cos\left(\frac{k_x a}{2} - \frac{k_y a}{2}\right) & 4 - \cos\left(\frac{k_y a}{2} + \frac{k_z a}{2}\right) - \cos\left(\frac{k_y a}{2} + \frac{k_x a}{2}\right) - \cos\left(\frac{k_y a}{2} - \frac{k_x a}{2}\right) - \cos\left(\frac{k_y a}{2} - \frac{k_z a}{2}\right) - \frac{m\omega^2}{\sqrt{2}C} & -\cos\left(\frac{k_y a}{2} + \frac{k_z a}{2}\right) + \cos\left(\frac{k_y a}{2} - \frac{k_z a}{2}\right) \\ -\cos\left(\frac{k_x a}{2} + \frac{k_z a}{2}\right) + \cos\left(\frac{k_x a}{2} - \frac{k_z a}{2}\right) & -\cos\left(\frac{k_y a}{2} + \frac{k_z a}{2}\right) + \cos\left(\frac{k_y a}{2} - \frac{k_z a}{2}\right) & 4 - \cos\left(\frac{k_z a}{2} + \frac{k_x a}{2}\right) - \cos\left(\frac{k_z a}{2} + \frac{k_y a}{2}\right) - \cos\left(\frac{k_z a}{2} - \frac{k_x a}{2}\right) - \cos\left(\frac{k_z a}{2} - \frac{k_y a}{2}\right) - \frac{m\omega^2}{\sqrt{2}C} \end{vmatrix} = 0$$

<http://lamp.tu-graz.ac.at/~hadley/ss1/phonons/fcc/fcc.html>

For every k there are 3 solutions for ω .



Phonon dispersion Au

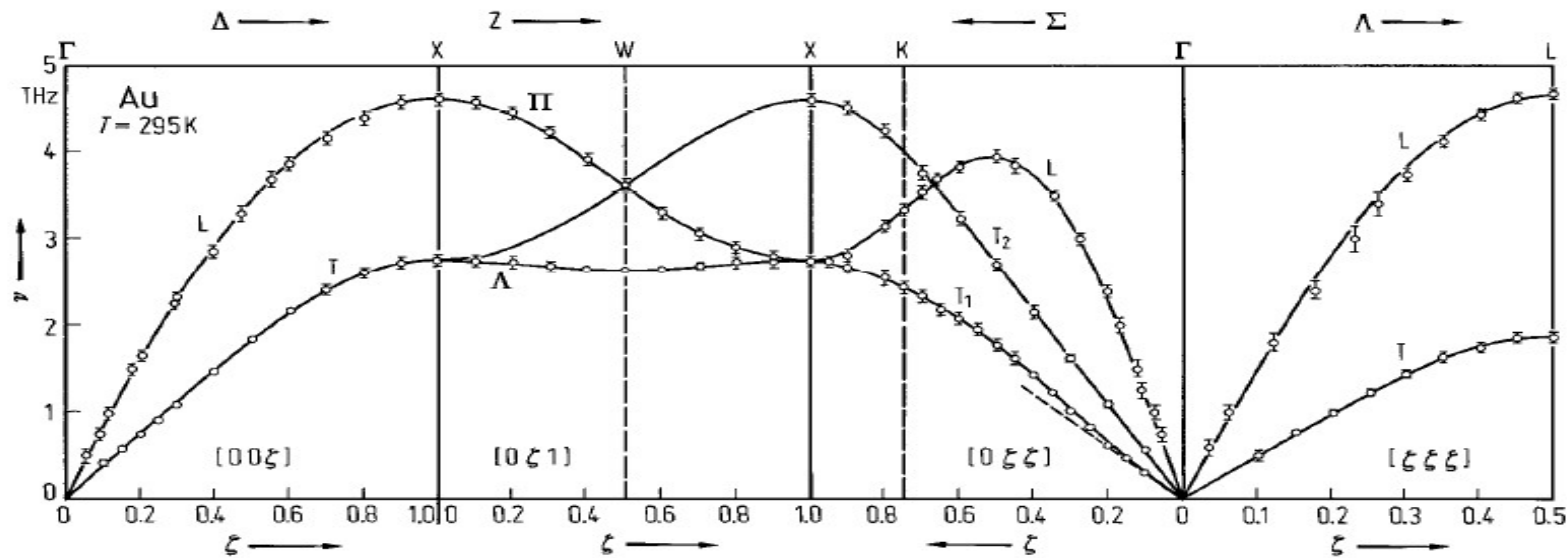
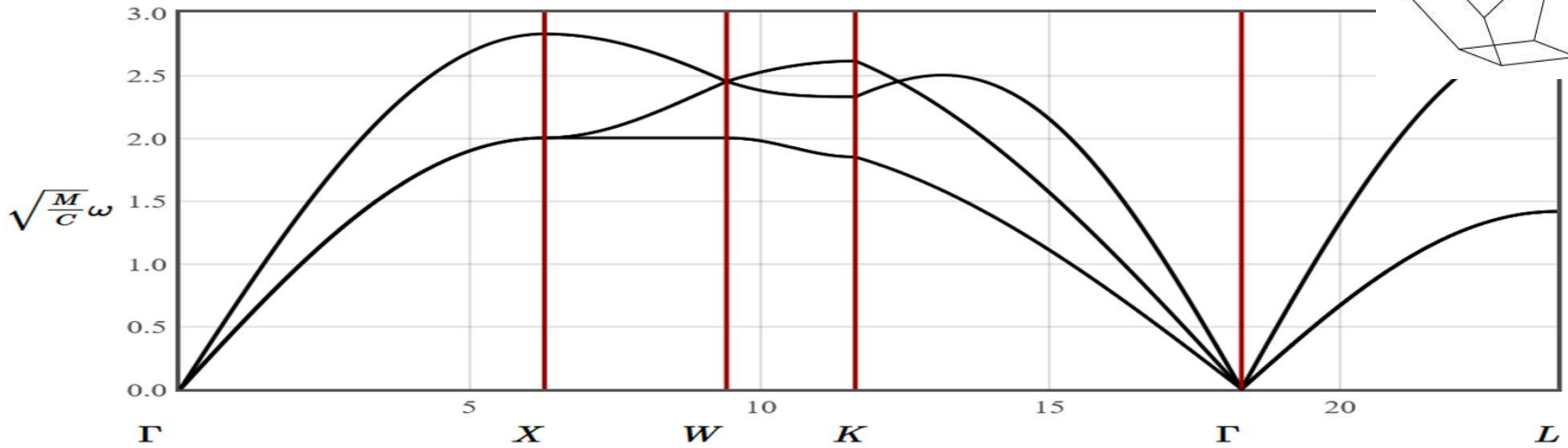
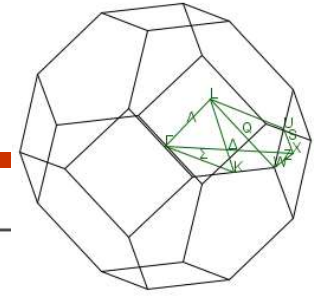
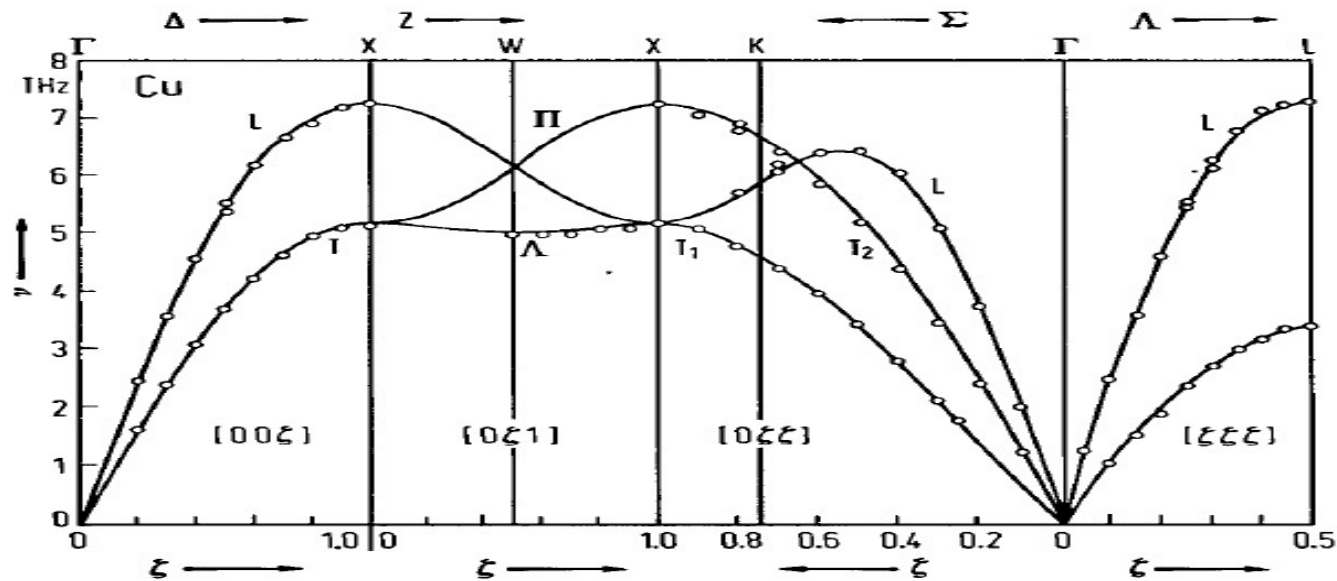
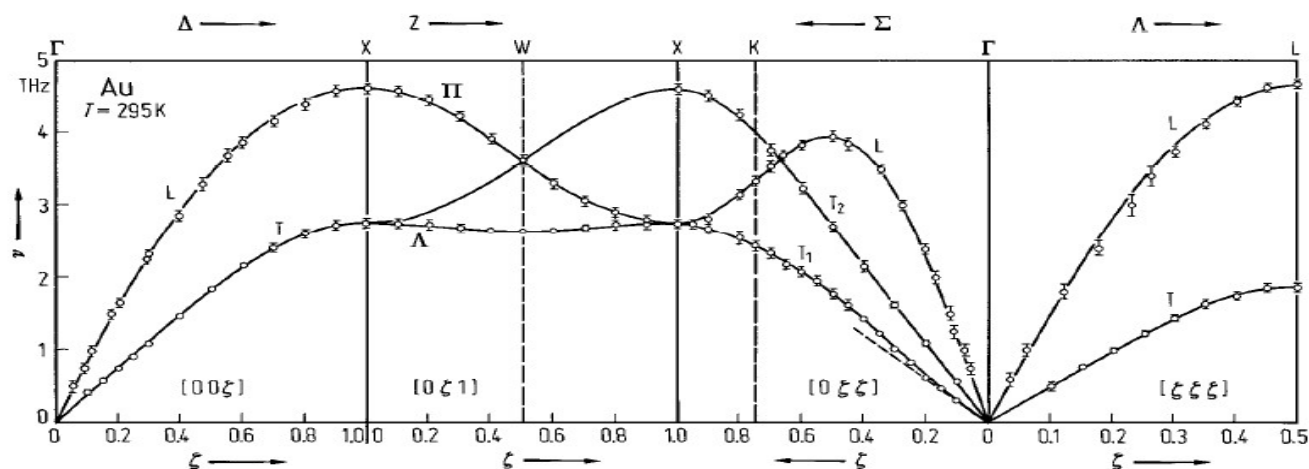


