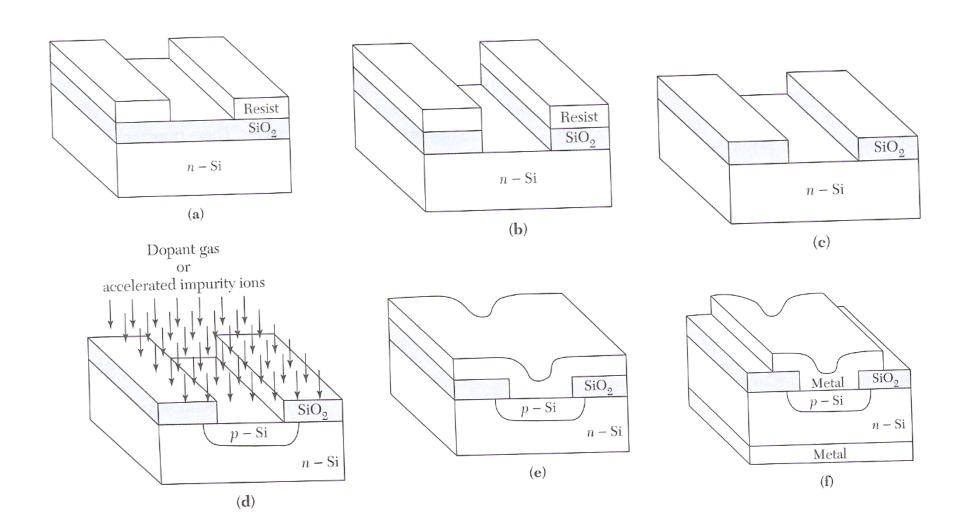


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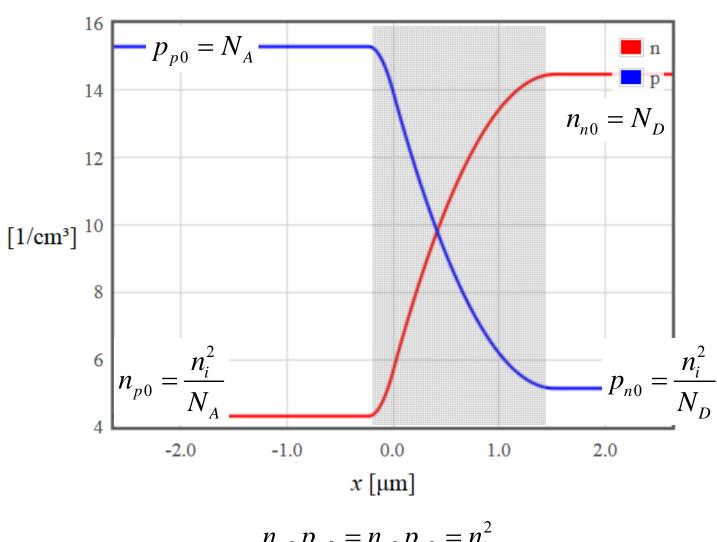
pn-junctions

diode fabrication



Equilibrium concentrations, V = 0

log(Carrier densities)

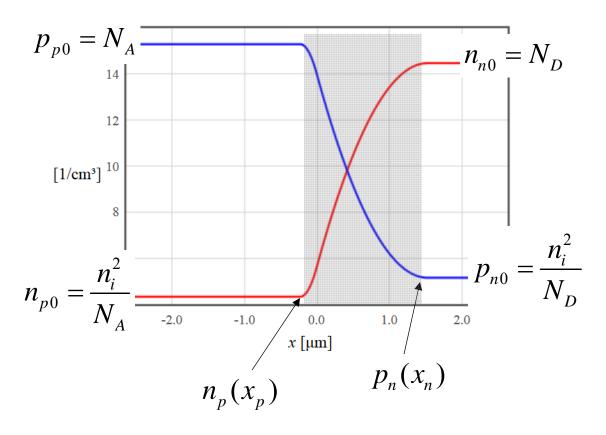


$$n_{p0}p_{p0} = n_{n0}p_{n0} = n_i^2$$

Bias voltage, V = 0

$$eV_{bi} = k_B T \ln \left(\frac{N_D N_A}{n_i^2}\right) = k_B T \ln \left(\frac{N_D}{n_{p0}}\right) = k_B T \ln \left(\frac{N_A}{p_{n0}}\right)$$

