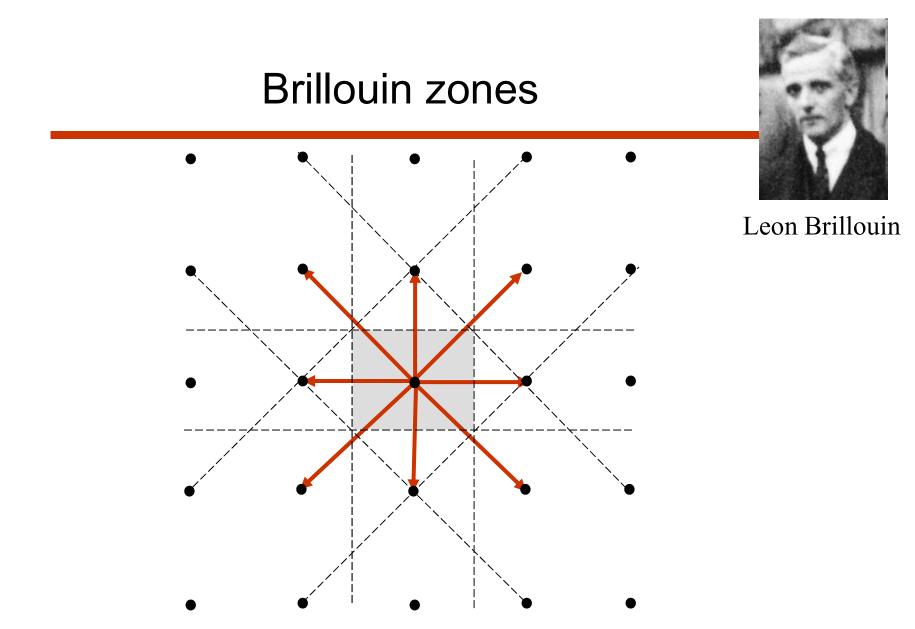


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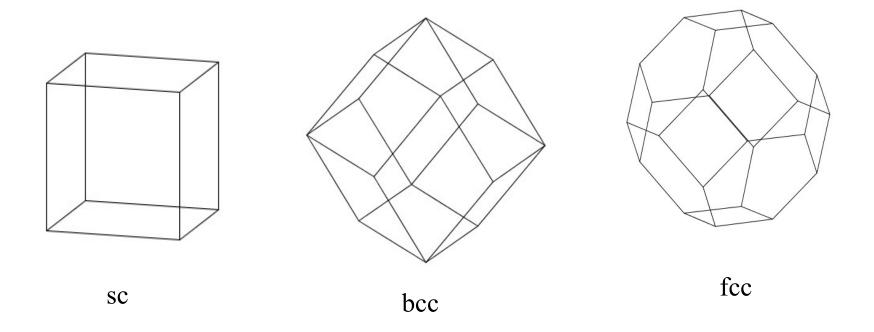
Institute of Solid State Physics

Diffraction



1st Brillouin zone consists of the *k*-states around the origin that can be reached without crossing a plane.

1st Brillouin zones



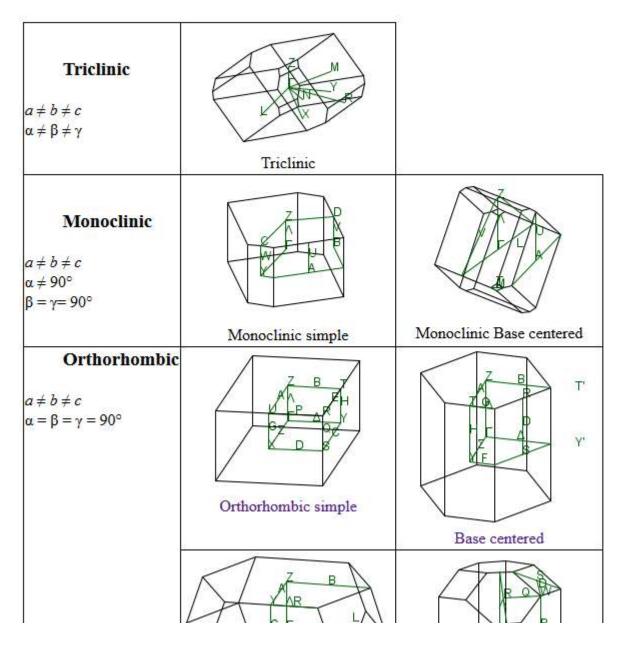
1st Brillouin is the Wigner-Seitz cell in reciprocal space.



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Brillouin zones



Electron density of an atom

Most of the electrons are concentrated around the nucleus. The integral over the electron density is proportional to the number of electrons.

$$n_j(\vec{r}) \propto \exp\left(-\frac{\left(\vec{r}-\vec{r}_j\right)^2}{r_0^2}\right)$$

Approximately a Gaussian centered at r_i

Electron density

Write the electron density as a Fourier series

$$n(\vec{r}) = \sum_{\vec{G}} n_{\vec{G}} e^{i\vec{G}\cdot\vec{r}} = \sum_{\vec{T}} \sum_{j} n_{j} \left(\vec{r} - \vec{r}_{j} + \vec{T}\right),$$

Translation of atom *i* of the basis

position of atom j of the basis

on vector

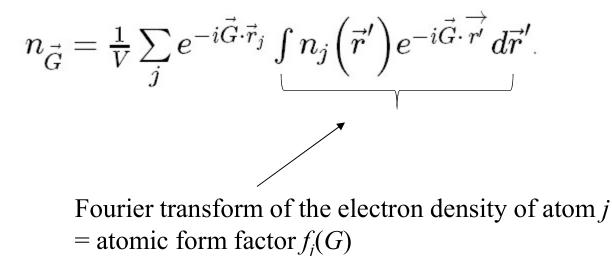
Multiply by $e^{-i\vec{G}'\cdot\vec{r}}$ and integrate over a unit cell.

$$\sum_{\vec{G}} \int_{\mathbf{u.c.}} n_{\vec{G}} e^{i\vec{G}\cdot\vec{r}} e^{-i\vec{G}'\cdot\vec{r}} d\vec{r} = \sum_{j} \int_{\mathbf{u.c.}} n_j (\vec{r}-\vec{r}_j) e^{-i\vec{G}'\cdot\vec{r}} d\vec{r}.$$

Electron density

$$n_{\vec{G}}V = \sum_{j} \int n_{j} \left(\vec{r} - \vec{r}_{j}\right) e^{-i\vec{G}\cdot\vec{r}} d\vec{r}$$

Make a substitution $\vec{r}' = \vec{r} - \vec{r}_j$.