Fig. 1. Au. Phonon dispersion relations in the principal symmetry directions according to [73Ly1]. The solid curves represent both the fourth neighbour general force model (M1) and the fifth neighbour axially symmetric model (M2) of Table 3 Au. The dotted line in the Σ direction is corresponding to the velocity of sound appropriate to the $[0\xi\xi] \Gamma_1$ branch.

Materials with the same crystal structure will have similar phonon dispersion relations



Cu



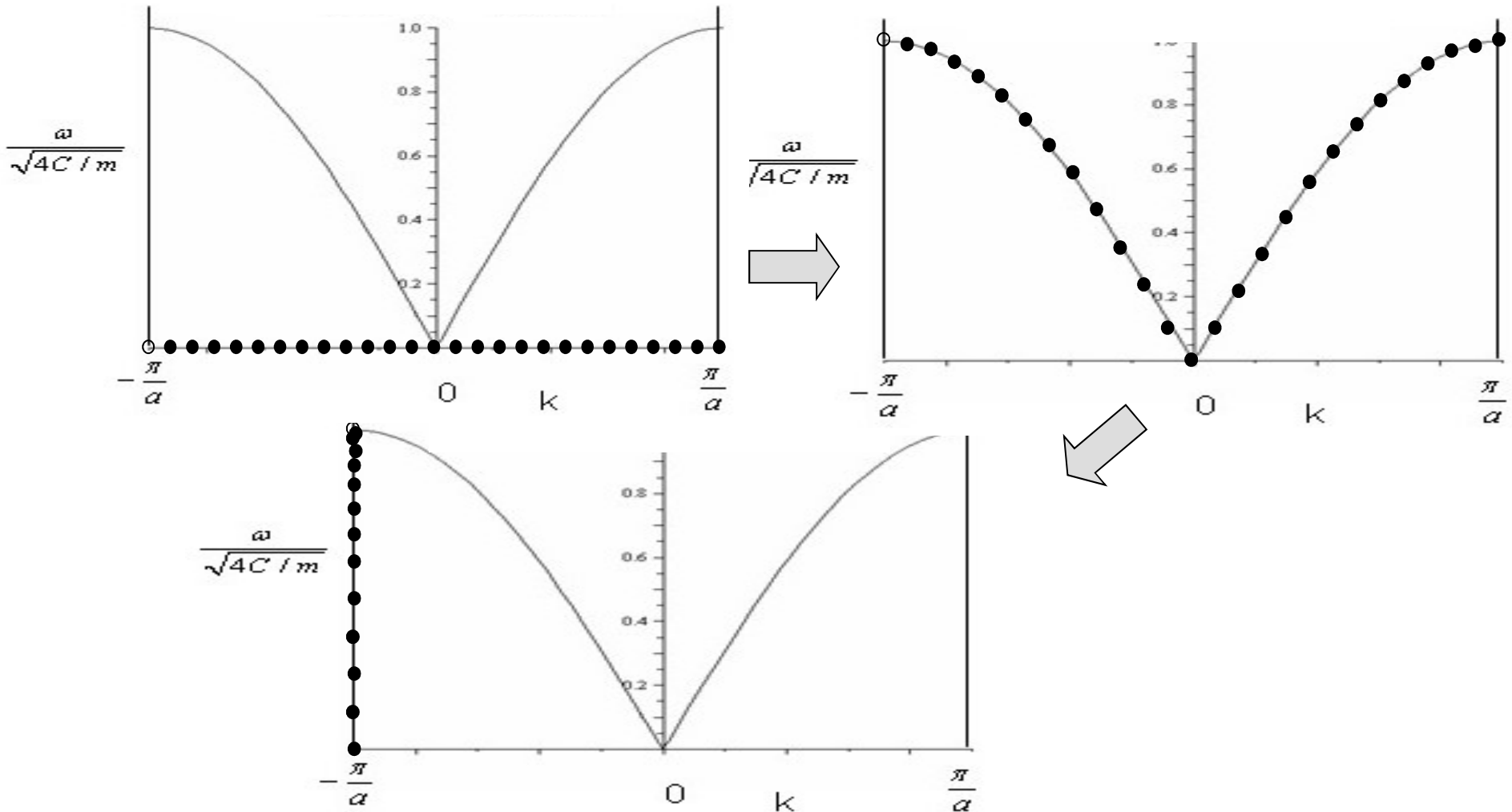
Au

Fig. 1. Au. Phonon dispersion relations in the principal symmetry directions according to [73Ly1]. The solid curves represent both the fourth neighbour general force model (M1) and the fifth neighbour axially symmetric model (M2) of Table 3 Au. The dotted line in the Σ direction is corresponding to the velocity of sound appropriate to the [0 $\zeta\zeta$] T₁ branch.