log(Carrier densities)



$$n_{p0}p_{p0} = n_{n0}p_{n0} = n_i^2$$

$$V = 0$$

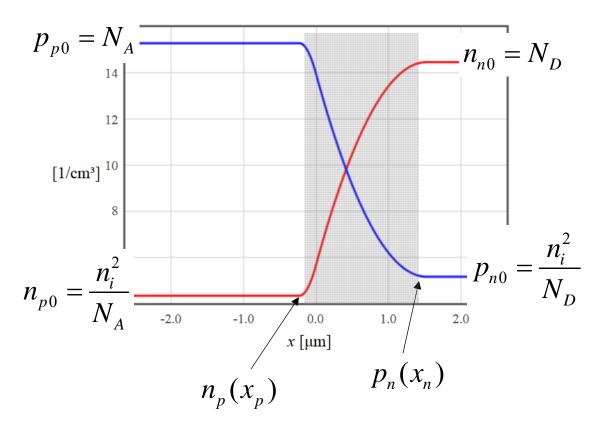
$$n_{p0} = N_D \exp\left(\frac{-eV_{bi}}{k_B T}\right)$$

$$p_{n0} = N_A \exp\left(\frac{-eV_{bi}}{k_B T}\right)$$

Bias voltage, $V \neq 0$

$$eV_{bi} = k_B T \ln \left(\frac{N_D N_A}{n_i^2}\right) = k_B T \ln \left(\frac{N_D}{n_{p0}}\right) = k_B T \ln \left(\frac{N_A}{p_{n0}}\right)$$

log(Carrier densities)



$$n_{p0}p_{p0} = n_{n0}p_{n0} = n_i^2$$

$$V = 0$$

$$n_{p0} = N_D \exp\left(\frac{-eV_{bi}}{k_B T}\right)$$

$$p_{n0} = N_A \exp\left(\frac{-eV_{bi}}{k_B T}\right)$$

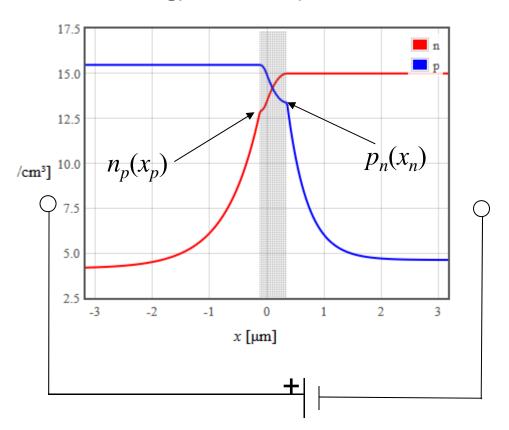
$$p_{n0} = \frac{n_i^2}{N_D}$$

$$N_D = N_D \exp\left(\frac{-e(V_{bi} - V)}{k_B T}\right)$$

$$p_n(x_n) = N_A \exp\left(\frac{-e(V_{bi} - V)}{k_B T}\right)$$

Forward bias, V > 0

log(Carrier Densities)



Electrons and holes are driven towards the junction.

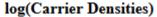
The depletion region becomes narrower

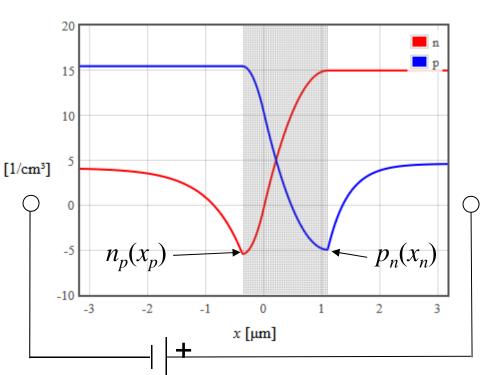
$$n_p(x_p) = N_D \exp\left(\frac{-e(V_{bi} - V)}{k_B T}\right)$$

$$p_n(x_n) = N_A \exp\left(\frac{-e(V_{bi} - V)}{k_B T}\right)$$

Minority electrons are injected into the p-region Minority holes are injected into the n-region

Reverse bias, V < 0





Electrons and holes are driven away from the junction.

