

# Electrons

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# Free electron Fermi gas

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A simple model for a metal is electrons confined to box with periodic boundary conditions.

Like the problem of photons in a box except:

Solve the Schrödinger equation instead of the wave equation.

Electrons are fermions not bosons.

# Free electrons

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$$\left( - \sum_i^N \frac{\hbar^2}{2m} \nabla_i^2 \right) \Psi(\vec{r}_1, \vec{r}_2, \dots, \vec{r}_N) = E \Psi(\vec{r}_1, \vec{r}_2, \dots, \vec{r}_N)$$

No potential, no  $e$ - $e$  interactions, periodic boundary conditions.  
This separates into “molecular orbital’ Hamiltonians

$$-\frac{\hbar^2}{2m} \nabla_i^2 \psi_i(\vec{r}_i) = E_i \psi_i(\vec{r}_i)$$

$$\Psi(\vec{r}_1, \vec{r}_2, \dots, \vec{r}_N) = \mathcal{A}(\psi_1(\vec{r}_1) \uparrow \psi_1(\vec{r}_1) \downarrow \psi_2(\vec{r}_2) \uparrow \psi_2(\vec{r}_2) \downarrow \dots),$$

$$E = E_1 \uparrow + E_1 \downarrow + E_2 \uparrow + E_2 \downarrow + \dots$$

# Free particles in 1-d

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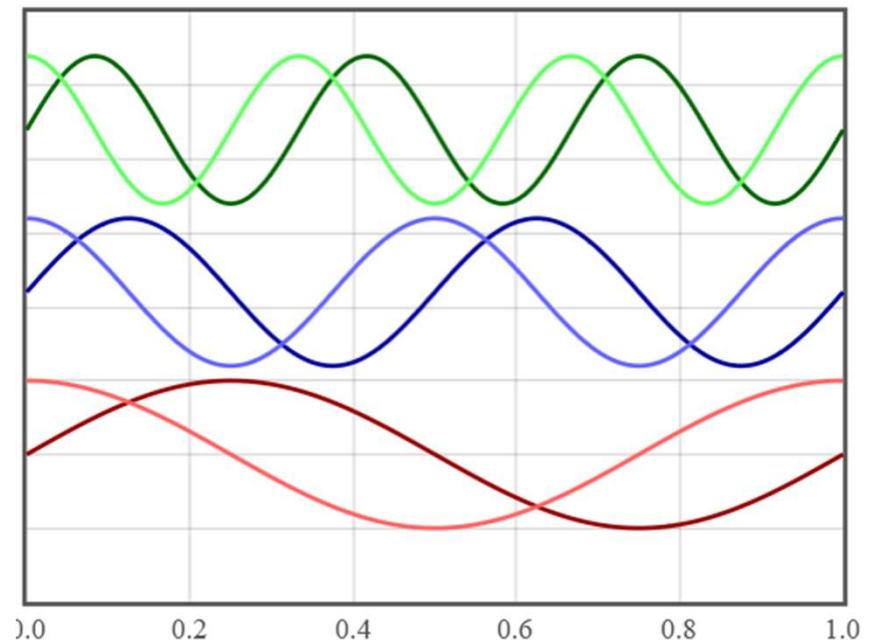
Fill the electrons states like in an atom.

$$-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} + V(x)\psi = E\psi$$

$$V(x) = 0$$

$$\psi_k = \frac{e^{i(kx - \omega t)}}{\sqrt{L}}$$

$$\psi_k^* \psi_k = \frac{1}{L}$$



# Non interacting fermions

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$$\psi = \frac{1}{\sqrt{V}} e^{i\vec{k}\cdot\vec{r}}$$