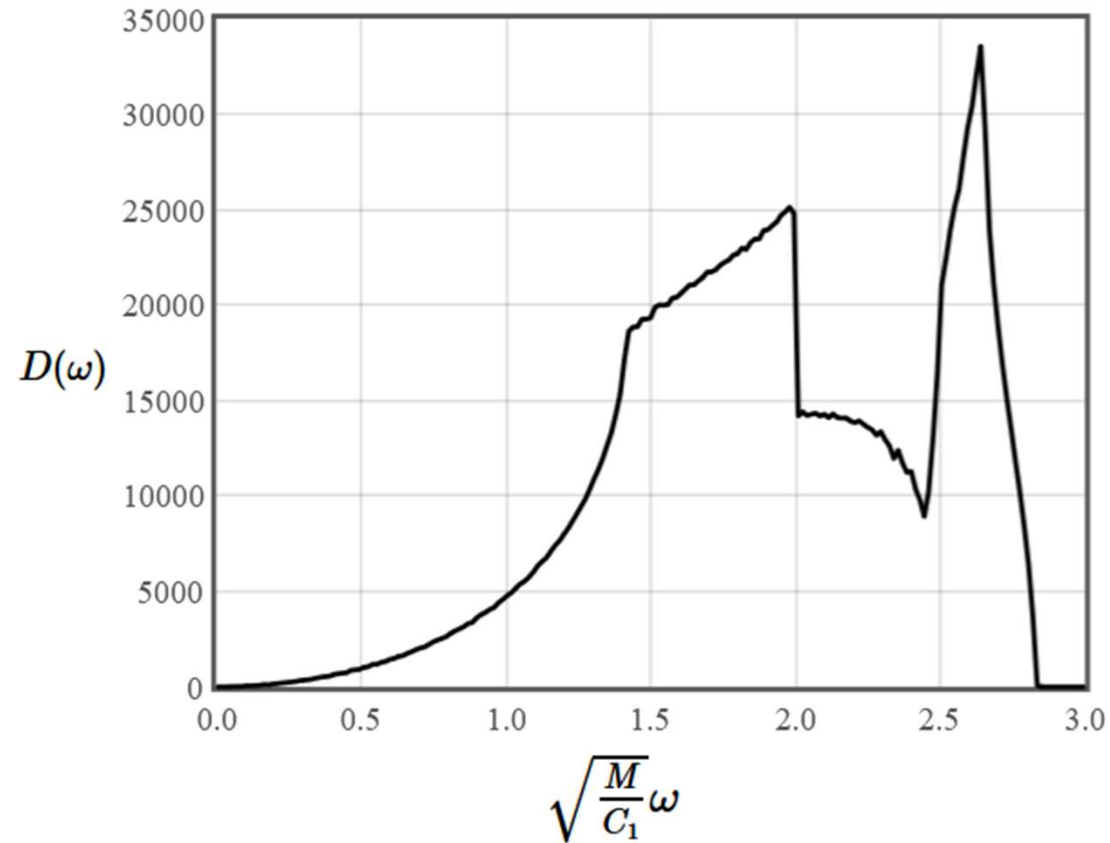
Linear Chain - density of states

Determine the density of states numerically

$$\omega = \sqrt{\frac{4C}{m}} \left| \sin\left(\frac{ka}{2}\right) \right|$$



Phonon density of states of a fcc crystal



http://lampx.tugraz.at/~hadley/ss1/phonons/fcc/fcc_dos.php

Phonon DOS fcc

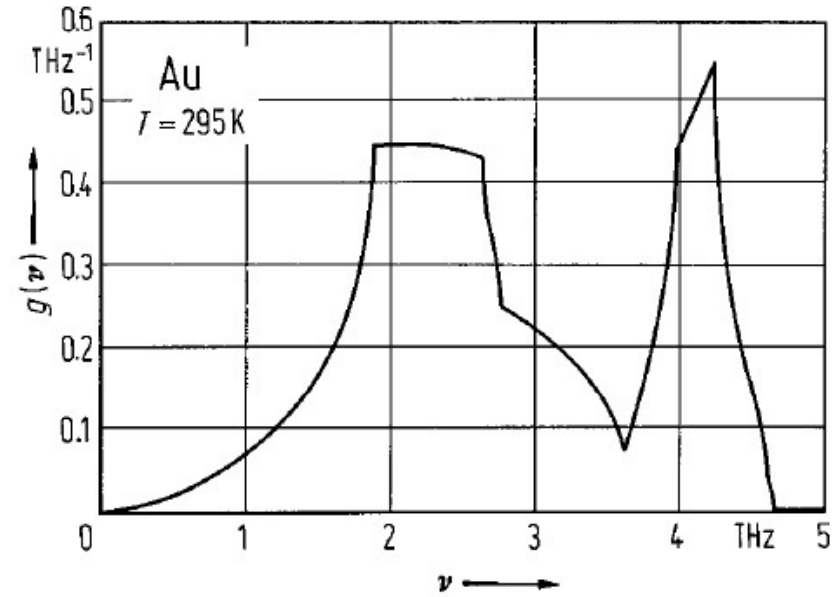
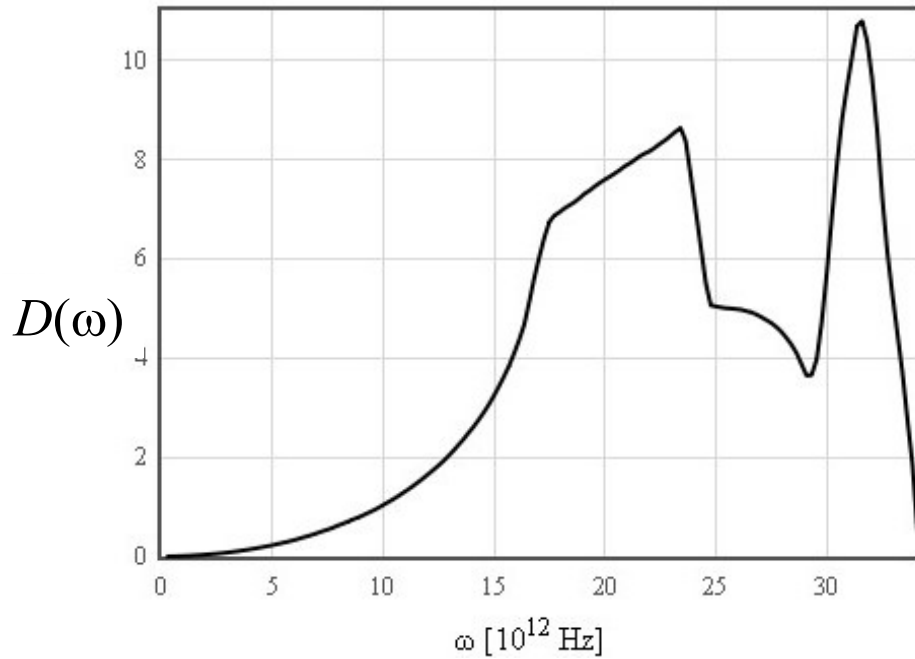
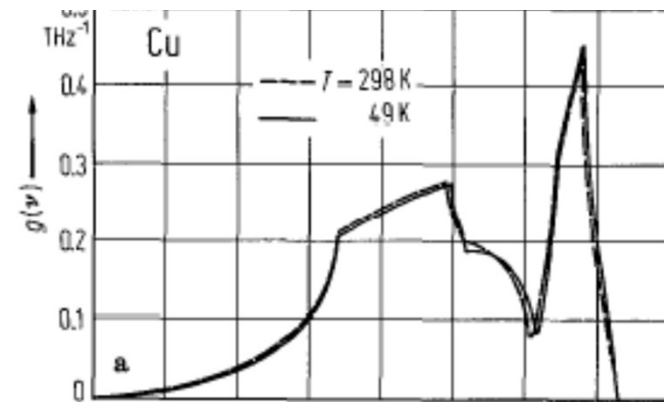
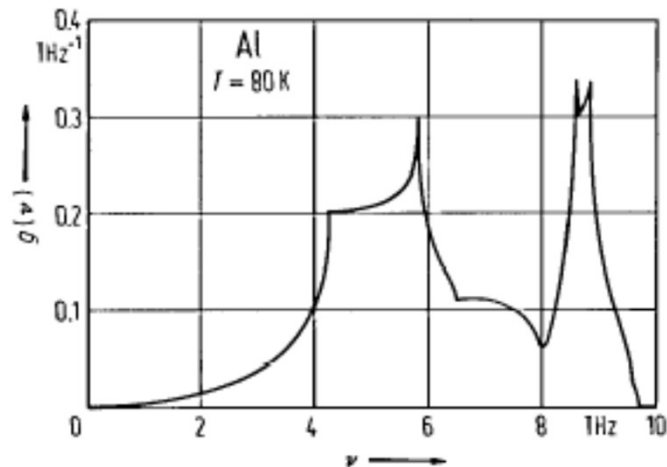


Fig. 2. Au. Frequency distribution calculated from the fourth neighbour general force constant model (M1) of Table 3 Au.



Phonon dispersion

1. Determine the dispersion relation:

Write down the equations of motion (masses and springs).

The solutions to these equations will be

$$A_{\mathbf{k}} \exp \left(i \left(\vec{k} \cdot \vec{a}_1 + \vec{k} \cdot \vec{a}_2 + \vec{k} \cdot \vec{a}_3 - \omega t \right) \right)$$

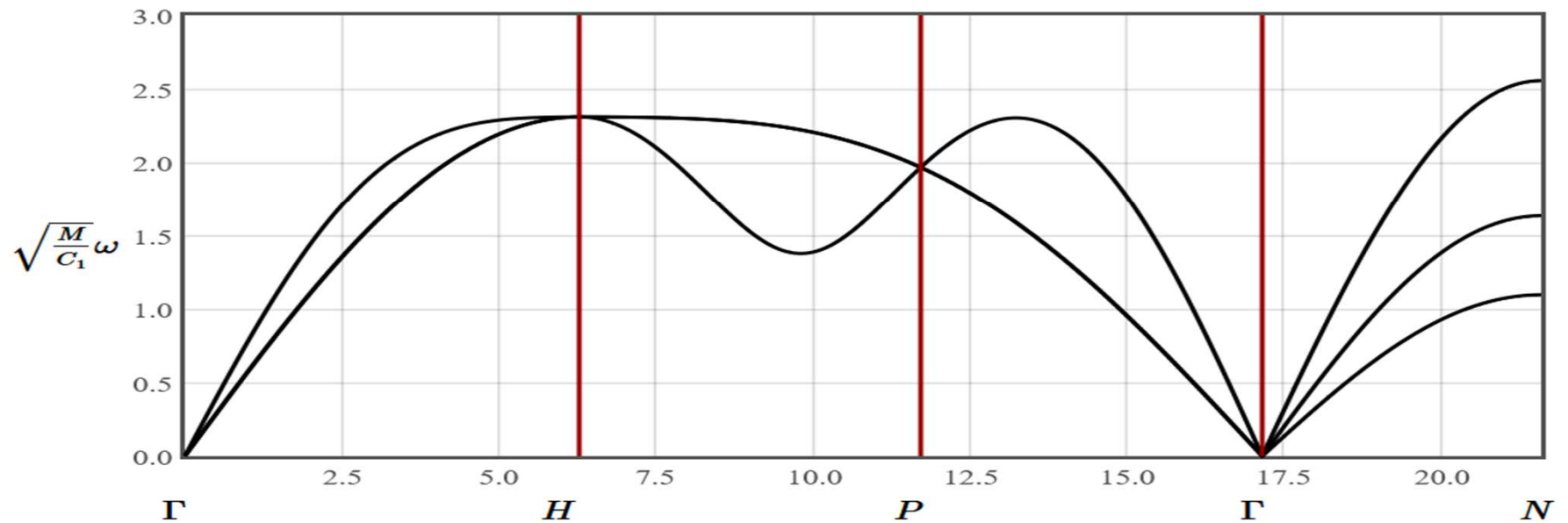
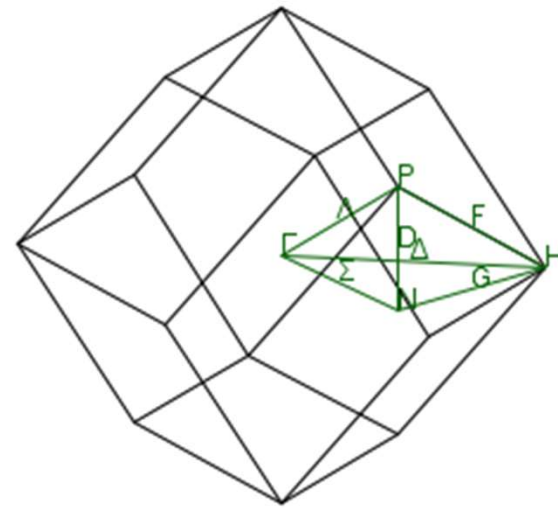
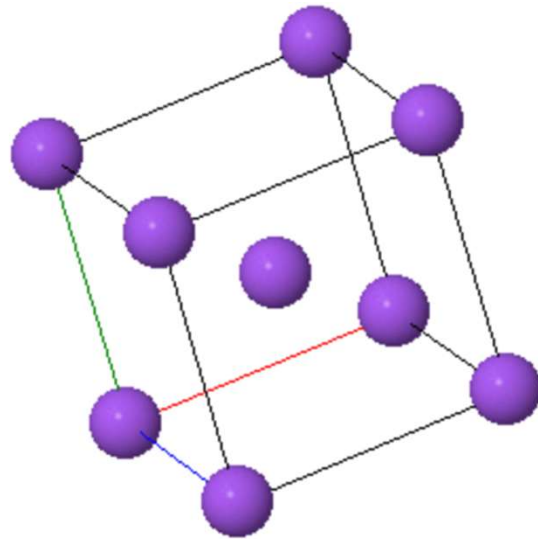
Substitute the solutions into the equations of motion to determine the dispersion relation.