The depletion region becomes wider

$$n_p(x_p) = N_D \exp\left(\frac{-e(V_{bi} - V)}{k_B T}\right)$$

$$n_{p}(x_{p}) = N_{D} \exp\left(\frac{-e(V_{bi} - V)}{k_{B}T}\right)$$

$$p_{n}(x_{n}) = N_{A} \exp\left(\frac{-e(V_{bi} - V)}{k_{B}T}\right)$$

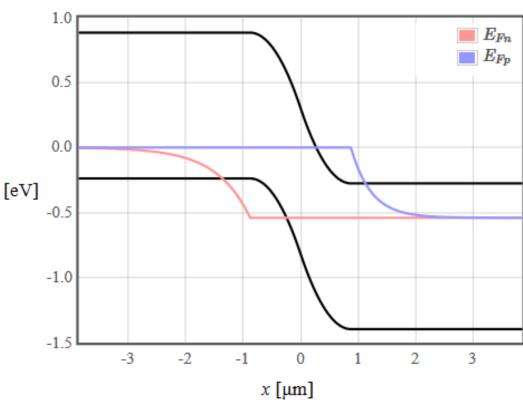
Minority electrons are extracted from the p-region by the electric field Minority holes are extracted from the n-region by the electric field

Quasi Fermi level

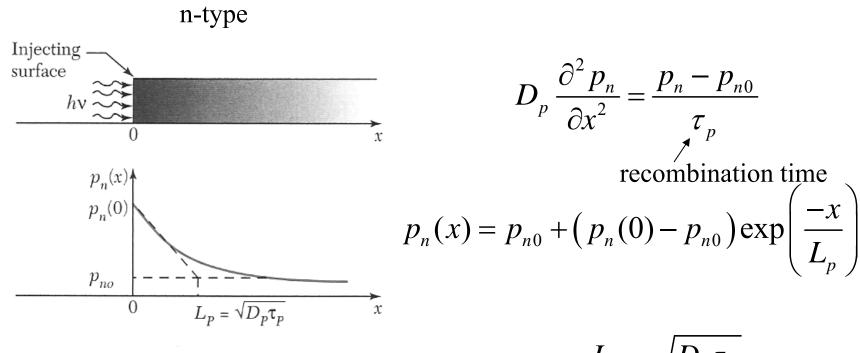
When the charge carriers are not in equilibrium the Fermi energy can be different for electrons and holes.

$$n = N_c \exp\left(\frac{E_{Fn} - E_c}{k_B T}\right)$$

$$p = N_{v} \exp\left(\frac{E_{v} - E_{Fp}}{k_{B}T}\right)$$



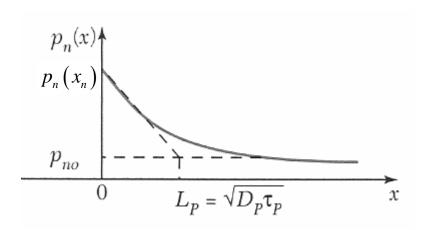
Review of Diffusion



Injection only occurs at the surface. There the minority carrier density is $p_n(0)$.

$$L_p = \sqrt{D_p \tau_p}$$
 diffusion length

n-type



$$p_n(x) = p_{n0} + (p_n(x_n) - p_{n0}) \exp\left(\frac{-x}{L_p}\right)$$

$$J_{diff,p} = -eD_p \frac{dp}{dx}$$

$$J_{diff,p} = -eD_p \frac{dp}{dx} = \left(p_n(x_n) - p_{n0}\right) \frac{eD_p}{L_p} \exp\left(\frac{-x}{L_p}\right)$$

