

Carrier Transport

- Ballistic transport
- Drift
- Diffusion
- Generation and recombination
- The continuity equation
- High field effects

Review

n-type

$$N_D > N_A \quad n = N_D - N_A$$

$$p = \frac{n_i^2}{n} = \frac{n_i^2}{N_D - N_A}$$

$$n = N_c \exp\left(\frac{E_F - E_c}{k_B T}\right) = N_D - N_A$$



$$E_F = E_c - k_B T \ln\left(\frac{N_c}{N_D - N_A}\right)$$

p-type

$$N_A > N_D \quad p = N_A - N_D$$

$$n = \frac{n_i^2}{p} = \frac{n_i^2}{N_A - N_D}$$

$$p = N_v \exp\left(\frac{E_v - E_F}{k_B T}\right) = N_A - N_D$$



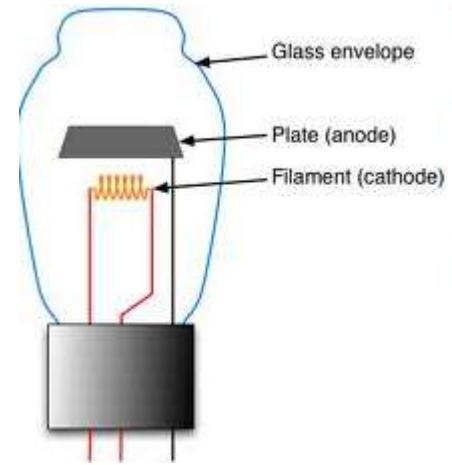
$$E_F = E_v + k_B T \ln\left(\frac{N_v}{N_A - N_D}\right)$$

Ballistic transport

$$\vec{F} = m\vec{a} = -e\vec{E} = m \frac{d\vec{v}}{dt}$$

$$\vec{v} = \frac{-e\vec{E}t}{m} + \vec{v}_0$$

$$\vec{x} = \frac{-e\vec{E}t^2}{2m} + \vec{v}_0 t + \vec{x}_0$$



Electrons moving in an electric field follow parabolic trajectories like a ball in a gravitational field.

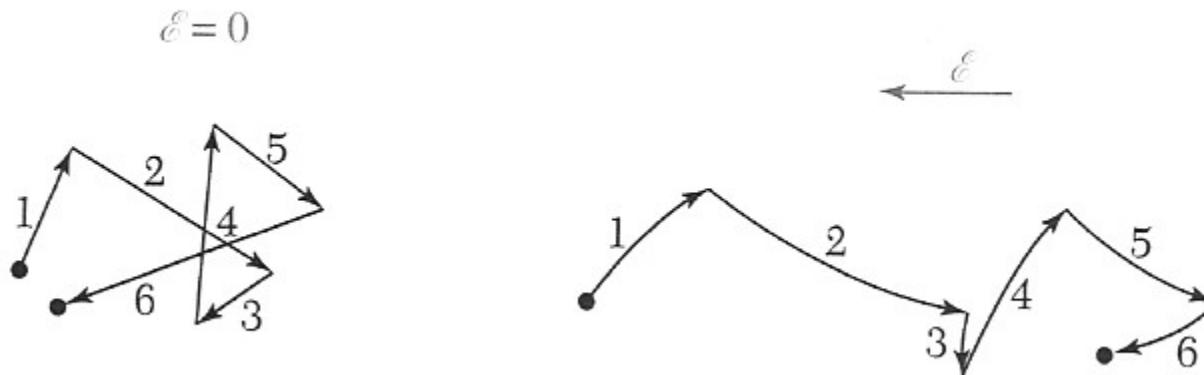
Drift

The electrons scatter and change direction after a time τ_{sc} .

Classical equipartition: $\frac{1}{2}mv_{th}^2 = \frac{3}{2}k_B T$

At 300 K, $v_{th} \sim 10^7$ cm/s.

mean free path: $\ell = v_{th}\tau_{sc} \sim 10$ nm ~ 200 atoms



Drift (diffusive transport)

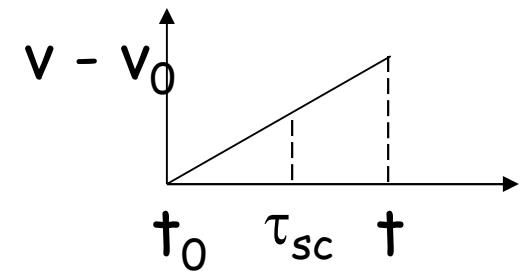
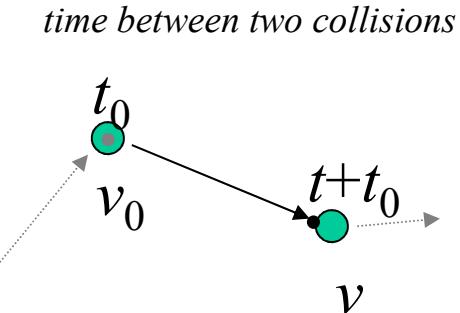
$$\vec{F} = -e\vec{E} = m^* \vec{a} = m^* \frac{d\vec{v}}{dt}$$

$$\vec{v} = \vec{v}_0 - \frac{e\vec{E}}{m^*} (t - t_0)$$

$$\langle v_0 \rangle = 0$$

$\langle t - t_0 \rangle = \tau_{sc}$ < average time between scattering events

$$\vec{v}_d = \frac{-e\vec{E}\tau_{sc}}{m^*} = \frac{-e\vec{E}\ell}{m^* v}$$



drift velocity:

$$\vec{v}_{d,n} = -\mu_n \vec{E}$$

$$\vec{v}_{d,p} = \mu_p \vec{E}$$

Drift

drift velocity:

$$\vec{v}_{d,n} = -\mu_n \vec{E} \quad \vec{v}_{d,p} = \mu_p \vec{E}$$

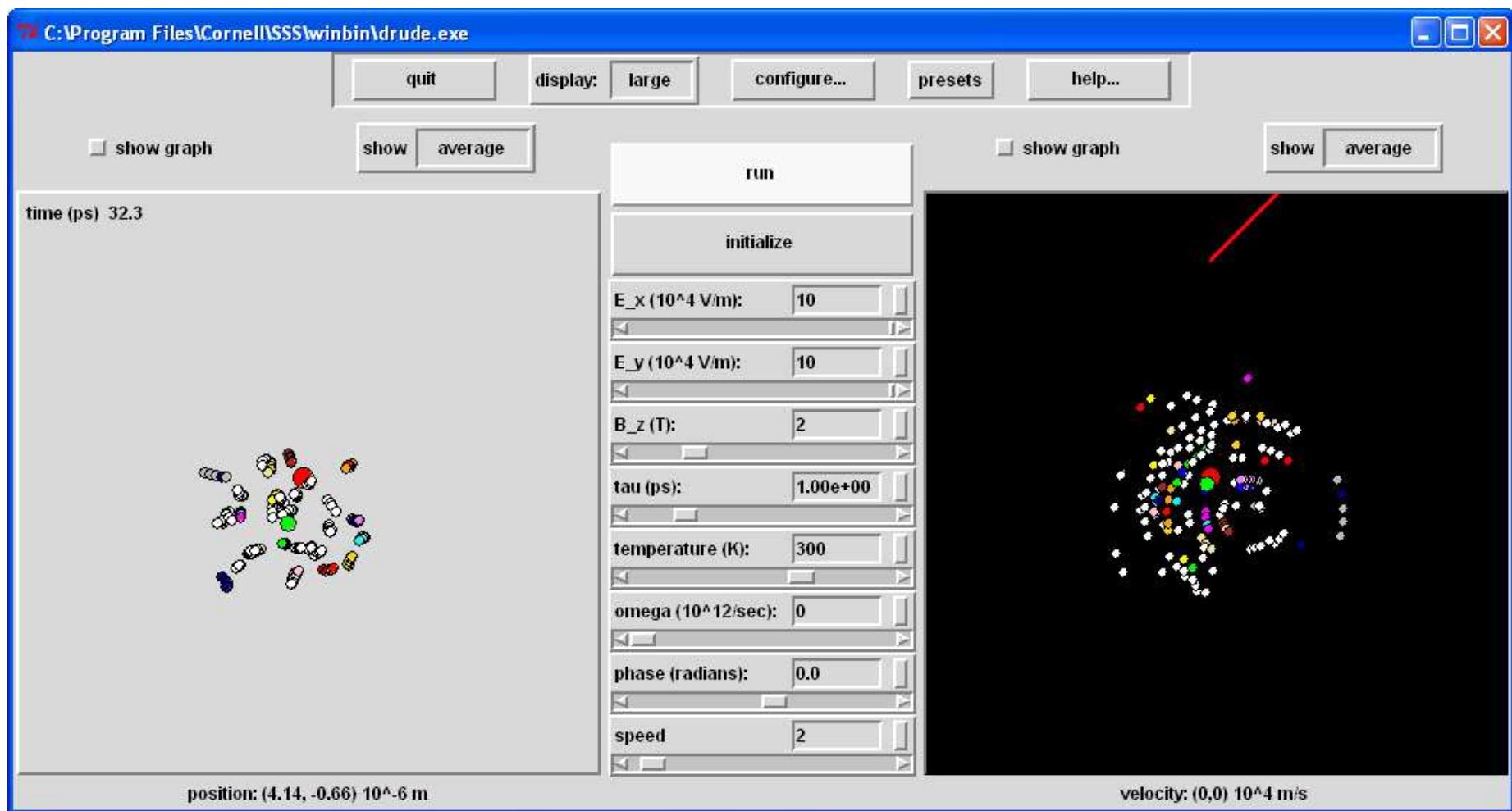
$$\vec{j} = -ne\vec{v}_{d,n} + pe\vec{v}_{d,p} = (ne\mu_n + pe\mu_p) \vec{E} = \sigma \vec{E}$$

$$\mu = \frac{-e\tau_{sc}}{m^*} = \frac{-e\ell}{m^* v}$$

for Si:

$$\begin{aligned}\mu_n &= 1500 \text{ cm}^2/\text{Vs} \\ \mu_p &= 450 \text{ cm}^2/\text{Vs}\end{aligned}$$

For $E = 1000 \text{ V/cm}$ $v_d = 10^6 \text{ cm/s}$



Drift

Solid state electronic devices, Streetman and Banerjee

		E_g (eV)	μ_n (cm ² /V-s)	μ_p (cm ² /V-s)	m_n^*/m_o (m_l, m_h)	m_p^*/m_o (m_{lh}, m_{hh})	a (Å)	ϵ_r	Density (g/cm ³)	Melting point (°C)
Si	(i/D)	1.11	1350	480	0.98, 0.19	0.16, 0.49	5.43	11.8	2.33	1415
Ge	(i/D)	0.67	3900	1900	1.64, 0.082	0.04, 0.28	5.65	16	5.32	936
SiC (α)	(i/W)	2.86	500	—	0.6	1.0	3.08	10.2	3.21	2830
AlP	(i/Z)	2.45	80	—	—	0.2, 0.63	5.46	9.8	2.40	2000
AlAs	(i/Z)	2.16	1200	420	2.0	0.15, 0.76	5.66	10.9	3.60	1740
AlSb	(i/Z)	1.6	200	300	0.12	0.98	6.14	11	4.26	1080
GaP	(i/Z)	2.26	300	150	1.12, 0.22	0.14, 0.79	5.45	11.1	4.13	1467
GaAs	(d/Z)	1.43	8500	400	0.067	0.074, 0.50	5.65	13.2	5.31	1238
GaN	(d/Z, W)	3.4	380	—	0.19	0.60	4.5	12.2	6.1	2530
GaSb	(d/Z)	0.7	5000	1000	0.042	0.06, 0.23	6.09	15.7	5.61	712
InP	(d/Z)	1.35	4000	100	0.077	0.089, 0.85	5.87	12.4	4.79	1070
InAs	(d/Z)	0.36	22600	200	0.023	0.025, 0.41	6.06	14.6	5.67	943
InSb	(d/Z)	0.18	10 ⁵	1700	0.014	0.015, 0.40	6.48	17.7	5.78	525
ZnS	(d/Z, W)	3.6	180	10	0.28	—	5.409	8.9	4.09	1650*
ZnSe	(d/Z)	2.7	600	28	0.14	0.60	5.671	9.2	5.65	1100*
ZnTe	(d/Z)	2.25	530	100	0.18	0.65	6.101	10.4	5.51	1238*
CdS	(d/W, Z)	2.42	250	15	0.21	0.80	4.137	8.9	4.82	1475
CdSe	(d/W)	1.73	800	—	0.13	0.45	4.30	10.2	5.81	1258
CdTe	(d/Z)	1.58	1050	100	0.10	0.37	6.482	10.2	6.20	1098
PbS	(i/H)	0.37	575	200	0.22	0.29	5.936	17.0	7.6	1119
PbSe	(i/H)	0.27	1500	1500	—	—	6.147	23.6	8.73	1081
PbTe	(i/H)	0.29	6000	4000	0.17	0.20	6.452	30	8.16	925

$$\vec{v}_{d,n} = -\mu_n \vec{E} \quad \vec{v}_{d,p} = \mu_p \vec{E}$$

$$\vec{j} = -ne\vec{v}_{d,n} + pe\vec{v}_{d,p} = (ne\mu_n + pe\mu_p) \vec{E} = \sigma \vec{E}$$