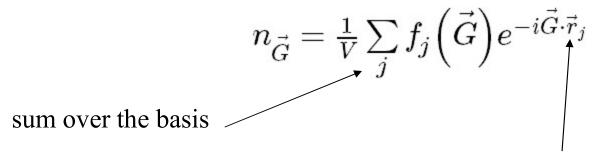


Atomic form factor

$$f_j\left(\vec{G}\right) = \int n_j\left(\vec{r}\right) e^{-i\vec{G}\cdot\vec{r}} d\vec{r},$$

The atomic form factors can be looked up in a table.

The structure factors are given in terms of the atomic form factors.



position of atom *j* of the basis

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Element	Z		a ₂	a3	a4	a5	<i>b</i> 1	b2	<i>b</i> ₃	<i>b</i> ₄	<u>b5</u>
H	1	0.0349	0.1201	0.1970	0.0573	0.1195	0.5347	3.5867	12.3471	18.9525	38.6269
He	2	0.0317	0.0838	0.1526	0.1334	0.0164	0.2507	1.4751	4.4938	12.6646	31.1653
Li	3	0.0750	0.2249	0.5548	1.4954	0.9354	0.3864	2.9383	15.3829	53.5545	138.7337
Be	4	0.0780	0.2210	0.6740	1.3867	0.6925	0.3131	2.2381	10.1517	30.9061	78.3273
В	5	0.0909	0.2551	0.7738	1.2136	0.4606	0.2995	2.1155	8.3816	24.1292	63.1314
С	6	0.0893	0.2563	0.7570	1.0487	0.3575	0.2465	1.7100	6.4094	18.6113	50.2523
N	7	0.1022	0.3219	0.7982	0.8197	0.1715	0.2451	1.7481	6.1925	17.3894	48.1431
0	8	0.0974	0.2921	0.6910	0.6990	0.2039	0.2067	1.3815	4.6943	12.7105	32.4726
F	9	0.1083	0.3175	0.6487	0.5846	0.1421	0.2057	1.3439	4.2788	11.3932	28.7881
Ne	10	0.1269	0.3535	0.5582	0.4674	0.1460	0.2200	1.3779	4.0203	9.4934	23.1278
		1									
Na	11	0.2142	0.6853	0.7692	1.6589	1.4482	0.3334	2.3446	10.0830	48.3037	138.2700
Mg	12	0.2314	0.6866	0.9677	2.1882	1.1339	0.3278	2.2720	10.9241	39.2898	101.9748
Al	13	0.2390	0.6573	1.2011	2.5586	1.2312	0.3138	2.1063	10.4163	34.4552	98.5344

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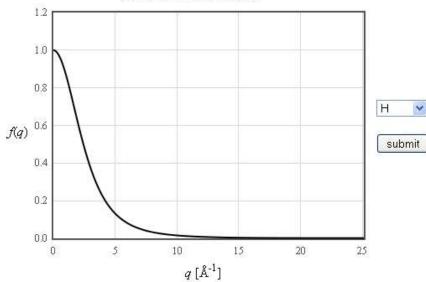
513.001 Molecular and Solid State Physics

Atomic form factors

In the range of scattering vectors between $0 \le q \le 25 \text{ Å}^{-1}$, the atomic form factor is well approximated by the expression, [1]

$$f(q) = \sum_{i=1}^{4} a_i \exp\left(-b_i \left(\frac{q}{4\pi}\right)^2\right) + c_i$$

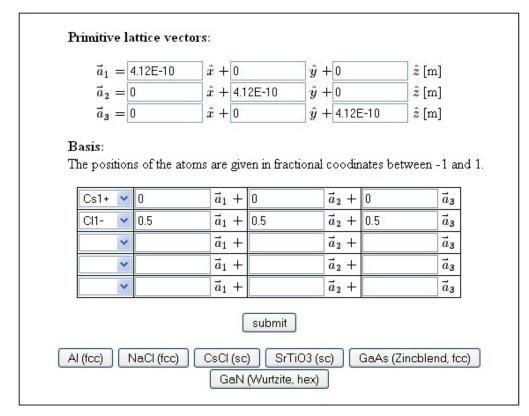
where the values of a_i , b_i , and c are tabulated below. The different atomic form factors for the elements can be plotted using the form below.