At the edge of the depletion region:

$$J_{diff,p} = -eD_p \frac{dp}{dx} = (p_n(x_n) - p_{n0}) \frac{eD_p}{L_p}$$

$$J_{diff,p} = \left(p_n(x_n) - p_{n0}\right) \frac{eD_p}{L_p}$$

$$p_n(x_n) = p_{p0} \exp\left(-\frac{e(V_{bi} - V)}{k_B T}\right)$$

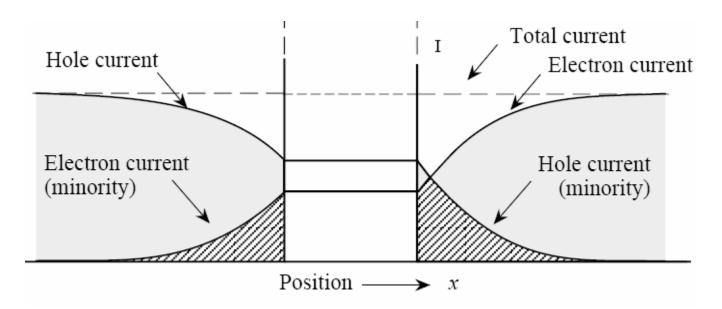
$$J_{diff,p} = \left(p_{p0} \exp\left(-\frac{e(V_{bi} - V)}{k_B T}\right) - p_{n0}\right) \frac{eD_p}{L_p}$$

$$p_{p0} = p_{n0} \exp\left(\frac{eV_{bi}}{k_B T}\right)$$

$$J_{diff,p} = p_{n0} \frac{eD_p}{L_p} \left(\exp\left(\frac{eV}{k_B T}\right) - 1\right)$$

$$J_{diff,p} = \frac{p_{n0}eD_p}{L_p} \left(\exp\left(\frac{eV}{k_BT}\right) - 1 \right)$$

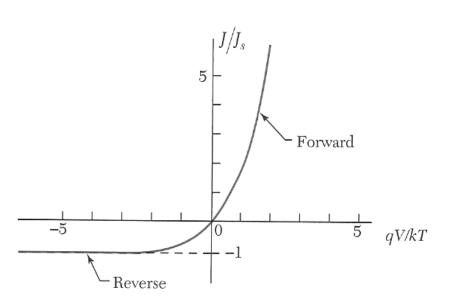
$$J_{diff,n} = \frac{n_{p0}eD_n}{L_n} \left(\exp\left(\frac{eV}{k_BT}\right) - 1 \right)$$



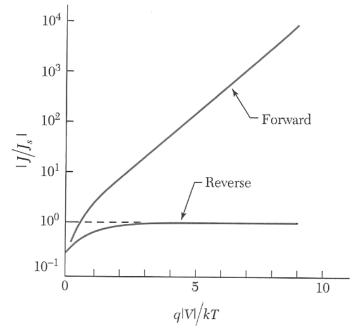
Diode current

$$I = eA \left(\frac{p_{n0}D_p}{L_p} + \frac{n_{p0}D_n}{L_n} \right) \left(\exp\left(\frac{eV}{k_BT}\right) - 1 \right) = I_s \left(\exp\left(\frac{eV}{k_BT}\right) - 1 \right)$$