2. Determine the density of states numerically from the dispersion relation

$$D(\omega)$$

For every allowed k , find all corresponding values of ω .

Phonon dispersion bcc



Phonon dispersion Fe

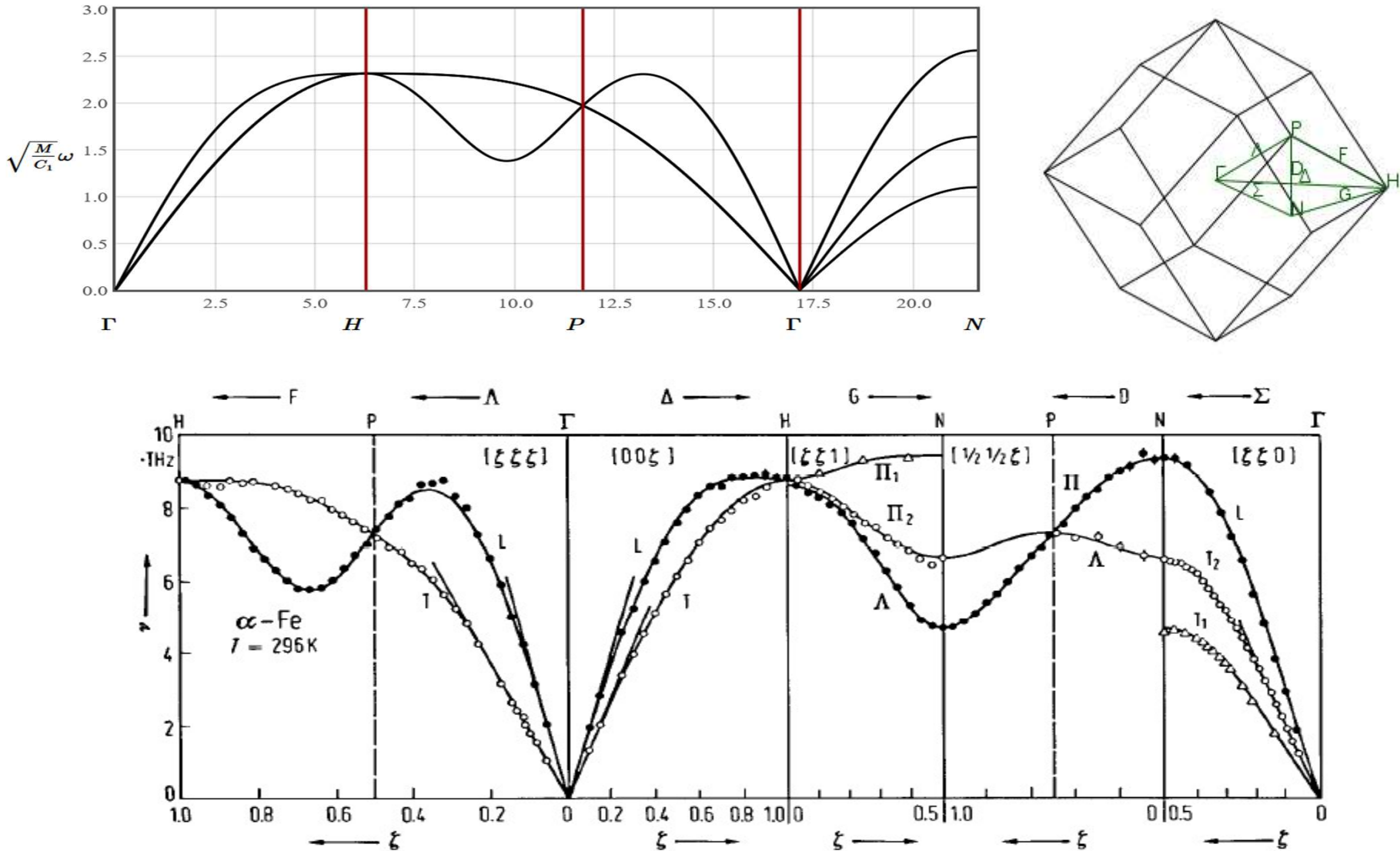


Fig. 2. Fe. Phonon dispersion curves in α -iron at 296 K. Experimental points: [68Va2]. Solid curve: fifth neighbour Born-von Karman model (Table 3 Fe [68Va2]).

From Springer Materials: Landholt Boernstein Database

Phonon DOS Fe

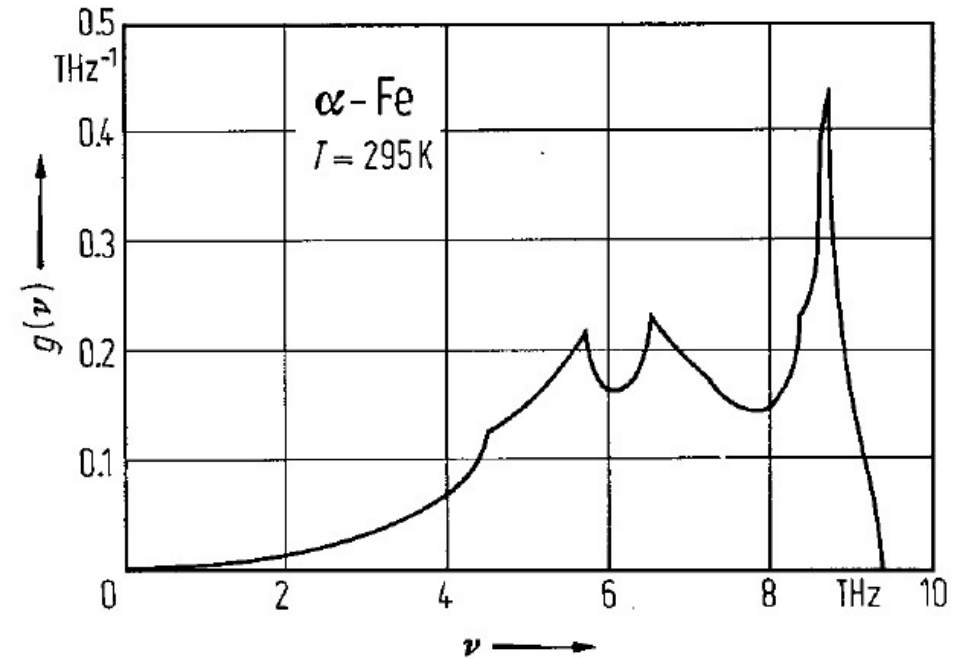
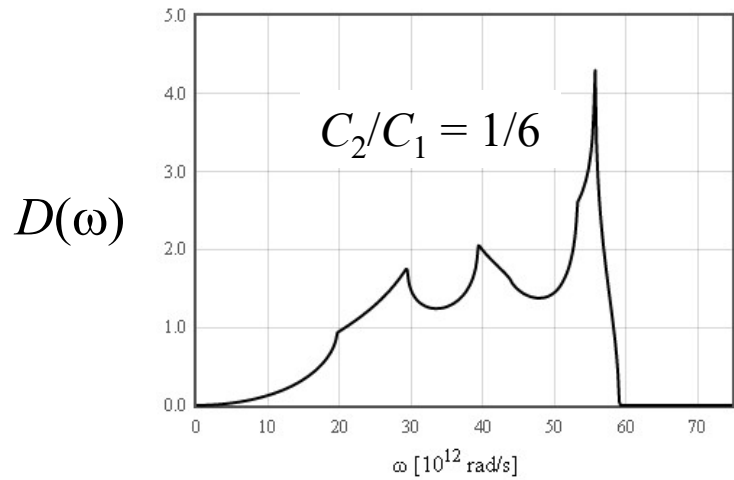
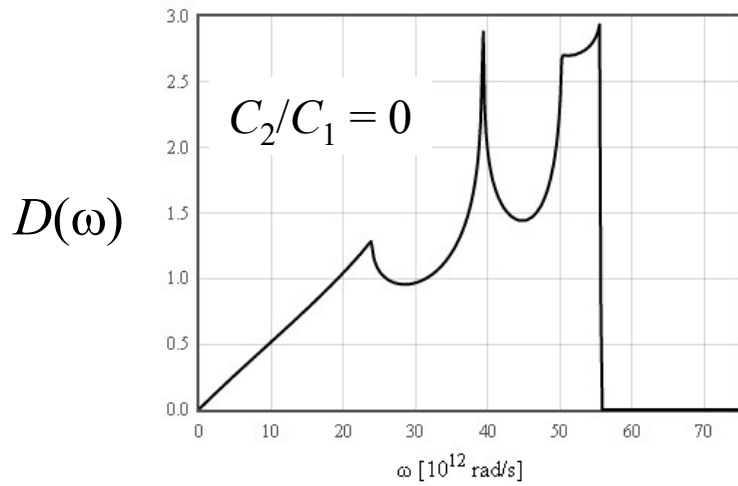
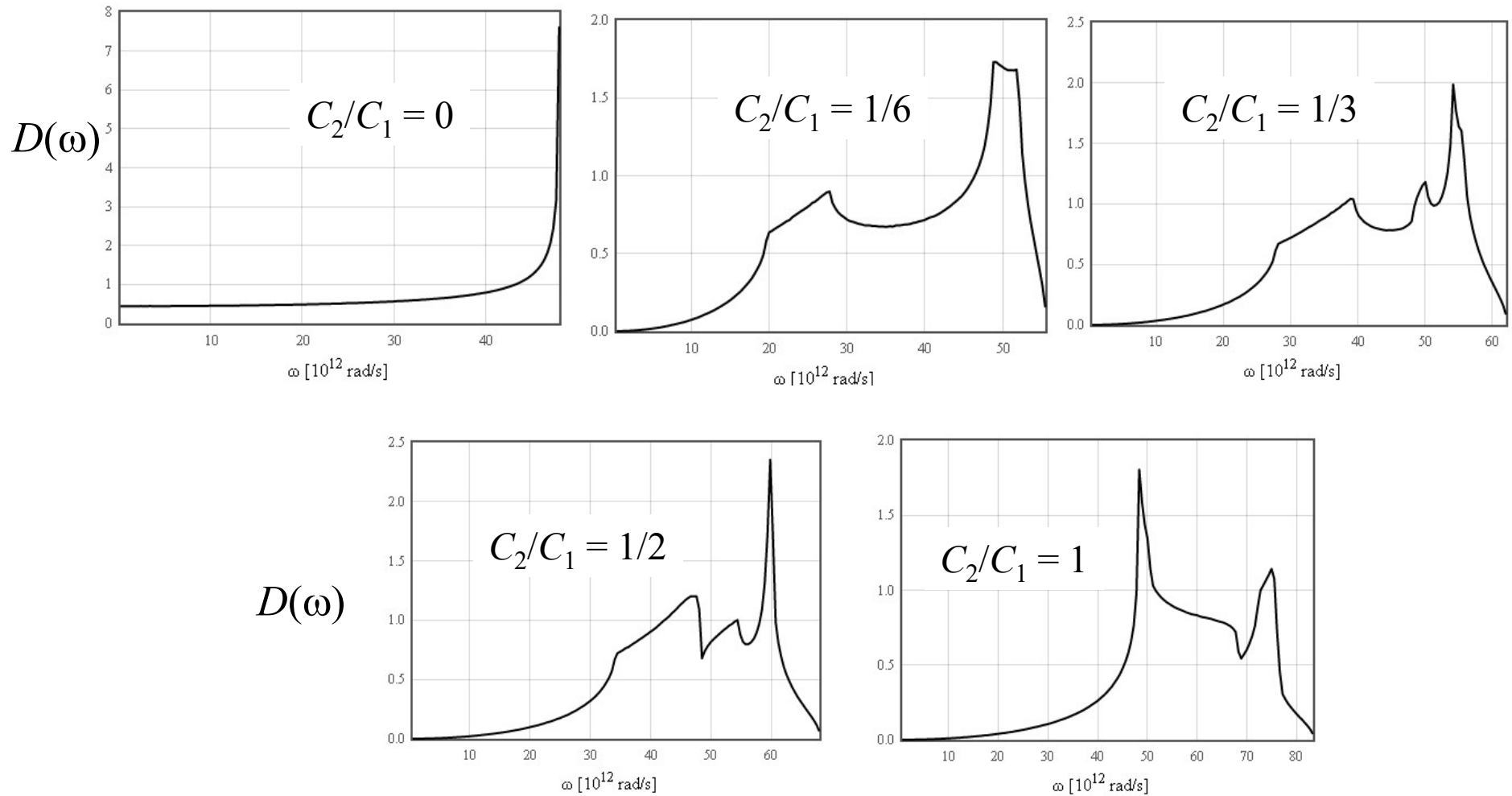