*b*₁ b2 b4 Element b3 a a_2 a_3 a4 с H 0.489918 20.6593 0.262003 7.74039 0.196767 49.5519 2.20159 0.049879 0.001305 53.1368 15.187 0.415815 186.576 3.56709 H1-0.897661 0.565616 0.116973 0.002389 He 0.8734 9.1037 0.6309 3.3568 22.9276 0.178 0.9821 0.3112 0.0064 ·τ. 1 1000 A 3640 1.0004 A 6196 00.0000 A 4000 120.021 A A000

Home Outline Introduction Molecules Crystal Structure **Crystal Diffraction** Crystal Binding Photons Phonons Electrons Energy bands Crystal Physics Semiconductors Magnetism Exam questions Appendices Lectures **TUG** students Student projects Skriptum Books Making presentations < hide <

Atomic form factor for H



Primitive reciprocal lattice vectors

$$\begin{split} \vec{b}_1 &= 2\pi \, \frac{\vec{a}_2 \times \vec{a}_3}{\vec{a}_1 \cdot (\vec{a}_2 \times \vec{a}_3)} \, = & 1.525 \text{e}{\pm}10 \, \hat{k}_x \pm 0.000 \, \hat{k}_y \pm 0.000 \, \hat{k}_z \, [\text{m}^{-1}] \\ \vec{b}_2 &= 2\pi \, \frac{\vec{a}_3 \times \vec{a}_1}{\vec{a}_1 \cdot (\vec{a}_2 \times \vec{a}_3)} \, = & 0.000 \, \hat{k}_x \pm 1.525 \text{e}{\pm}10 \, \hat{k}_y \pm 0.000 \, \hat{k}_z \, [\text{m}^{-1}] \\ \vec{b}_3 &= 2\pi \, \frac{\vec{a}_1 \times \vec{a}_2}{\vec{a}_1 \cdot (\vec{a}_2 \times \vec{a}_3)} \, = & 0.000 \, \hat{k}_x \pm 0.000 \, \hat{k}_y \pm 1.525 \text{e}{\pm}10 \, \hat{k}_z \, [\text{m}^{-1}] \end{split}$$

Structure factors

The value of $|n_{\vec{G}}|$ for the 000 diffraction peak is the total number of electrons in the primitive unit cell. The intensities of the peaks in an x-ray diffraction experiment

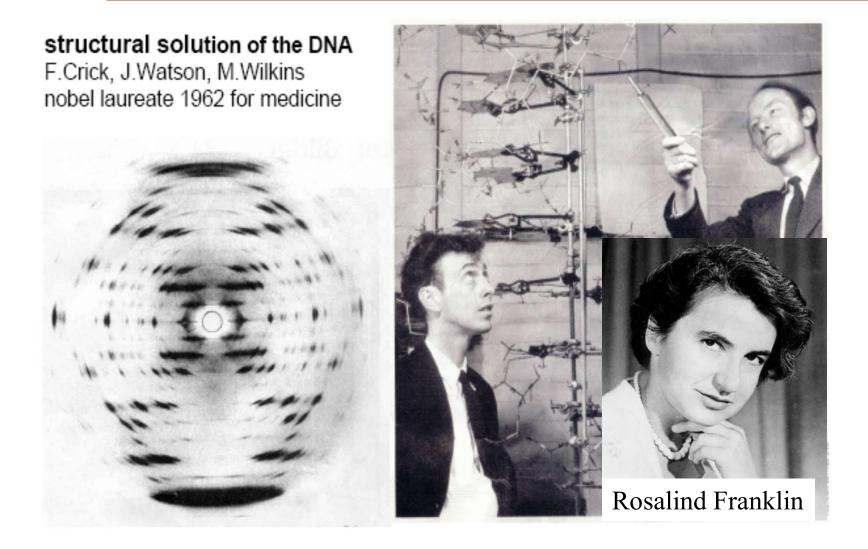
hkl	$ ec{G} $ Å ⁻¹	$n_{\vec{G}}$	$ n_{\tilde{G}} ^2$	$\operatorname{Re}\{n_{\tilde{G}}\}$	$\operatorname{Im}\{n_{\tilde{G}}\}$
000	0.000	72.00	5184	72.00	0.000
-100	1.525	34.43	1185	34.43	5.333e-8
0-10	1.525	34.43	1185	34.43	5.333e-8
00-1	1.525	34.43	1185	34.43	5.333e-8
001	1.525	34.43	1185	34.43	-5.333e-8
	1 000		1100		

Structure factor

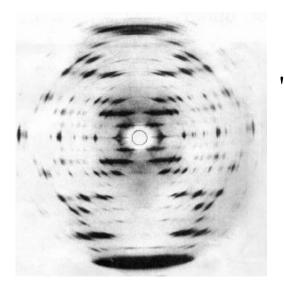
$$n_{ec{G}} = rac{1}{V} \sum_j f_j(G) e^{-iec{G}\cdotec{r}_j}$$

An x-ray experiment measures the scattered intensity $|F_G|^2$. The phase information is lost. This is proportional to $|n_G|^2$.

crystal structure solution



crystal structure solution



"Guess" the crystal structure

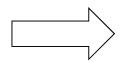


Table 4.3.2.2 |pdf |

Elastic atomic scattering factors of electrons for neutral atoms and s up to 2.0 $\mbox{\AA}^{-1}$

Element	Z	a_1	a2	a3	a4	a5	b_1	b_2	b_3
Н	1	0.0349	0.1201	0.1970	0.0573	0.1195	0.5347	3.5867	12.3471
He	2	0.0317	0.0838	0.1526	0.1334	0.0164	0.2507	1.4751	4.4938
Li	3	0.0750	0.2249	0.5548	1.4954	0.9354	0.3864	2.9383	15.3829
Be	4	0.0780	0.2210	0.6740	1.3867	0.6925	0.3131	2.2381	10.1517
В	5	0.0909	0.2551	0.7738	1.2136	0.4606	0.2995	2.1155	8.3816
с	6	0.0893	0.2563	0.7570	1.0487	0.3575	0.2465	1.7100	6.4094
N	7	0.1022	0.3219	0.7982	0.8197	0.1715	0.2451	1.7481	6.1925
0	8	0.0974	0.2921	0.6910	0.6990	0.2039	0.2067	1.3815	4.6943
F	9	0.1083	0.3175	0.6487	0.5846	0.1421	0.2057	1.3439	4.2788
Ne	10	0.1269	0.3535	0.5582	0.4674	0.1460	0.2200	1.3779	4.0203
Na	11	0.2142	0.6853	0.7692	1.6589	1.4482	0.3334	2.3446	10.0830
Mg	12	0.2314	0.6866	0.9677	2.1882	1.1339	0.3278	2.2720	10.9241
Al	13	0.2390	0.6573	1.2011	2.5586	1.2312	0.3138	2.1063	10.4163
Si	14	0.2519	0.6372	1.3795	2.5082	1.0500	0.3075	2.0174	9.6746
Р	15	0.2548	0.6106	1.4541	2.3204	0.8477	0.2908	1.8740	8.5176
S	16	0.2497	0.5628	1.3899	2.1865	0.7715	0.2681	1.6711	7.0267
Cl	17	0.2443	0.5397	1.3919	2.0197	0.6621	0.2468	1.5242	6.1537
Ar	18	0.2385	0.5017	1.3428	1.8899	0.6079	0.2289	1.3694	5.2561
	1.0	laure		0.0004	0.0240	0.01.00	0.0700	0.0074	10.1000