Matthiessen's rule

$$\frac{1}{\tau_{sc}} = \frac{1}{\tau_{sc,lattice}} + \frac{1}{\tau_{sc,impurity}}$$

↑ ↗
phonons, temperature dependent mostly temperature independent

$$\frac{1}{\mu} = \frac{1}{\mu_{lattice}} + \frac{1}{\mu_{impurity}}$$

$$\sigma = \frac{1}{\rho} = ne\mu_n + pe\mu_p$$

↗
doping increases the conductivity
by increasing the carrier density
but decreases the mobility

Mobility calculator

$$\mu = \mu_{min} + \frac{\mu_{max} - \mu_{min}}{1 + (N/N_{ref})^\gamma}$$

For Electrons:

$$\mu_{min} = 47 \left(\frac{T}{300} \right)^{-1.23} \frac{\text{cm}^2}{\text{Vs}}$$

$$\Delta\mu = \mu_{max} - \mu_{min} = 1373 \left(\frac{T}{300} \right)^{-2.38} \frac{\text{cm}^2}{\text{Vs}}$$

$$N_{ref} = 1,05 \cdot 10^{17} \left(\frac{T}{300} \right)^{3.65} \text{ cm}^{-3}; \gamma = 0,68 \left(\frac{T}{300} \right)^{-0.32}$$

For Holes:

$$\mu_{min} = 36 \left(\frac{T}{300} \right)^{-0.87} \frac{\text{cm}^2}{\text{Vs}}$$

$$\Delta\mu = \mu_{max} - \mu_{min} = 438 \left(\frac{T}{300} \right)^{-2.01} \frac{\text{cm}^2}{\text{Vs}}$$

$$N_{ref} = 2,85 \cdot 10^{17} \left(\frac{T}{300} \right)^{2.93} \text{ cm}^{-3}; \gamma = 0,65 \left(\frac{T}{300} \right)^{0.26}$$

INPUTS

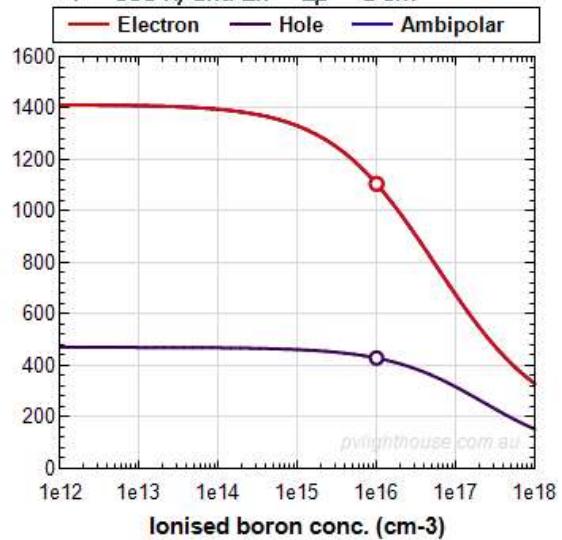
Semiconductor material	c-silicon	Excess electron conc.	Δn	1	cm ⁻³
Dopant atom	boron	Excess hole conc.	Δp	1	cm ⁻³
Ionised dopant conc.	1E+16	Electron eff. lifetime	$\tau_{eff,e}$	1E-4	s
Temperature	300 K	Hole eff. lifetime	$\tau_{eff,h}$	1E-4	s

OUTPUTS

Carrier concentrations			Carrier mobility etc.		
Equilibrium n_0, p_0 (cm ⁻³)	Excess $\Delta n, \Delta p$ (cm ⁻³)	Net n, p (cm ⁻³)	Mobility μ_e, μ_h, μ_a (cm ² V ⁻¹ s ⁻¹)	Diffusivity D_e, D_h, D_a (cm ² s ⁻¹)	Diff Length L_e, L_h, L_a (cm)
ions	9300	1.0	9300	1107	28.61
	1.0E+16	1.0	1.0E+16	429.3	11.10
olar				1107	28.61
					5.349E-2
					3.331E-2

Resistivity ($\Omega\text{-cm}$)	
rium	ρ_0 1.454
i-state	ρ 1.454

Mobility vs ionised dopant concentration
for boron-doped c-silicon with
 $T = 300 \text{ K}$, and $\Delta n = \Delta p = 1 \text{ cm}^{-3}$