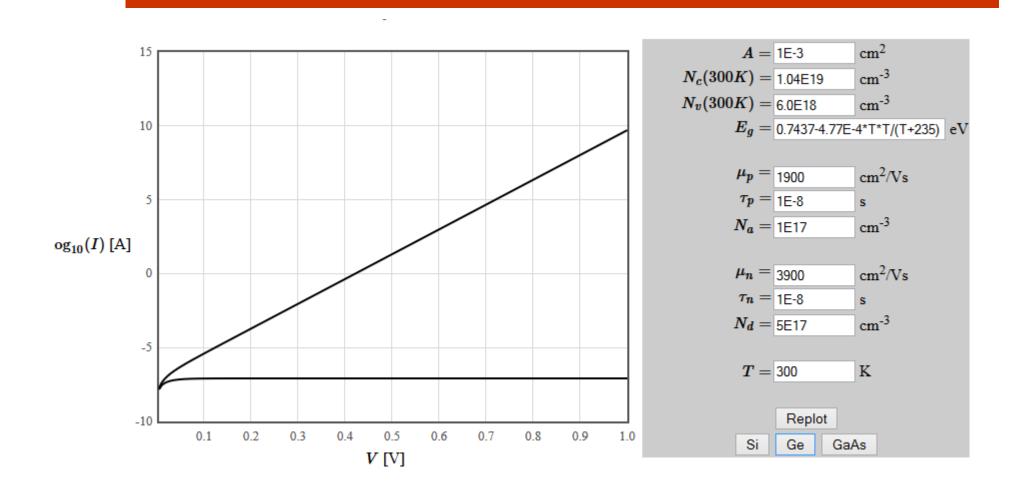
Area



Saturation current



Diode I-V charateristics



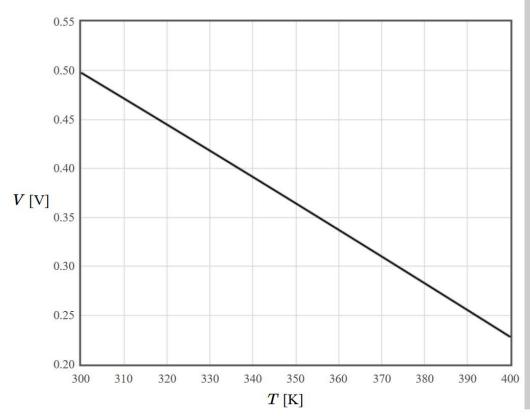
http://lamp.tu-graz.ac.at/~hadley/psd/L6/pnIV.php

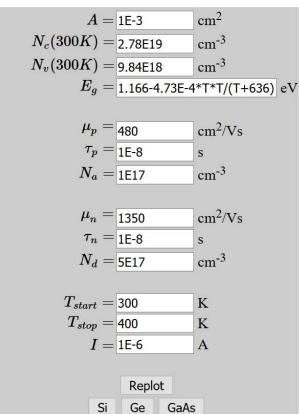
Thermometer

$$I_S = Aen_i^2igg(rac{D_p}{L_pN_d} + rac{D_n}{L_nN_a}igg)$$

$$n_i = \sqrt{N_c igg(rac{T}{300}igg)^{3/2} N_v igg(rac{T}{300}igg)^{3/2}} \expigg(rac{-E_g}{2k_B T}igg)$$

$$D_n = rac{\mu_n k_B T}{e}$$

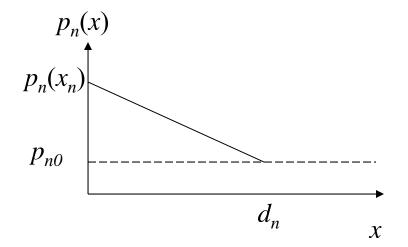




$n_{p}(x_{p}) = N_{D} \exp\left(\frac{-e(V_{bi} - V)}{k_{B}T}\right)$ $p_{n}(x_{n}) = N_{A} \exp\left(\frac{-e(V_{bi} - V)}{k_{B}T}\right)$ Forward Reverse npn $q(V_{bi} + V)$ $J_{\it diff} > J_{\it drift}$ $J_{diff} < J_{drift}$ p_{po} p_{po} 10¹⁸ L 9.65×10^{9} $p_{\underline{no}}$ $9.3 \times 10^{1} n_{i}$ $|p_{no}(e^{q\dot{V}/kT}-1)$ x_n Injection Extraction

Short diode





$$d_n \ll L_p$$

Metal contact is much closer to the depletion region than the diffusion length

$$J_{diff,p} = -eD_p \frac{dp}{dx}$$

$$J_{diff,p} = -eD_p \frac{dp}{dx} = \frac{eD_p}{d_n} (p_n(x_n) - p_{n0})$$

$$J_{diff,p} = \left(p_n(x_n) - p_{n0}\right) \frac{eD_p}{d_n}$$

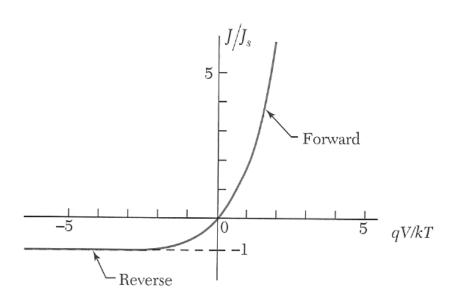
$$J_{diff,p} = \left(p_{n0} \exp\left(\frac{e(V)}{k_B T}\right) - p_{n0}\right) \frac{eD_p}{d_n}$$

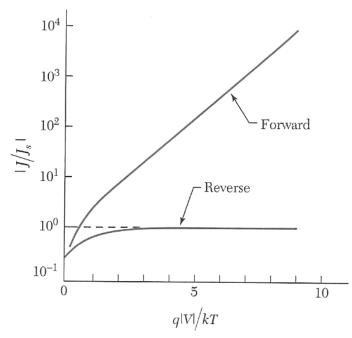
$$J_{diff,p} = \frac{p_{n0}eD_p}{d_n} \left(\exp\left(\frac{eV}{k_B T}\right) - 1 \right)$$