Compare $|n_G|^2$ to the measurements

From the atomic form factors, calculate the structure factors n_G .

$$n_{ec{G}} = \sum\limits_{j} f_j(G) e^{-iec{G}\cdotec{r}_j}$$

Sum over basis

position of atom *j* of the basis

x-ray diffraction

The shape and the dimensions of the unit cell can be deduced from the positions of the Bragg reflections; the content of the unit cell, on the other hand, must be determined from the intensities of the reflections.

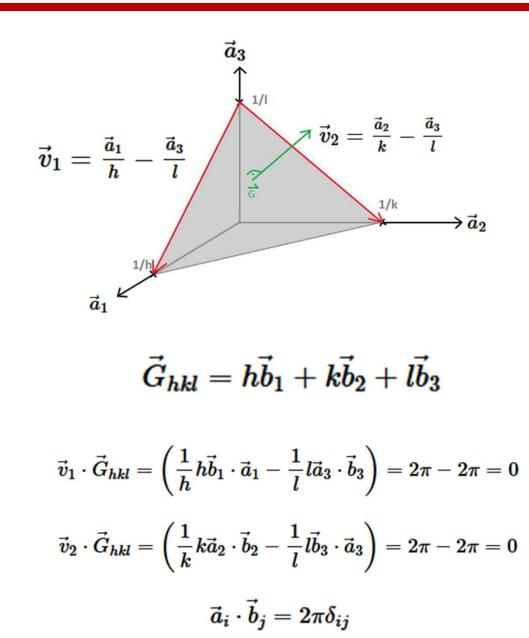
Solid State Physics, Ibach and Lüth

Diffraction condition:
$$\Delta \vec{k} = \vec{G}$$

The intensity of the peaks is proportional to the squared Fourier coefficients of the electron density.

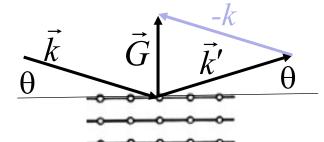
$$n_{ec{G}} = \sum_j f_j(G) e^{-iec{G}\cdotec{r}_j}$$

The reciprocal lattice vector \vec{G}_{hkl} is orthogonal to the (hkl) plane

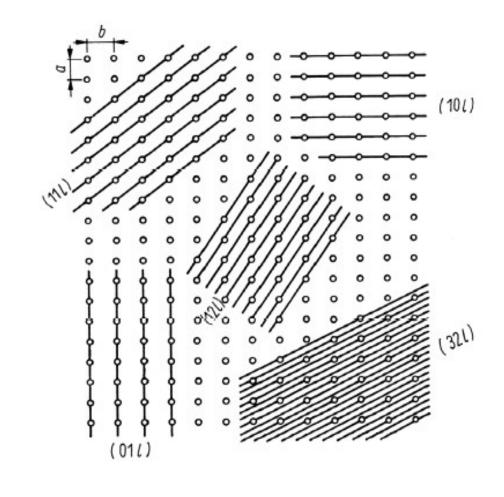


x-ray diffraction

 $\vec{G}_{hkl} \perp (hkl)$



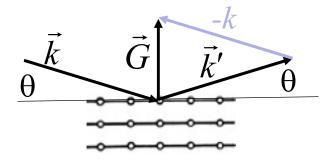
 $\left|\vec{G}_{hkl}\right| = \frac{2\pi}{d_{hkl}}$



distance between the net planes

x-ray diffraction

$$\vec{G}_{hkl} \perp (hkl)$$



$$\left|\vec{G}_{hkl}\right| = \frac{2\pi}{d_{hkl}}$$
$$\left|\vec{k}\right| = \frac{2\pi}{\lambda}$$

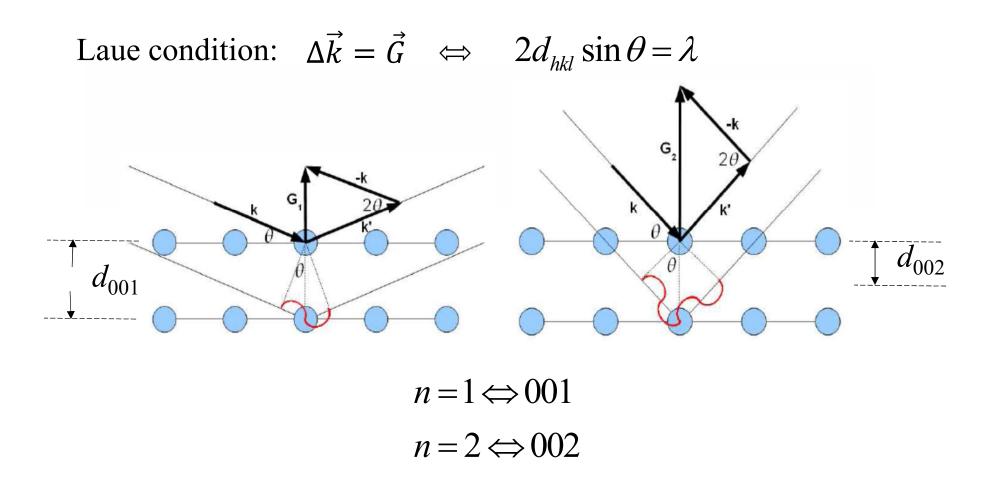
$$\frac{2\pi}{d_{hkl}} = 2\frac{2\pi}{\lambda}\sin\theta$$