Fig. 3. Fe. Frequency spectrum of α -iron at 295 K calculated from the Born-von Karman force constants of Table 3 Fe [67Mi1].

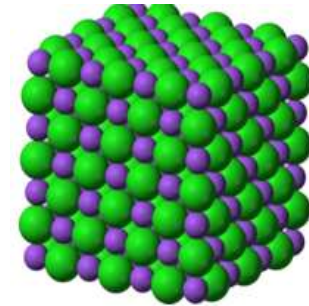
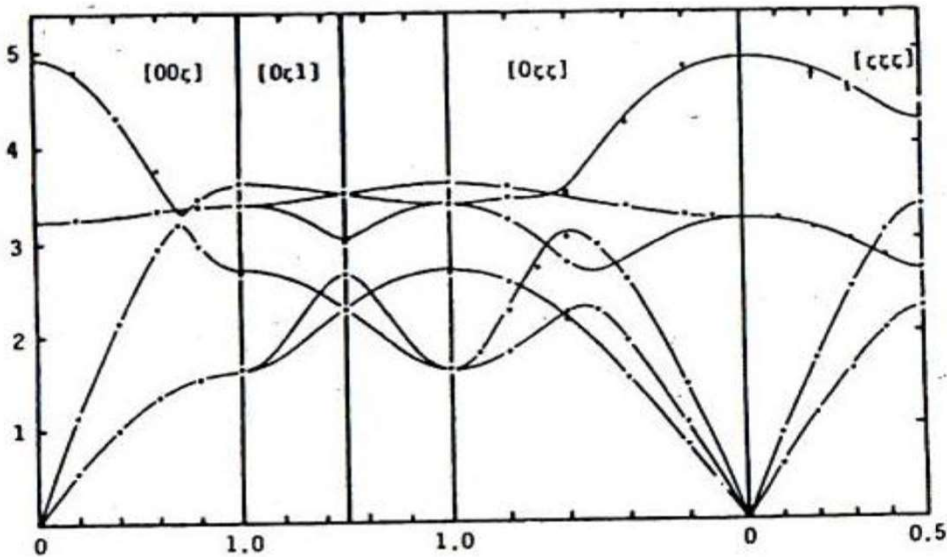
From Springer Materials: Landholt Boernstein Database

Next nearest neighbors (simple cubic)

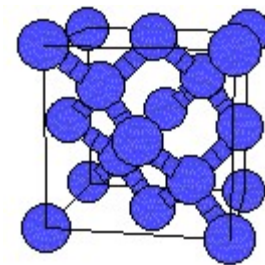
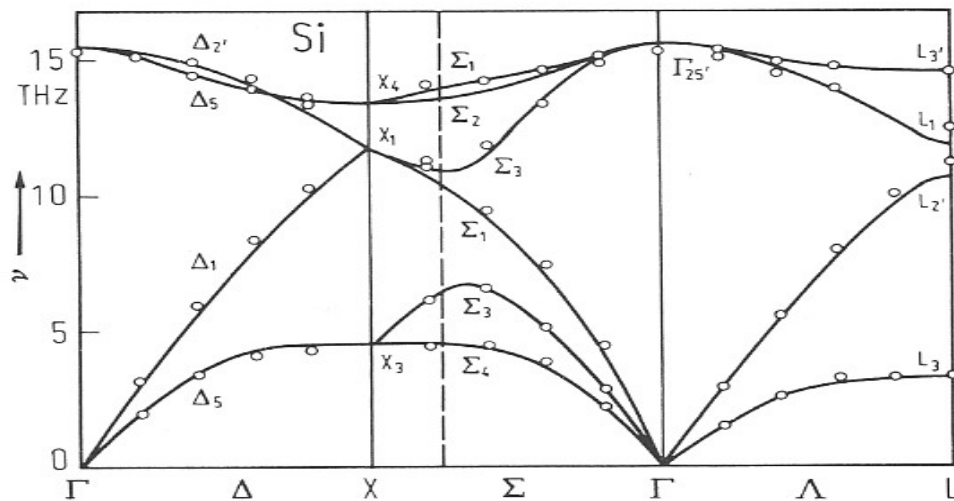


Sometimes the 5th neighbors are included.

Two atoms per primitive unit cell

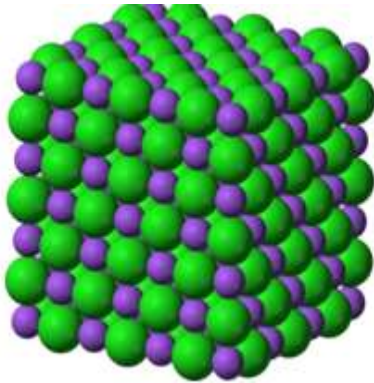


NaCl



Si

NaCl



2 atoms/unit cell

6 equations

3 acoustic and
3 optical branches

x - Richtung:

$$M_1 \frac{d^2 u_{nml}^x}{dt^2} = C \left(-2u_{nml}^x + v_{(n-1)m(l-1)}^x + v_{n(m-1)l}^x \right)$$

$$M_2 \frac{d^2 v_{nml}^x}{dt^2} = C \left(-2v_{nml}^x + u_{(n+1)m(l+1)}^x + u_{n(m+1)l}^x \right)$$

y - Richtung:

$$M_1 \frac{d^2 u_{nml}^y}{dt^2} = C \left(-2u_{nml}^y + v_{(n-1)(m-1)l}^y + v_{nm(l-1)}^y \right)$$

$$M_2 \frac{d^2 v_{nml}^y}{dt^2} = C \left(-2v_{nml}^y + u_{(n+1)(m+1)l}^y + u_{nm(l+1)}^y \right)$$

z - Richtung:

$$M_1 \frac{d^2 u_{nml}^z}{dt^2} = C \left(-2u_{nml}^z + v_{n(m-1)(l-1)}^z + v_{(n-1)ml}^z \right)$$