ture inputs

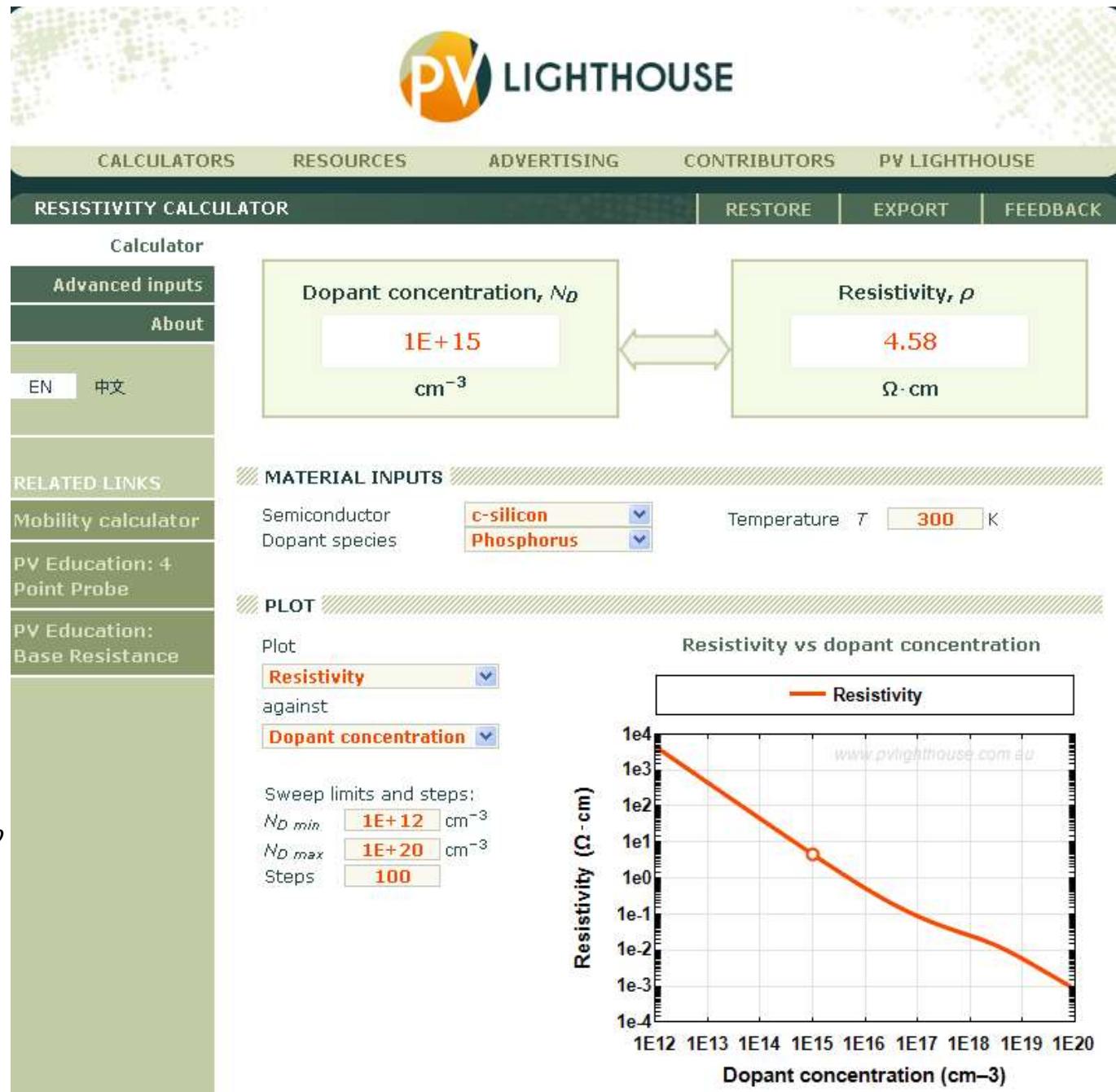
Mobility	1E+12	cm ⁻³
Ionised dopant conc.	1E+18	cm ⁻³
points	50	

Resistivity calculator

$$n = N_c \exp\left(\frac{E_F - E_c}{k_B T}\right)$$

$$p = N_v \exp\left(\frac{E_v - E_F}{k_B T}\right)$$

$$\sigma = \frac{1}{\rho} = ne\mu_n + pe\mu_p$$

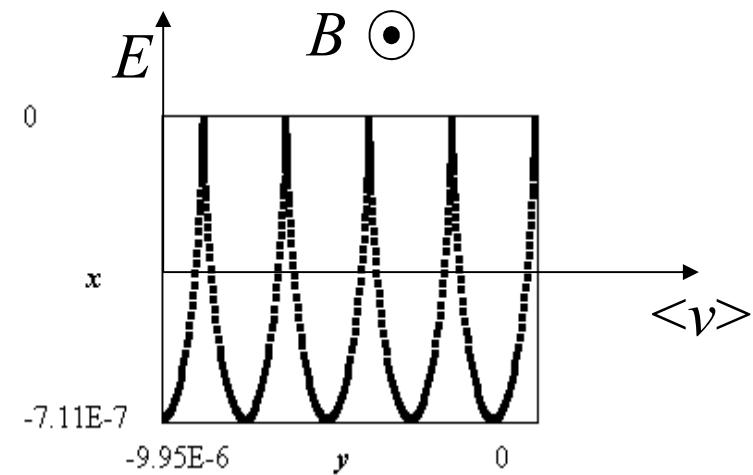


<http://www.pvlighthouse.com.au/calculators/Resistivity%20calculator/Resistivity%20calculator.aspx>

Crossed E and B fields

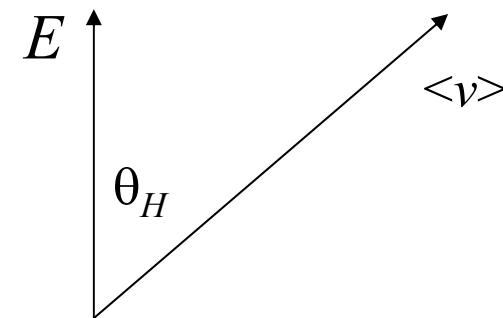
Ballistic transport

$$\vec{F} = m\vec{a} = -e(\vec{E} + \vec{v} \times \vec{B})$$



Diffusive transport

Hall angle:



$$\theta_H = \tan^{-1} \left(-\frac{eB_z \tau_{sc}}{m^*} \right)$$

Magnetic field (diffusive transport)

$$\vec{F} = m\vec{a} = -e\vec{E} = e \frac{\vec{v}_d}{\mu}$$

$$\vec{F} = m\vec{a} = -e(\vec{E} + \vec{v}_d \times \vec{B}) = e \frac{\vec{v}_d}{\mu}$$