 $\left|\vec{G}_{hkl}\right| = \left|\Delta \vec{k}\right| = 2\left|\vec{k}\right| \sin \theta$

 $2d_{hkl}\sin\theta = \lambda$ another formulation of the diffraction condition

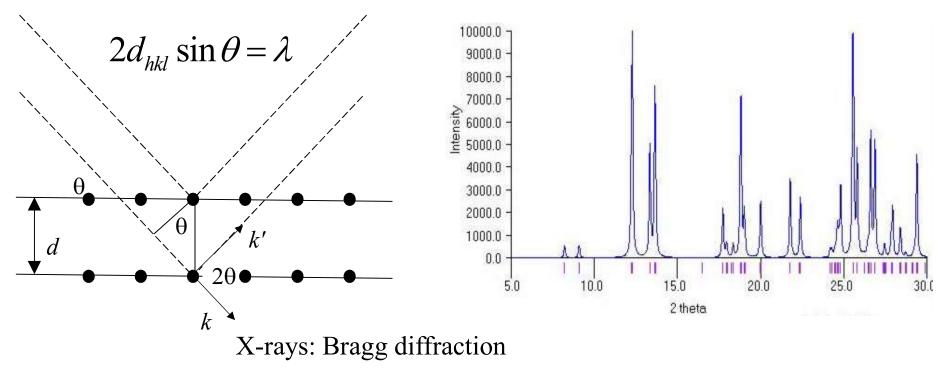
Bragg and Laue conditions

Bragg condition: $2d\sin\theta = n\lambda$



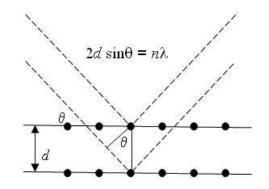
Powder diffraction

Powder diffraction is performed on a powder of many small crystals. Ideally, every possible crystalline orientation is represented equally in a powdered sample. The relative intensities of the diffraction peaks indicate which crystal structures are present.

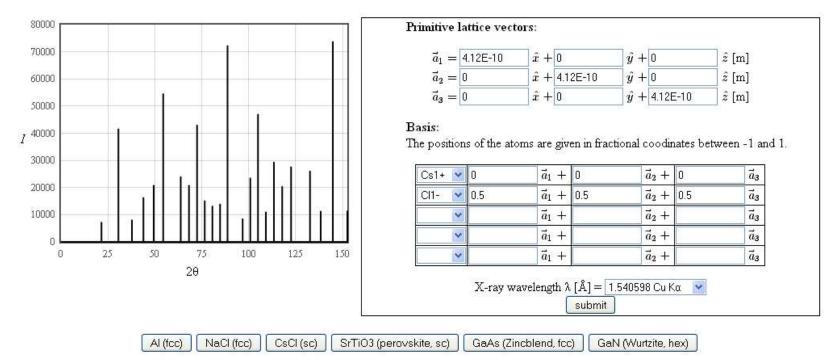


Powder diffraction

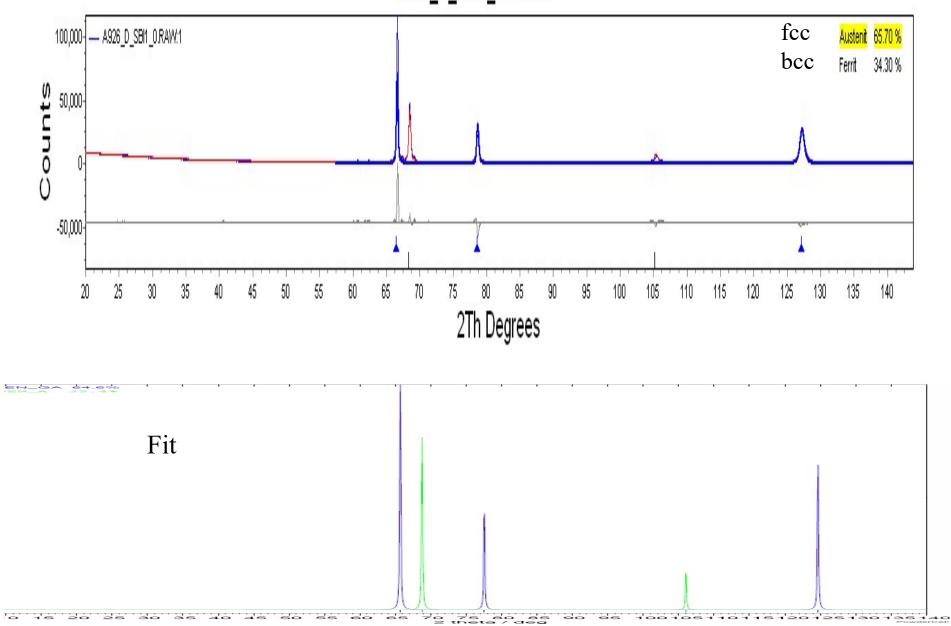
For powder diffraction, a crystal is ground into a fine powder so that there are many small crystals with random orientations. X-rays strike the surface of the sample at an angle θ and an x-ray detector is placed at an angle θ to the surface. Only planes parallel to the surface will diffract x-rays to the detector.