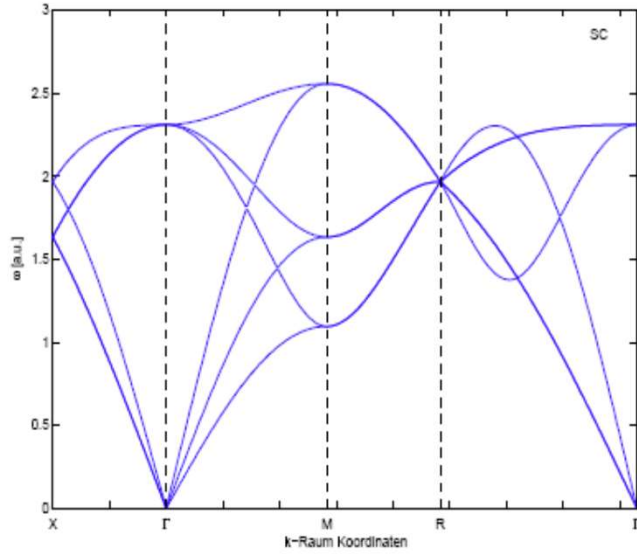
$$M_2 \frac{d^2 v_{nml}^z}{dt^2} = C \left(-2v_{nml}^z + u_{n(m+1)(l+1)}^z + u_{(n+1)ml}^z \right)$$

$$u_{nml}^x = u_{\vec{k}}^x \exp\left(i\left(\vec{k} \cdot \vec{a}_1 + \vec{k} \cdot \vec{a}_2 + \vec{k} \cdot \vec{a}_3 - \omega t\right)\right) \quad v_{nml}^x = v_{\vec{k}}^x \exp\left(i\left(\vec{k} \cdot \vec{a}_1 + \vec{k} \cdot \vec{a}_2 + \vec{k} \cdot \vec{a}_3 - \omega t\right)\right)$$

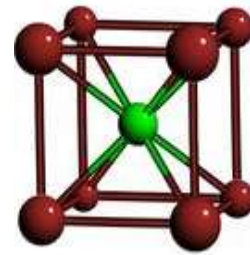
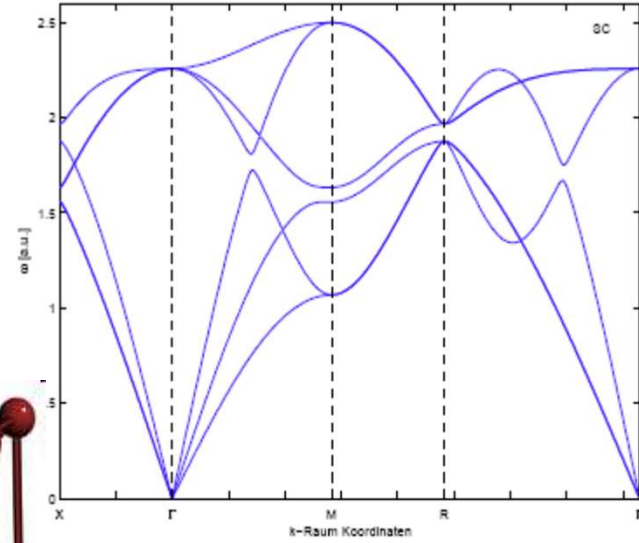
CsCl

Hannes Brandner

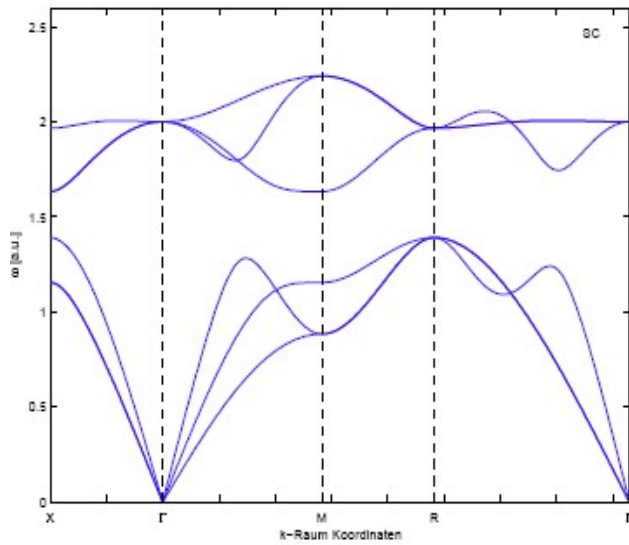
$$M_1 = M_2$$



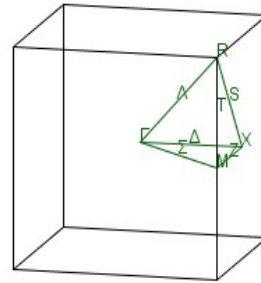
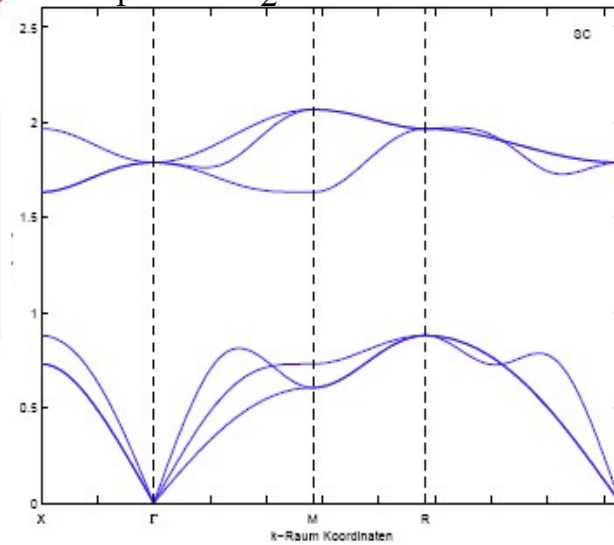
$$M_1 = 1.1 M_2$$



$$M_1 = 2M_2$$



$$M_1 = 5M_2$$



3 dimensions

N atoms

$3N$ normal modes

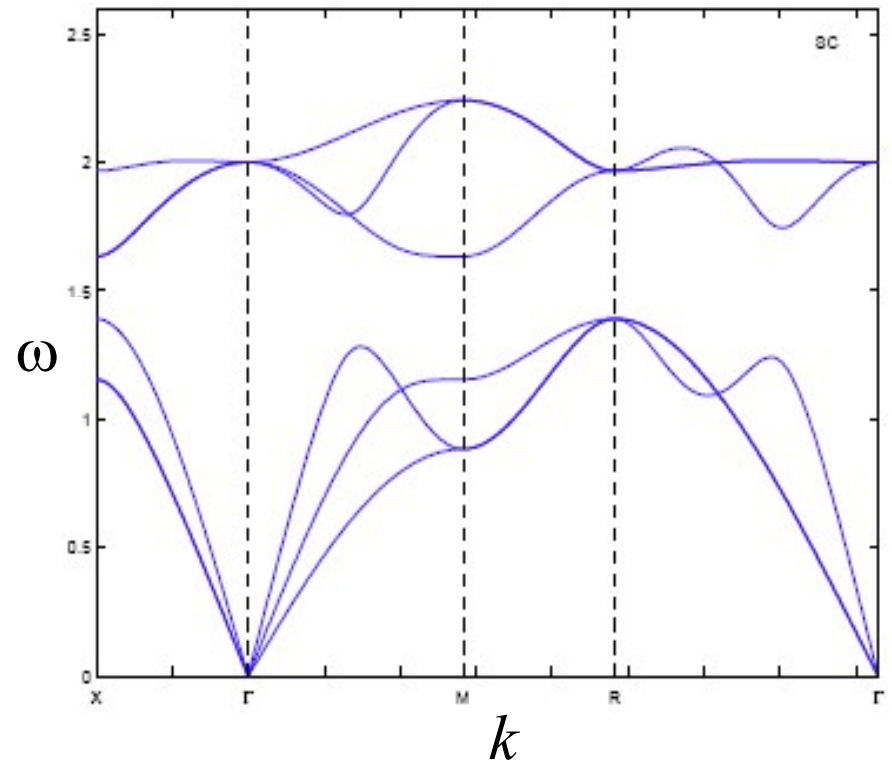
p atoms per unit cell

N/p unit cells = k vectors

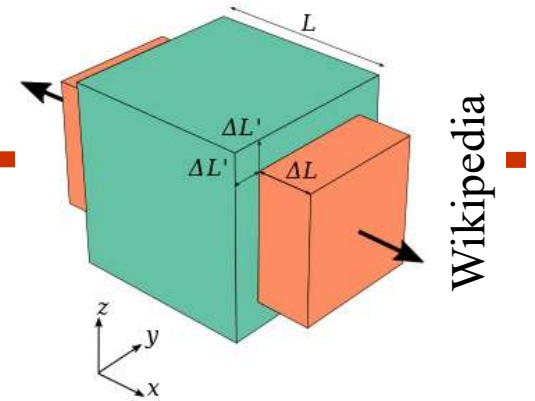
$3p$ branches to the dispersion relation

3 acoustic modes (1 longitudinal, 2 transverse)