If B is in the z -direction, the three components of the force are

$$-\mu(E_x + v_{dy}B_z) = v_{dx}$$

$$-\mu(E_y - v_{dx}B_z) = v_{dy}$$

$$-\mu E_z = v_{dz}$$

Magnetic field

$$v_{d,x} = -\mu E_x - \mu B_z v_{d,y}$$

$$v_{d,y} = -\mu E_y + \mu B_z v_{d,x}$$

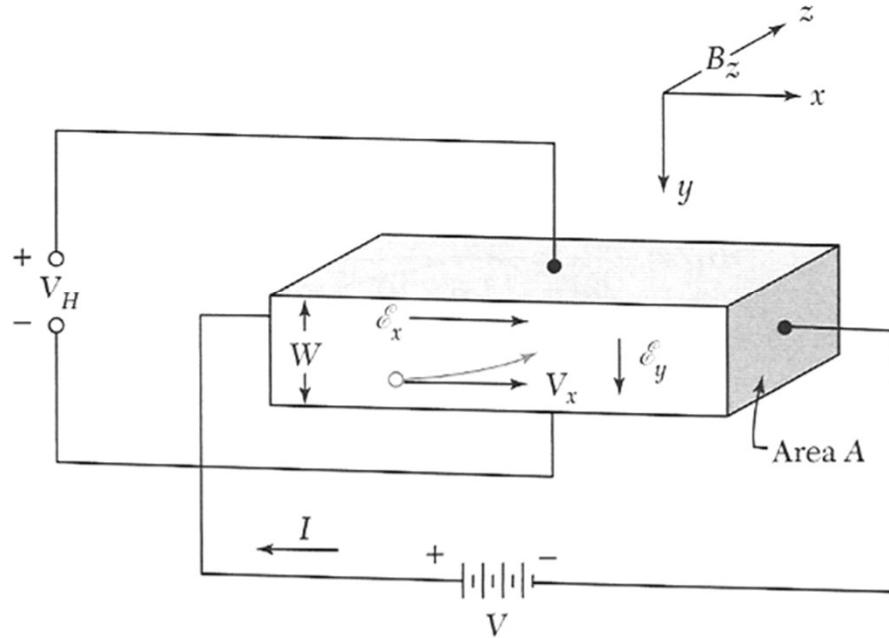
$$v_{d,z} = -\mu E_z$$

If $E_y = 0$,

$$v_{d,y} = -\mu B_z v_{d,x}$$

$$\tan \theta_H = -\mu B_z$$

The Hall Effect (diffusive regime)



$$v_{d,x} = -\mu E_x - \mu B_z v_{d,y}$$

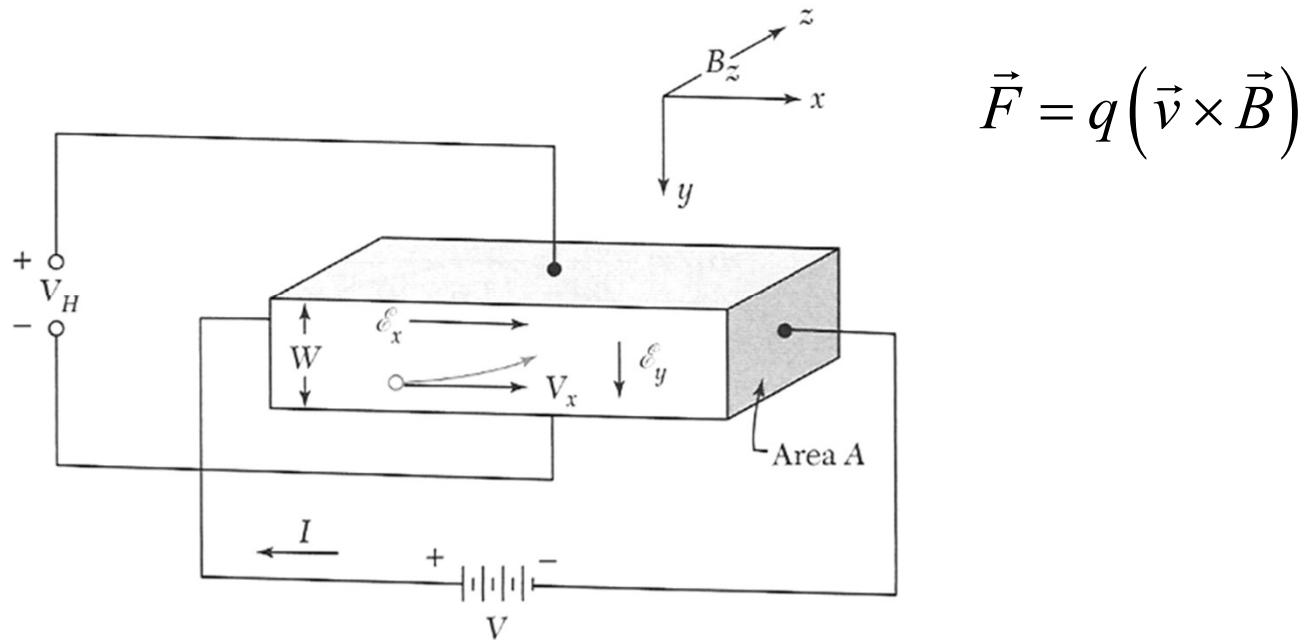
$$v_{d,y} = -\mu E_y + \mu B_z v_{d,x}$$

$$v_{d,z} = -\mu E_z$$

If $v_{d,y} = 0$,

$$E_y = v_x B_z = V_H/W = R_H j_x B_z \quad V_H = \text{Hall voltage}, R_H = \text{Hall Constant}$$

The Hall Effect (diffusive regime)



$$\vec{F} = q(\vec{v} \times \vec{B})$$

$$E_y = v_x B_z = V_H/W = R_H j_x B_z$$

V_H = Hall voltage, R_H = Hall Constant

$$j_x = I/A$$

$$v_x = -j_x/ne \quad \text{for n-type}$$

$$v_x = j_x/pe \quad \text{for p-type}$$

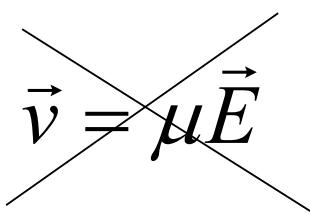
$$R_H = -1/ne \quad \text{for n-type}$$

$$R_H = 1/pe \quad \text{for p-type}$$

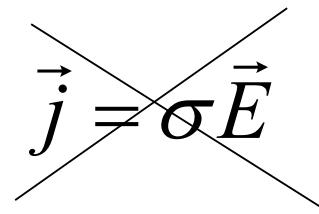
Ballistic transport in transistors

The mean free path ~ 100 nm > gate length ~ 20 nm

v not proportional to E

$$\vec{v} = \mu \vec{E}$$


j not proportional to E

$$\vec{j} = \sigma \vec{E}$$


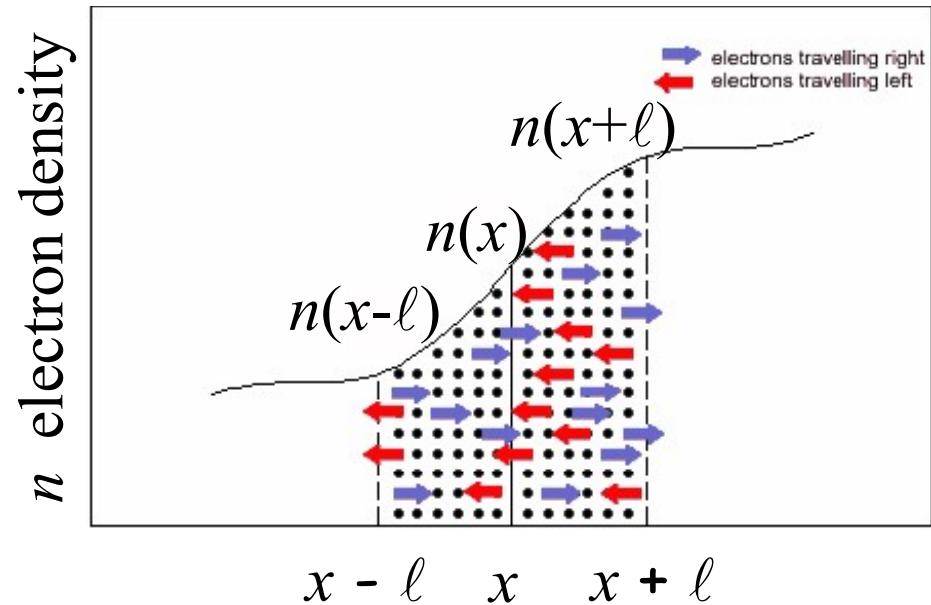
nonlocal response

Electrons bend in a magnetic field like they do in vacuum.