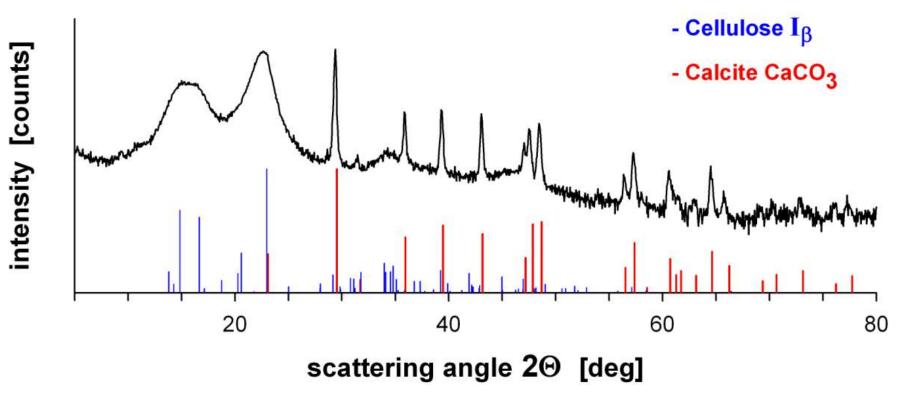
Since there are many small crystals with random orientations in the sample, all possible crystal planes that can diffract the x-rays will contribute to the measured signal when θ satisifies the Bragg condition. The form below can calculate the powder diffraction pattern for any crystal with up to five atoms per primitive unit cell. Some buttons are provided that load the form with the data for certain crystals.



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http://rruff.geo.arizona.edu/AMS/all_minerals.php

American Mineralogist Crystal Structure Database

Abellaite	Abelsonite	Abenakiite-(Ce)	Abernathyite	Abhurite						
Abswurmbachite	Acanthite	Acetamide	Acetylene-hydrate	Achavalite						
Actinium	Actinolite	Acuminite	Adachiite	Adamantane						
Adamantane-methane- hydrate	<u>Adamite</u>	Adamsite-(Y)	Adelite	Admontite						
Adolfpateraite	Adranosite	Adranosite-(Fe)	Aegirine	Aenigmatite						
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Algodonite	Alinite Crysta.	l Structures 1 (196	53) 7-83							
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		•	4	4.76	47.49	2.0248	2	0	0	6
			6	5.16	28.01	1.4317	2	2	0	12
			(i)	8.30	30.71	1.2210	3	1	1	24
			8	2.52	8.74	1.1690	2	2	2	8
				-	Bob Downs, Ran et al. (1993)	-				

Powder diffraction

Phase identification

Every crystal has a specific "fingerprint" given by the positions and intensities of the diffraction peaks. The composition of a multi-phase specimen can be determined by fitting its diffraction pattern to the diffraction patterns of pure crystals which can be looked up in a database.

International Centre for Diffraction Data www.icdd.com

Release 2024: 1186076 material data sets

Phase transitions, thermal expansion, piezoelectricity, piezomagnetism, bulk modulus, compliance tensor can be measured.

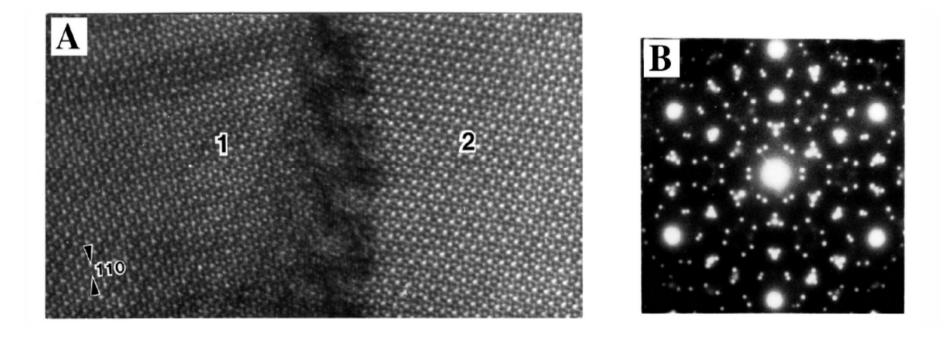
Particle beams

Particles moving in vacuum have the following energymomentum relation.

$$E = \frac{1}{2}mv^{2} = \frac{p^{2}}{2m} = \frac{\hbar^{2}k^{2}}{2m} = \frac{\hbar^{2}}{2m\lambda^{2}}$$

If λ is much smaller than the distance between atoms, you can generate a diffraction pattern.

Electron diffraction in a TEM



The wavelength of the electrons is typically much smaller than the lattice spacing. The diffraction peaks in the plane perpendicular to k are observed.

Electron diffraction

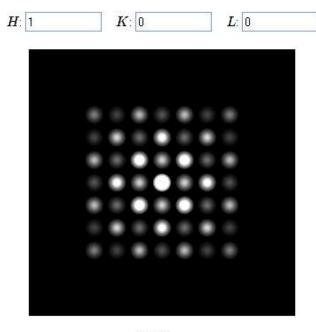
In electron diffraction, the intensity of a diffraction peak at reciprocal lattice vector \vec{G} is the square of the structure factor, $n_{\vec{\sigma}}$.

$$n_{\vec{G}} = \frac{1}{V} \sum_{j} f_j \Big(\vec{G}\Big) e^{-i\vec{G}\cdot\vec{r}_j} = \frac{1}{V} \sum_{j} f_j \Big(\vec{G}\Big) \Big(\cos\Bigl(\vec{G}\cdot\vec{r}_j\Bigr) - i \sin\Bigl(\vec{G}\cdot\vec{r}_j\Bigr) \Big)$$

Here V is the volume of the primitive unit cell, j sums over the atoms in the basis, \vec{r}_j are the positions of the atoms in the basis, and $f_j(\vec{G})$ are the electron atomic form factors evaluated at \vec{G} .

The form below calculates the electron structure factors based on this formula. The crystal structure is specified by providing the primitive lattice vectors and the positions of the atoms in the basis. A basis of up to five atoms can be calculated. The script first calculates the primitive reciprocal lattice vectors and from them calculates the reciprocal lattice vectors $\vec{G}_{hkl} = h\vec{b}_1 + k\vec{b}_2 + l\vec{b}_3$.