Short diode current

$$I = eA \left(\frac{p_{n0}D_p}{d_n} + \frac{n_{p0}D_n}{d_p} \right) \left(\exp\left(\frac{eV}{k_BT}\right) - 1 \right) = I_s \left(\exp\left(\frac{eV}{k_BT}\right) - 1 \right)$$

Area



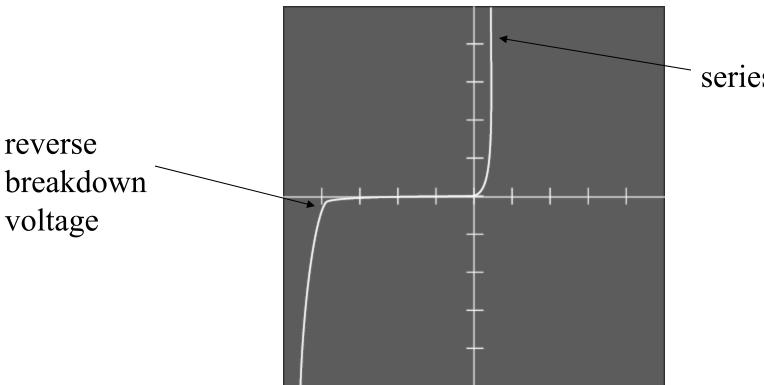


Real diodes

$$I = I_s \left(\exp\left(\frac{eV}{nk_BT}\right) - 1 \right)$$

n = nonideality factor

n = 1 for an ideal diode



series resistance

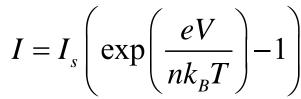
Real diodes

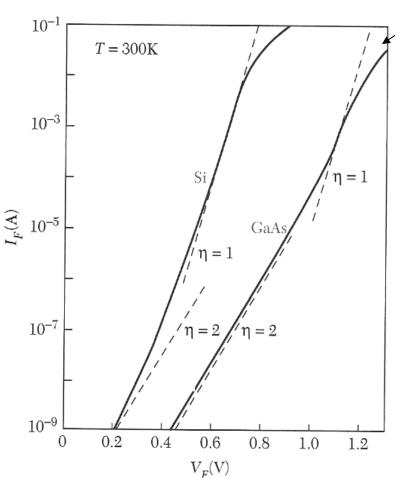
There is constant generation/recombination of electron hole pairs.

In forward bias there is an extra current from recombination.

In reverse bias there is an extra current from generation.