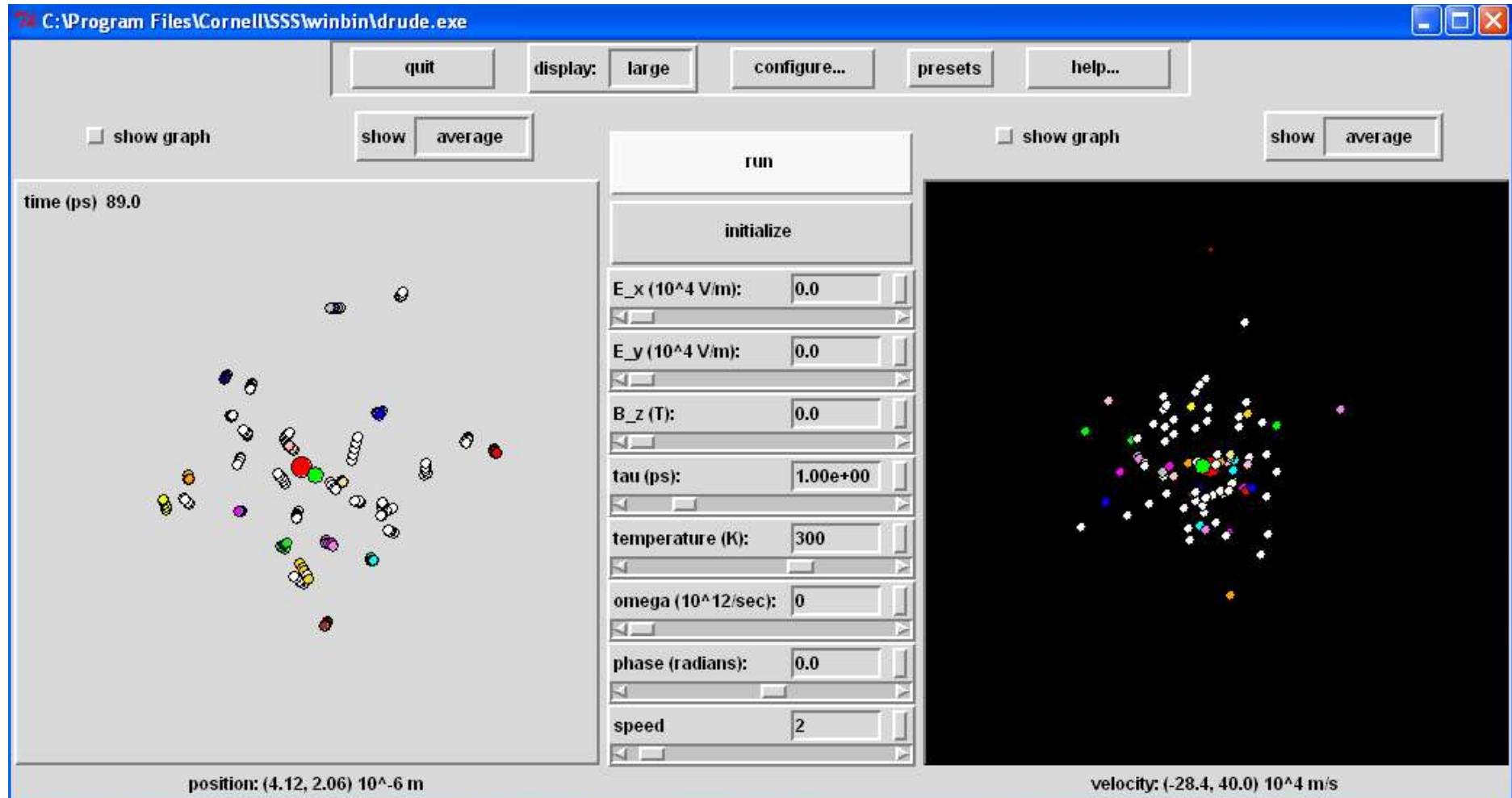
Diffusion

$$j_{n,diff} = |e| D_n \frac{dn}{dx}$$

$$j_{p,diff} = -|e| D_p \frac{dp}{dx}$$



Diffusion is from high concentration to low concentration.



If no forces are applied, the electrons diffuse.
The average velocity moves against an electric field.
In just a magnetic field, the average velocity is zero.
In an electric and magnetic field, the electrons move in a straight line at the Hall angle.

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Diffusion current

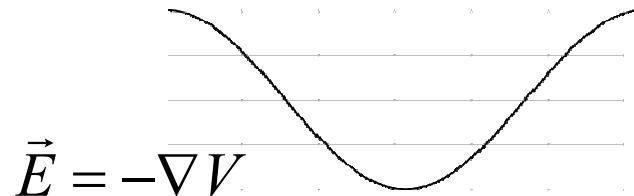
The phosphorous concentration in a region of a silicon crystal varies linearly from a concentration of $5\text{E}+16 \text{ cm}^{-3}$ at $x = 0$ to a concentration of $7\text{E}+17 \text{ cm}^{-3}$ at $x = 1 \text{ mm}$. The diffusion constant for electrons is $22.5 \text{ cm}^2/\text{s}$, the diffusion constant for holes is $5.2 \text{ cm}^2/\text{s}$, and the temperature is 300 K . What is the diffusion current density in the positive x -direction?