On this page, the direction of the incoming electrons is given in terms of the primitive lattice vectors in reciprocal space, $H\vec{b}_1 + K\vec{b}_2 + L\vec{b}_3$. Usually the direction of the incoming electrons are given in terms of the conventional lattice vectors. Be aware that the [100] is a (usually) different direction if primitive lattice vectors are used than if conventional lattice vectors are used.



Primitive lattice vectors:

$\vec{a}_1 =$	4.12E-10	$\hat{x} +$	0	$\hat{y} +$	0	\hat{z} [m]
$\vec{a}_2 =$	0	$\hat{x} +$	4.12E-10	$\hat{y} +$	0	\hat{z} [m]
$\vec{a}_3 =$	0	\hat{x} +	Ó	$\hat{y} +$	4.12E-10	\hat{z} [m]

Basis:

The positions of the atoms are given in fractional coodinates between -1 and 1.

Cs 🔽	0	$\vec{a}_1 + $	0	$\vec{a}_2 + $	0	\vec{a}_3
СІ 🔽	0.5	$\vec{a}_1 + $	0.5	$\vec{a}_2 + $	0.5	\vec{a}_3
*	-	$\vec{a}_1 + $		$\vec{a}_2 +$		\vec{a}_3
*		$\vec{a}_1 + $		$\vec{a}_2 +$		\vec{a}_3
~		$\vec{a}_1 + $		$\vec{a}_2 +$		\vec{a}_3

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Neutron diffraction

Typically a nuclear reactor is used as the neutron source

There are different atomic form factors for neutrons than for x-rays.

Determine the positions of H in biological samples.

Can for example distinguish between Fe and Co which have similar atomic form factors for x-rays.

Structure factor for neutrons

The structure factor for neutrons can be calculated with the following formula,

$$F_{ec{G}} = \sum_{j} b_{j} e^{-iec{G}\cdotec{r}_{j}} = \sum_{j} b_{j} \left(\cos \Bigl(ec{G}\cdotec{r}_{j}\Bigr) - i \sin \Bigl(ec{G}\cdotec{r}_{j}\Bigr)
ight).$$

where \vec{r}_j defines the position of the atom j and \vec{G} is the reciprocal lattice vector. \vec{b}_j is called the neutron scattering length, it depends on the spin-state of the neutron-nucleus system and the isotope the neutron is scattered from. The scattering lengths can be looked up at the <u>NIST Center for Neutron Research</u>.

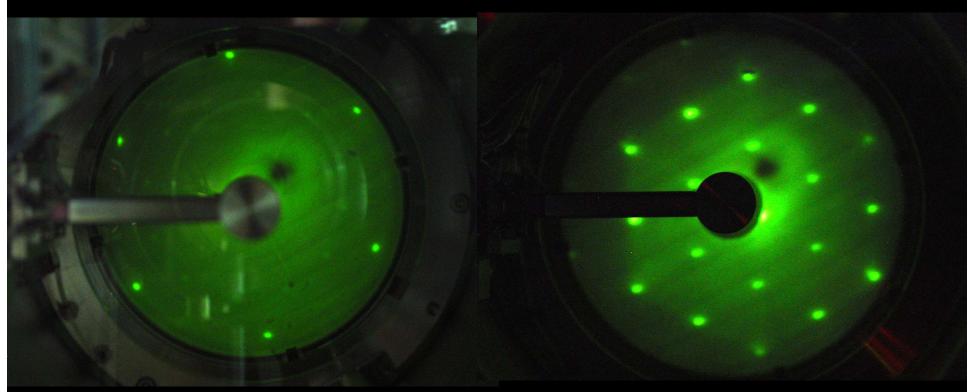
The form below calculates the neutron structure factors. The script first calculates the reciprocal lattice vectors and from them calculates the reciprocal lattice vectors $\vec{G}_{hkl} = h\vec{b}_1 + k\vec{b}_2 + l\vec{b}_3$. The structure factors are calculated for a few reciprocal lattice vectors and listed in a table.

	= 4.1	2E-10	\hat{x} +	0		\hat{y} +	0		<i>î</i> [m]
\vec{a}_2	= 0		\hat{x} +	4.128	-10	\hat{y} +	0		<i>î</i> [m]
\vec{a}_3	=0		\hat{x} +	0		\hat{y} +	4.12E-	10	<i>î</i> [m]
Pb Ti		0		$\vec{a}_1 + \vec{a}_1 +$			$\vec{a}_2 + \vec{a}_2 + \vec$	(
e positi	ons of	the ator	ns are	given	in fra	ctional	coodi	nates l	between
								(]
0		0		$\vec{a}_1 +$	-		$\vec{a}_2 +$	-	
0	~	0.5		$\vec{a}_1 +$	0		$\vec{a}_2 +$	0.5	
0	~	0.5		$\vec{a}_1 +$	0.5		$\vec{a}_2 +$	0	
				$\vec{a}_1 +$			$\vec{a}_2 +$		
	~			7 1			$\vec{a}_2 +$		
0				$\vec{a}_1 +$	L				

LEED

Low Energy Electron Diffraction

 $100 \text{ V} \rightarrow k \sim 5 \times 10^{10} \text{ m}^{-1}$



Clean Pd (111)

Pd (111) + 0.3 ML VO_x

LEED is surface sensitive

LEED

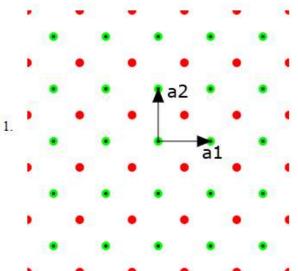
Energy of the electron beam: 100 [eV] Primitive lattice vectors:

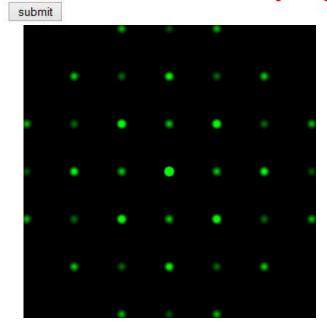
$\vec{a}_1 =$	4.12E-10	\hat{x} +	0	ŷ [m]
$\vec{a}_2 =$	0	x +	4.12E-10	ŷ [m]

Basis:

The positions of the atoms are given in fractional coodinates between -1 and 1.