Low bias: recombination dominates, n = 2

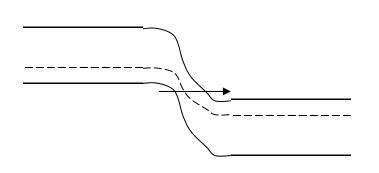


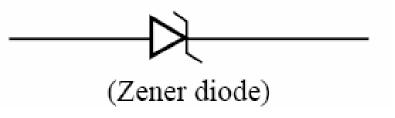


Very high bias: series resistance

High bias: ideal behavior, n = 1

Zener tunneling

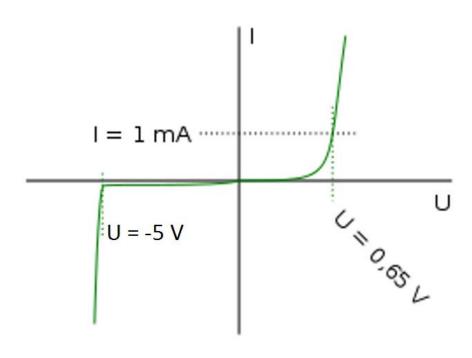




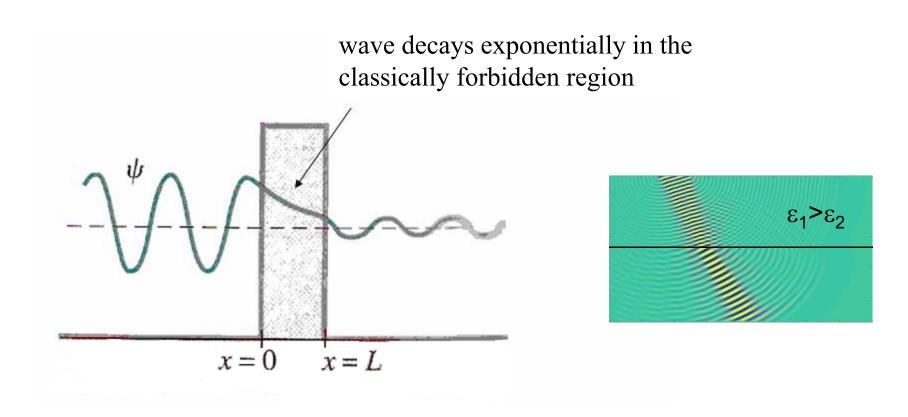
Electrons tunnel from valence band to conduction band

Occurs at high doping

$$|V_{\text{zener}}| < 5.6 \text{ V}$$



Tunneling



Tunneling is a wave phenomena. Tunneling and total internal reflection are used in a beam splitter.

Zener tunneling

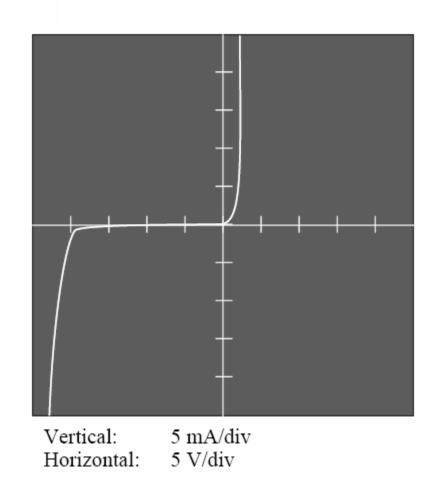
Breakdown voltage is typically much lower than the breakdown voltage of an avalanche diode and can be tuned by adjusting the width of the depletion layer.

Used to provide a reference voltage.

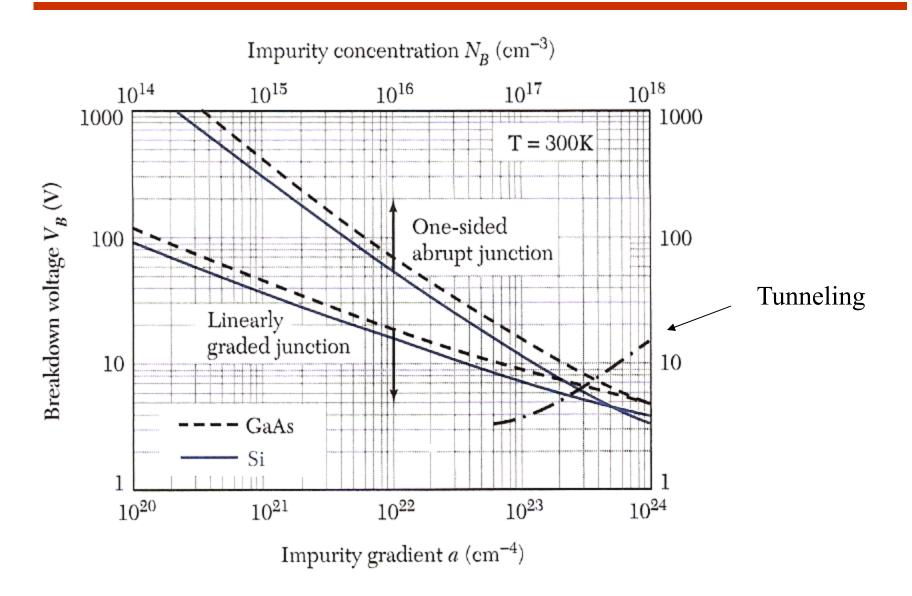
Avalanche breakdown

Impact ionization causes an avalanche of current

Occurs at low doping



Avalanche breakdown





Technische Universität Graz

metal - semiconductor contacts

Photoelectric effect
Workfunction
Electron affinity
Interface states
Schottky barriers
Schottky diodes
Ohmic contacts
Thermionic emission
Tunnel contacts