$3p - 3$ optical modes



Poisson's ratio



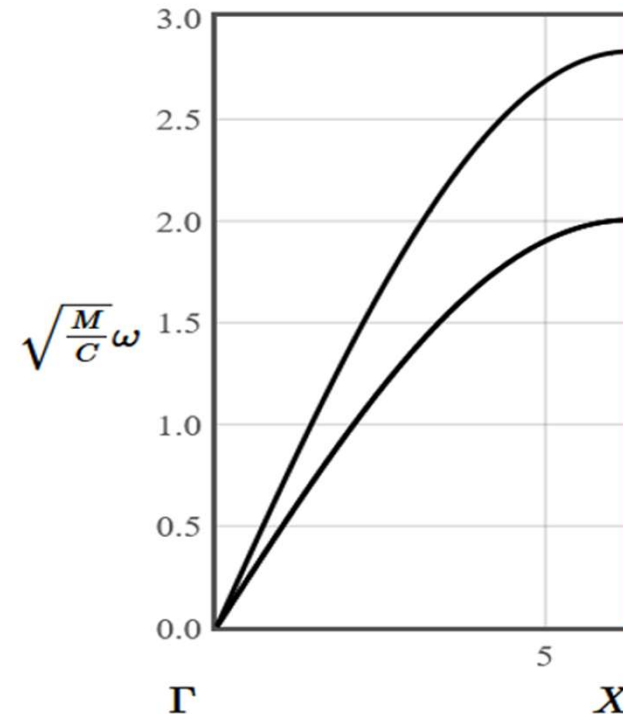
E - Elastic constant

ν - Poisson's ratio

ρ - density

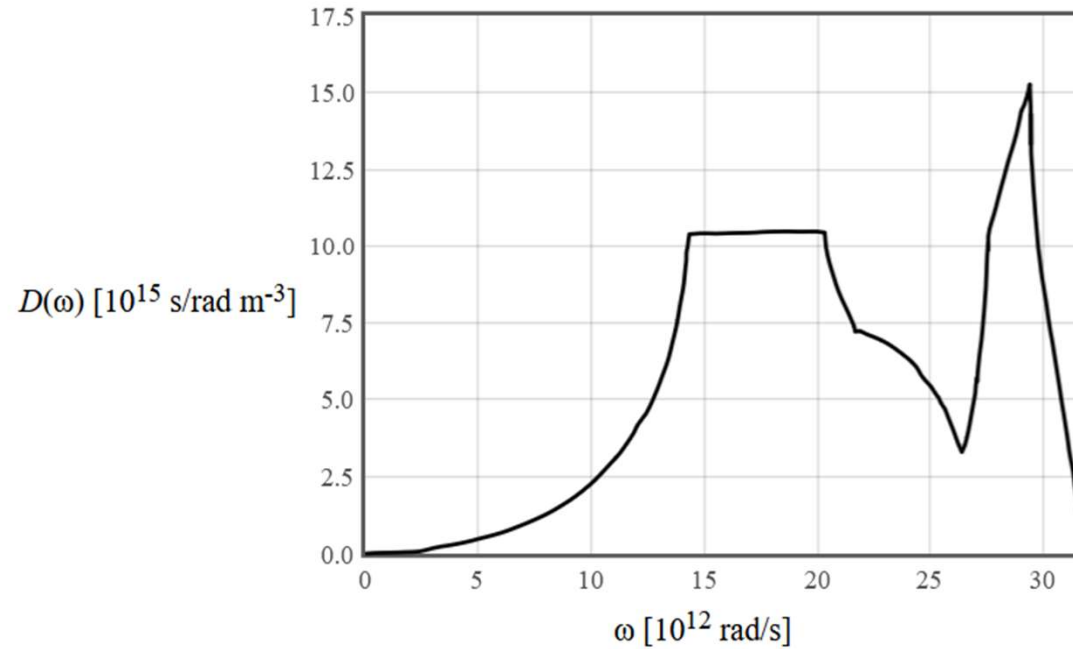
$$c_T = \sqrt{\frac{E(1-\nu)}{\rho(1-\nu-2\nu^2)}}$$

$$c_L = \sqrt{\frac{E}{2\rho(1+\nu)}}$$



If the density is known, you can determine E and ν .

Phonon density of states for fcc silver



The atomic density is taken to be $5.86 \times 10^{28} \text{ m}^{-3}$. Each atom has three degrees of freedom so the integral over all frequencies is $3 \times 5.86 \times 10^{28} \text{ m}^{-3}$. The data is from [doi: 10.1007/b19988](https://doi.org/10.1007/b19988).

T = 296 K

ω [rad/s]	$D(\omega)$ [s rad ⁻¹ m ⁻³]
0.0000	0.0000
5.7327e+10	6.8161e+12
4.0123e+11	2.3856e+13
7.4510e+11	3.0672e+13
1.0890e+12	3.4080e+13
1.4233e+12	4.0897e+13
1.7624e+12	5.1121e+13
2.0967e+12	5.7937e+13
2.4120e+12	7.4977e+13
2.7177e+12	1.2610e+14
3.0379e+12	1.8744e+14
3.3723e+12	2.3516e+14
3.7163e+12	2.7261e+14



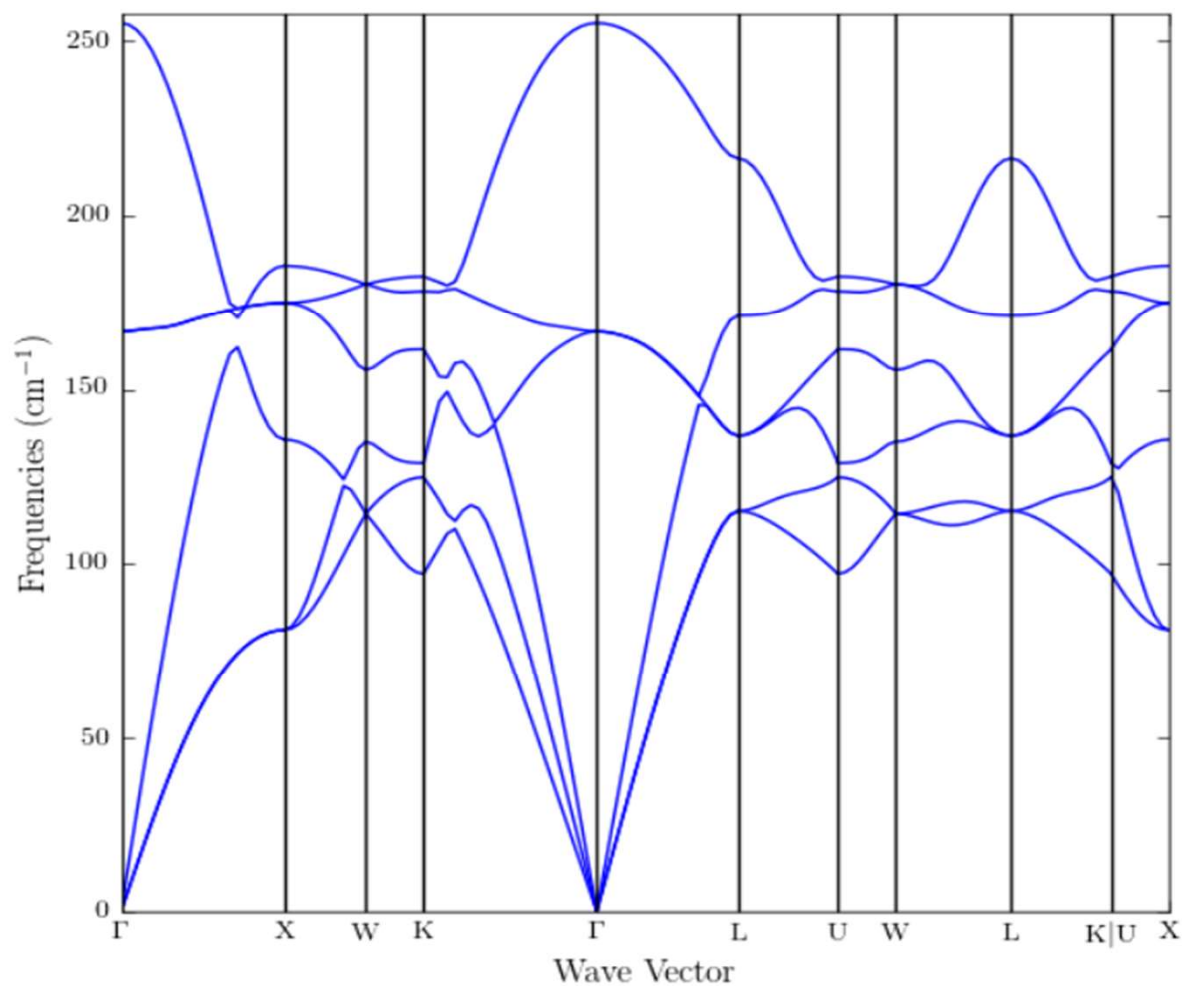
Vibrational Properties

Reference for phonon calculations and visualization:

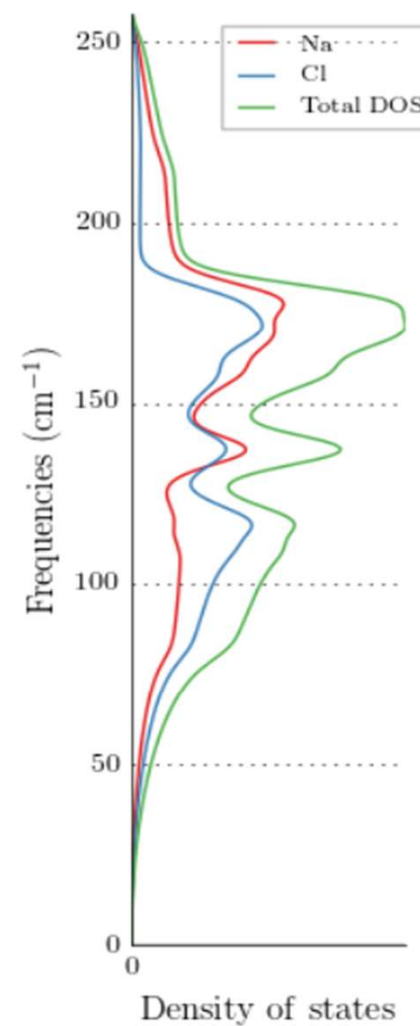


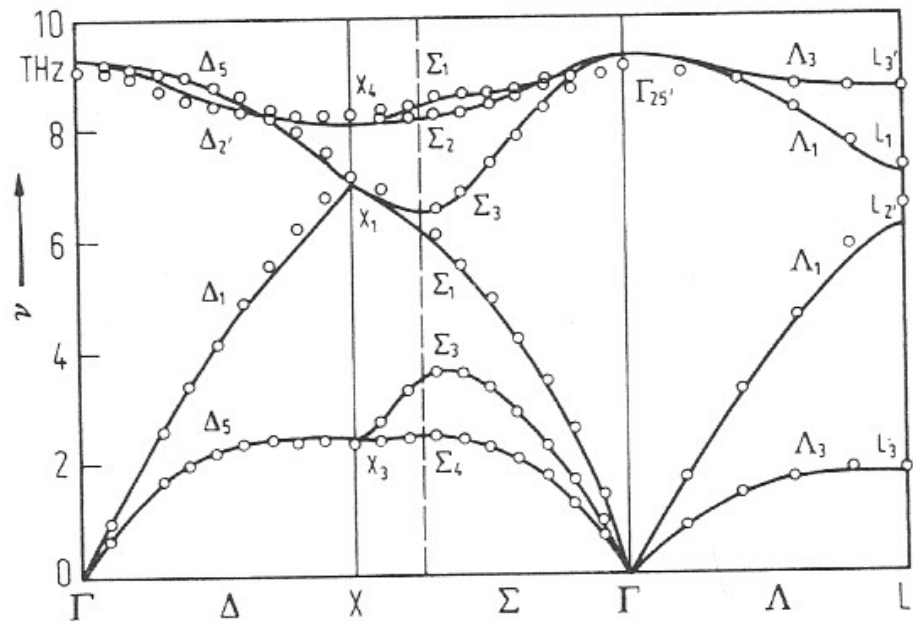
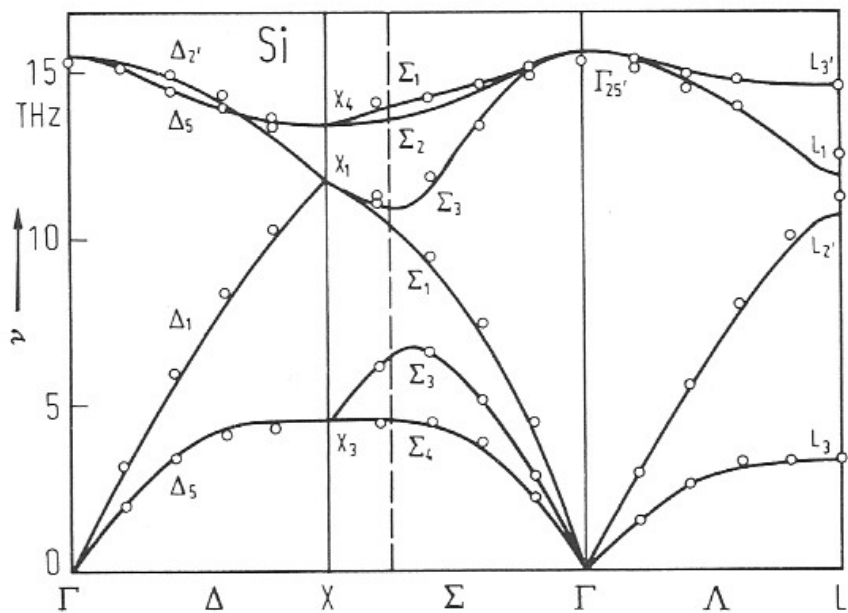
[Visualize with phononwebsite](#)

Phonon dispersion

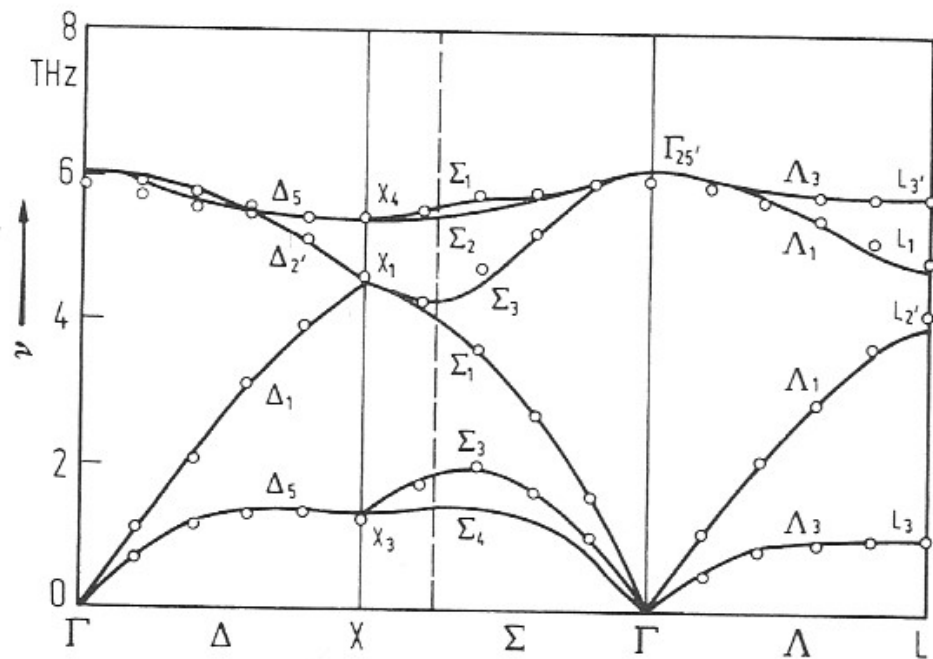
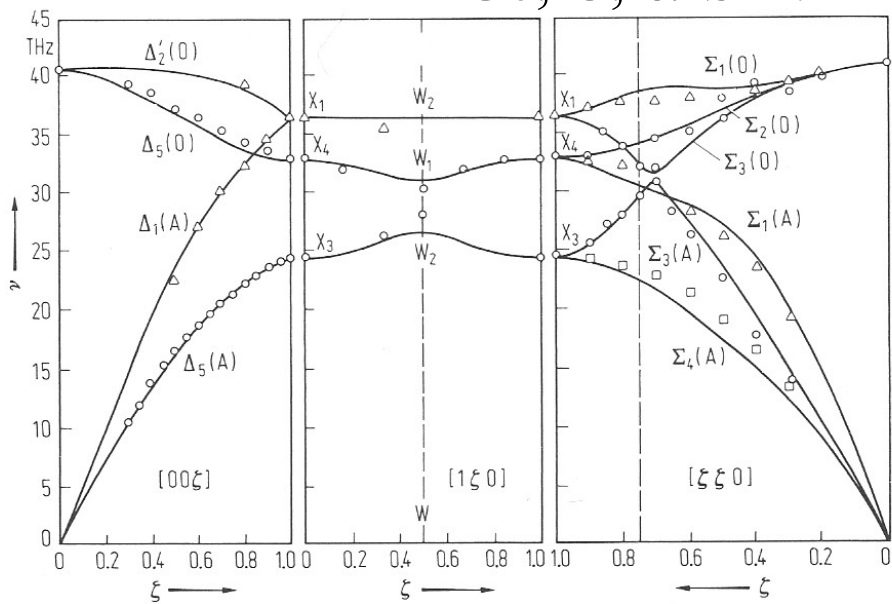


Density of States



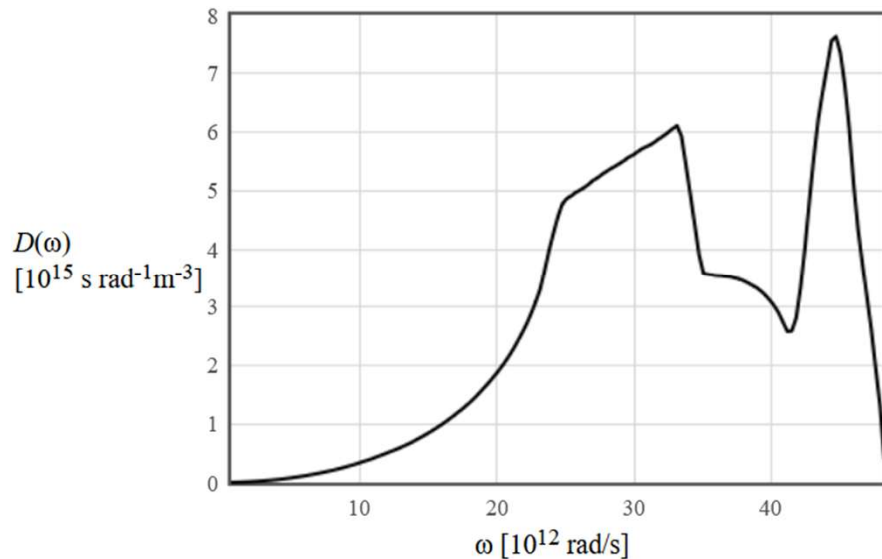


Ge, C, α -Sn ?

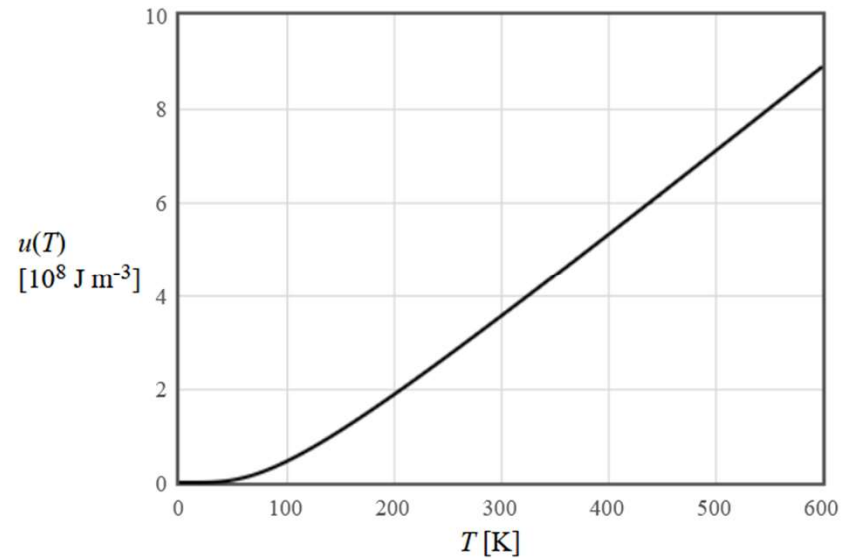


Density of states \rightarrow Internal energy density

$$u(T) = \int_0^{\infty} \frac{\hbar\omega D(\omega)}{\exp\left(\frac{\hbar\omega}{k_B T}\right) - 1} d\omega$$



DoS \rightarrow $u(T)$



<http://lampx.tugraz.at/~hadley/ss1/phonons/table/dos2ut.html>