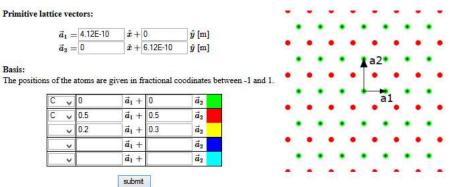
Cs ∨	0	$\vec{a}_1 + $	0	\vec{a}_2
CI 👻	0.5	$\vec{a}_1 +$	0.5	\vec{a}_2
¥		$\vec{a}_1 + $	6	\vec{a}_2
~		$\vec{a}_1 +$		\vec{a}_2
~	S	$\vec{a}_1 + $	6	\vec{a}_2







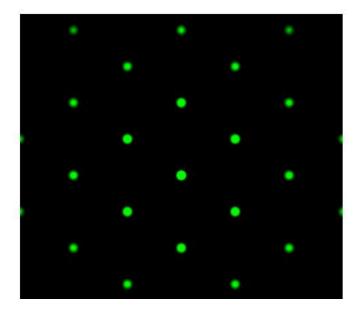
Forbidden reflections



Primitive reciprocal lattice vectors

 $\vec{b}_1 = 2\pi \frac{R \vec{a}_2}{\vec{a}_1 \cdot R \vec{a}_2} = 1.525 \text{e}^{+10} \hat{k}_x + 0.000 \hat{k}_y \text{ [m}^{-1]}$ $\vec{b}_2 = 2\pi \frac{R \vec{a}_1}{\vec{a}_1 \cdot R \vec{a}_2} = 0.000 \hat{k}_x + -1.027 \text{e}^{+10} \hat{k}_y \text{ [m}^{-1]}$ $\text{with} \qquad R = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}$

Low Energy Electron Diffraction



Forbidden reflections

 $n_{u.c.}(ec{r}) = \sum_j Z_j \delta(ec{r} - ec{r}_j).$

Primitive reciprocal lattice vectors

$$\begin{split} \vec{b}_1 &= 2\pi \frac{\vec{a}_2 \times \vec{a}_3}{\vec{a}_1 \cdot (\vec{a}_2 \times \vec{a}_3)} = 3.939 \ \hat{k}_x + -2.275 \text{e}{+}10 \ \hat{k}_y + 0.000 \ \hat{k}_z \ [\text{m}^{-1}] \\ \vec{b}_2 &= 2\pi \frac{\vec{a}_3 \times \vec{a}_1}{\vec{a}_1 \cdot (\vec{a}_2 \times \vec{a}_3)} = 3.939 \ \hat{k}_x + 2.275 \ \hat{k}_y + 0.000 \ \hat{k}_z \ [\text{m}^{-1}] \\ \vec{b}_3 &= 2\pi \frac{\vec{a}_1 \times \vec{a}_2}{\vec{a}_1 \cdot (\vec{a}_2 \times \vec{a}_3)} = 0.000 \ \hat{k}_x + 0.000 \ \hat{k}_y + 1.212 \text{e}{+}10 \ \hat{k}_z \ [\text{m}^{-1}] \end{split}$$

$$n_{ec{G}} = \sum_j Z_j \exp(-iec{G}\cdotec{r}_j).$$

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The value of $ n_{\vec{c}} $ for the
000 diffraction peak is
the total number of
electrons in the primitive
unit cell. The intensities
of the peaks in an x-ray
diffraction experiment
are proportional to
$ n_{\tilde{G}} ^2$. Note that
elements with more
electrons produce
stronger diffraction
intensities.

hkl	$ ec{G} $ Å ⁻¹	$ n_{\vec{G}} $	$ n_{\vec{G}} ^2$	$\operatorname{Re}\{n_{\vec{G}}\}$	$\operatorname{Im}\{n_{\vec{G}}\}$
000	0.000	75.94	5767	75.94	0.000
0-10	4.549e-10	37.87	1434	-37.87	0.02201
010	4.549e-10	37.87	1434	-37.87	-0.02201
0-20	9.098e-10	38.17	1457	-38.17	0.04379
020	9.098e-10	38.17	1457	-38.17	-0.04379
0-30	1.365e-9	75.94	5767	75.94	-0.1318
030	1.365e-9	75.94	5767	75.94	0.1318
0-3-1	1.212	0.3909	0.1528	0.02780	0.3899
0-31	1.212	0.3914	0.1532	-0.02727	0.3904
0-2-1	1.212	42.85	1836	-7.648	42.16
0-21	1.212	42.74	1827	7.551	42.07
0-1-1	1.212	43.01	1850	7.610	-42.33
0-11	1.212	42.96	1845	7,561	-42.29
00-1	1.212	8.896e-8	7.914e-15	-1.573e-8	8.756e-8
001	1.212	8.896e-8	7.914e-15	-1.573e-8	-8.756e-8
01-1	1.212	42.96	1845	-7.561	42.29
011	1.212	43.01	1850	7.610	42.33
02-1	1.212	42.74	1827	7.551	-42.07
021	1.212	42.85	1836	-7.648	-42.16
03_1	1 212	∩ 301⊿	0.1532	_0.02727	_0 3004

Structure factors

Atomic beams

Hydrogen and Helium are used for diffraction studies

$$E = \frac{1}{2}mv^{2} = \frac{p^{2}}{2m} = \frac{\hbar^{2}k^{2}}{2m} = \frac{\hbar^{2}}{2m\lambda^{2}}$$

Low energies can be used for delicate samples. Measure the surface like LEED.