$J =$ A/cm²

[Submit answer](#)[Clear](#)

$$J_{n,diff} = |e| D_n \frac{dn}{dx}$$

Einstein relation



$$\vec{E} = -\nabla V$$

$$n(x) = A \exp\left(\frac{-eV(x)}{k_B T}\right)$$
 Boltzmann factor

In equilibrium, drift = diffusion

$$-en\mu\vec{E} + D\nabla n = 0$$

$$\nabla n = -\frac{1}{k_B T} A \exp\left(\frac{-eV_{pot}}{k_B T}\right) \nabla V = -\frac{n}{k_B T} \nabla V = \frac{n\vec{E}}{k_B T}$$

$$-en\mu\vec{E} + D \frac{n\vec{E}}{k_B T} = 0$$

$$D = \frac{\mu k_B T}{e}$$

Über die von der molekularkinetischen Theorie der Wärme geforderte Bewegung von in ruhenden Flüssigkeiten suspendierten Teilchen

Current Density Equations

$$\vec{j}_n = -ne\mu_n \vec{E} + eD_n \nabla n$$
$$\vec{j}_p = pe\mu_p \vec{E} - eD_p \nabla p$$

Drift
↓

Diffusion
↖

$$\vec{j}_{total} = \vec{j}_n + \vec{j}_p$$

Current Density Equations

note: electron and hole currents have same direction

electric current = charge \times particle flow

drift



diffusion

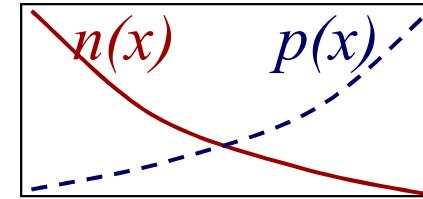


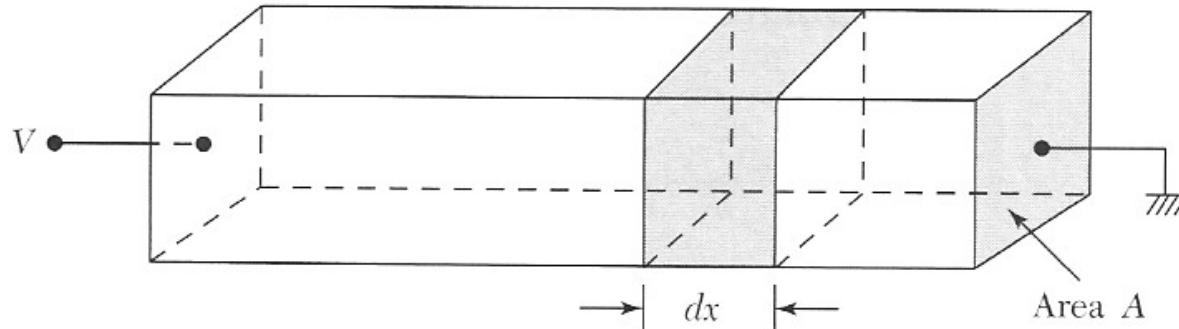
Diagram showing electron drift current. An arrow labeled "current" points to the right. To its left is another arrow labeled "flow" pointing to the left, and above it is a symbol consisting of a minus sign inside a circle: \ominus . Below the arrows is the equation $j_e = -e \times \text{flow}$.

Diagram showing electron diffusion current. An arrow labeled "current" points to the right. To its left is another arrow labeled "flow" pointing to the left, and above it is a symbol consisting of a minus sign inside a circle: \ominus . Below the arrows is the equation $-D_n(dn/dx)$.

Diagram showing hole drift current. An arrow labeled "current" points to the right. To its left is another arrow labeled "flow" pointing to the right, and below it is a symbol consisting of a plus sign inside a square: \oplus . Below the arrows is the equation $j_p = e \times \text{flow}$.

Diagram showing hole diffusion current. An arrow labeled "current" points to the right. To its left is another arrow labeled "flow" pointing to the right, and below it is a symbol consisting of a plus sign inside a square: \oplus . Below the arrows is the equation $D_p(dp/dx)$.

Continuity equations

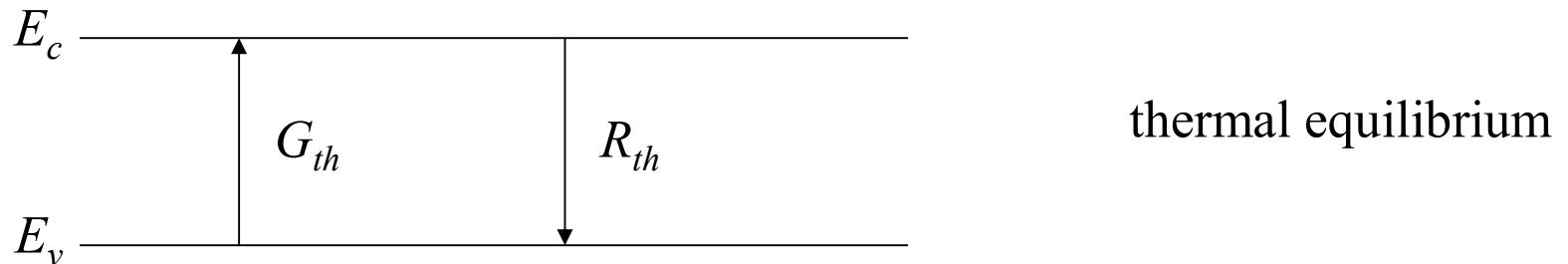


$$\frac{\partial n}{\partial t} = \frac{1}{e} \nabla \cdot \vec{j}_n + G_n - R_n$$

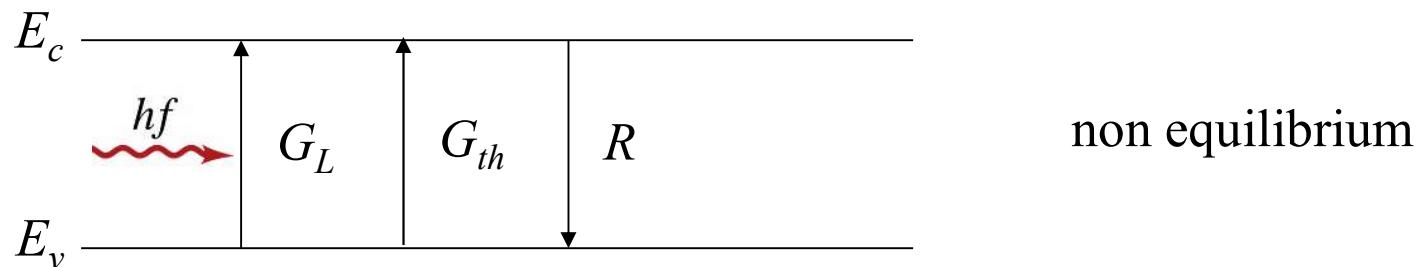
$$\frac{\partial p}{\partial t} = -\frac{1}{e} \nabla \cdot \vec{j}_p + G_p - R_p$$

j_n and j_p consist of drift and diffusion terms

Generation and Recombination



Shining light on a semiconductor or injecting electrons or holes from a contact can result in a **non-equilibrium** distribution $np \neq n_i^2$



Recombination

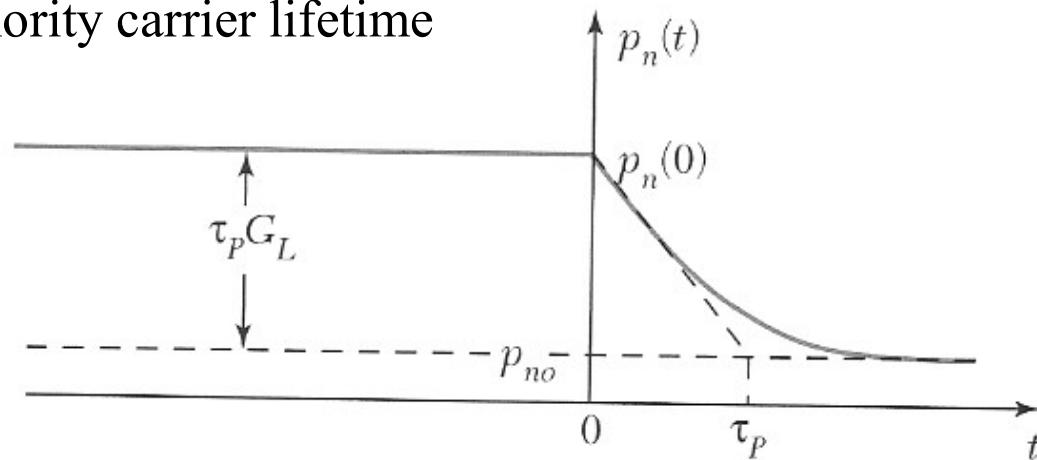


Recombination rate is limit by the density of minority carriers.
The majority carriers have to find a minority carrier to recombine.

p_n (or n_p) = minority carrier concentration

p_{n0} (or n_{p0}) = equilibrium minority carrier concentration

τ_p = minority carrier lifetime



minority carrier lifetimes

p-type

$$n_p(t) = n_{excess} \exp(-t / \tau_n) + n_{p0}$$

n-type

$$p_n(t) = p_{excess} \exp(-t / \tau_p) + p_{n0}$$

minority carrier
lifetimes

$$np = n_i^2$$

Continuity equations

$$\frac{\partial n}{\partial t} = \frac{1}{e} \nabla \cdot \vec{j}_n + G_n - R_n$$

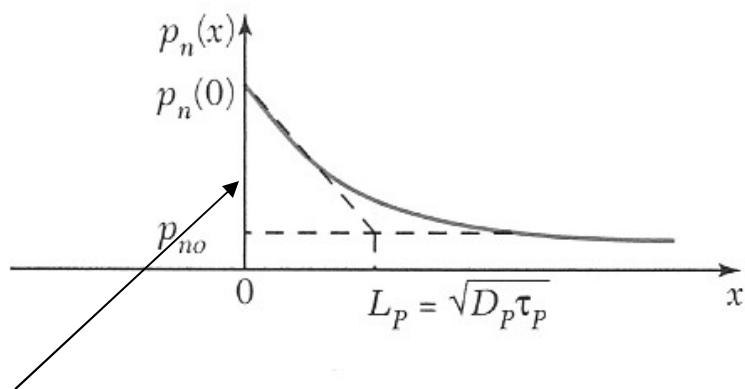
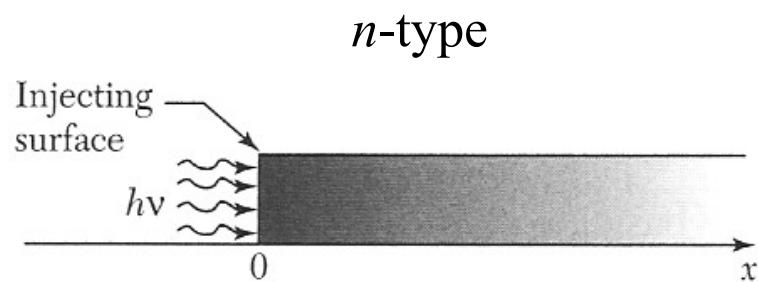
drift: $\vec{j}_n = -ne\mu_n \vec{E}$ $\nabla \cdot \vec{j}_n = -en\mu_n \nabla \cdot \vec{E} - e\nabla n \mu_n \vec{E}$

diffusion: $\vec{j}_{n,diff} = |e| D_n \nabla n$ $\nabla \cdot \vec{j}_{n,diff} = |e| D_n \nabla^2 n$

$$\frac{\partial n}{\partial t} = n\mu_n \nabla \cdot \vec{E} + \nabla n \mu_n \vec{E} + D_n \nabla^2 n + G_n - \frac{n - n_0}{\tau_n}$$

$$\frac{\partial p}{\partial t} = -n\mu_p \nabla \cdot \vec{E} - \nabla n \mu_p \vec{E} + D_p \nabla^2 p + G_p - \frac{p - p_0}{\tau_p}$$

Diffusion Length



Steady state

$$\frac{\partial p_n}{\partial t} = 0 = D_p \frac{\partial^2 p_n}{\partial x^2} - \frac{p_n - p_{n0}}{\tau_p}$$

$$p_n(x) = p_{n0} + (p_n(0) - p_{n0}) \exp\left(\frac{-x}{L_p}\right)$$

Generation only occurs at the surface. There the minority carrier density is $p_n(0)$.

Diffusion Length

$$0 = D_p \frac{\partial^2 p_n}{\partial x^2} - \frac{p_n - p_{n0}}{\tau_p} \quad \Leftrightarrow \quad p_n(x) = p_{n0} + (p_n(0) - p_{n0}) \exp\left(\frac{-x}{L_p}\right)$$

$$0 = \frac{D_p (p_n(0) - p_{n0})}{L_p^2} \exp\left(\frac{-x}{L_p}\right) - \frac{(p_n(0) - p_{n0})}{\tau_p} \exp\left(\frac{-x}{L_p}\right)$$

$$L_p = \sqrt{D_p \tau_p}$$

diffusion length,
typically microns

Haynes Shockley experiment

$$n_p(x,t) = \frac{n_{generated}}{\sqrt{4\pi D_n t}} \exp\left(-\frac{(x - \mu_n E t)^2}{4D_n t}\right) \exp\left(-\frac{t}{\tau_n}\right) + n_{p0}$$