$$\psi^* \psi = \frac{1}{V}$$

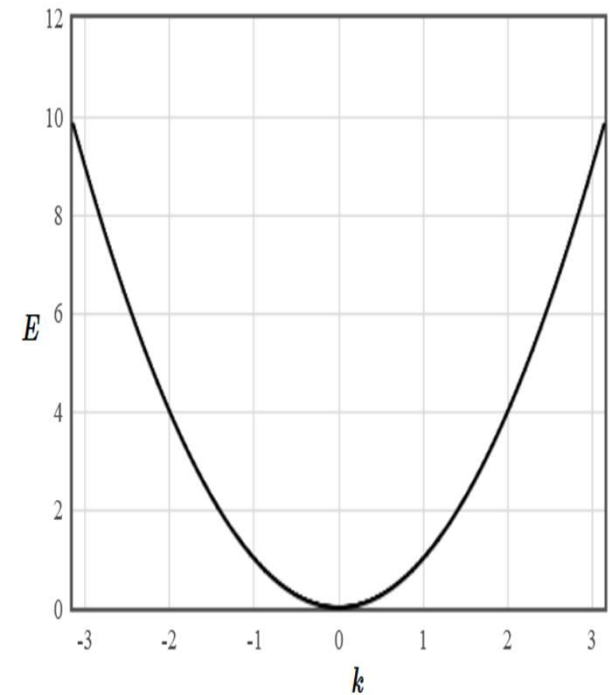
# Free particles in 1-d

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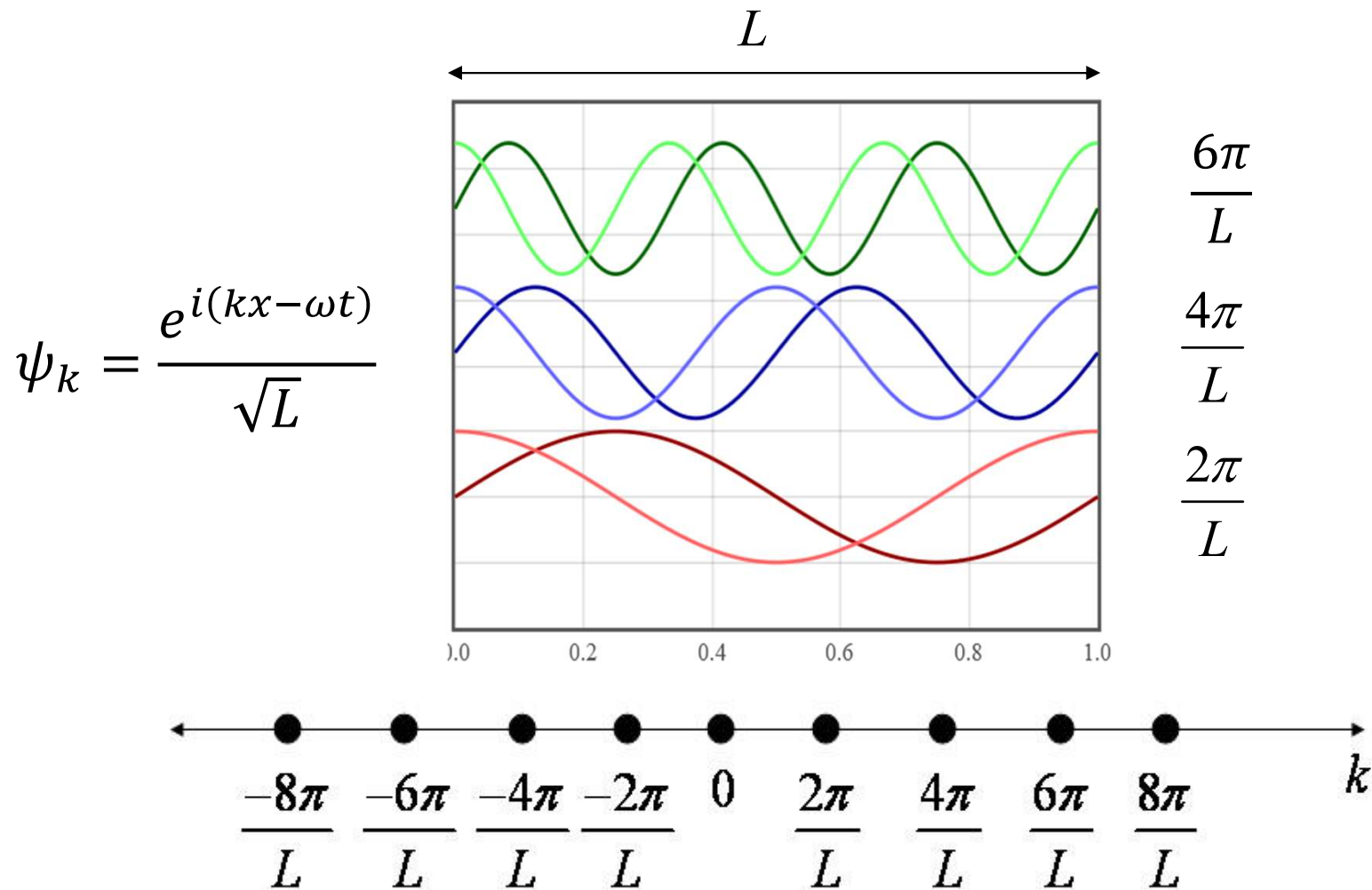
$$i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \frac{d^2 \psi}{dx^2} \quad V = 0$$

Eigen function solutions:  $\psi_k = \frac{e^{i(kx - \omega t)}}{\sqrt{L}}$

Dispersion relation:  $E = \hbar\omega = \frac{\hbar^2 k^2}{2m} = \frac{1}{2} m v^2$

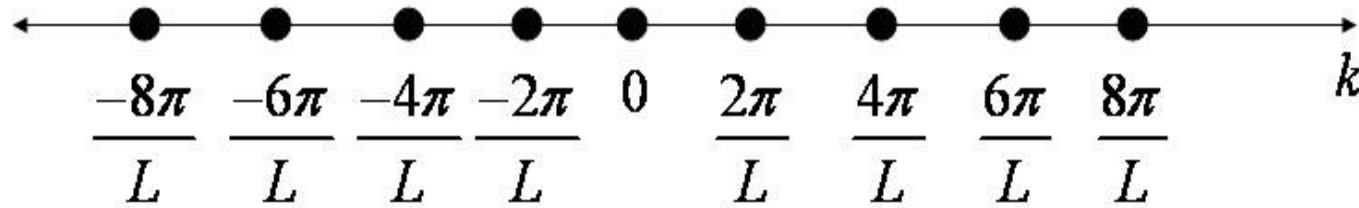


# Periodic boundary conditions



# Periodic boundary conditions

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Left and right                      Spin

$$LD(|k|)dk = \frac{2 \cdot 2dk}{\frac{2\pi}{L}}$$

Density of states in  $k$ :  $D(|k|) = \frac{2}{\pi}$

Number of states between  $|k|$  and  $|k|+dk$  is  $LD(|k|)dk$



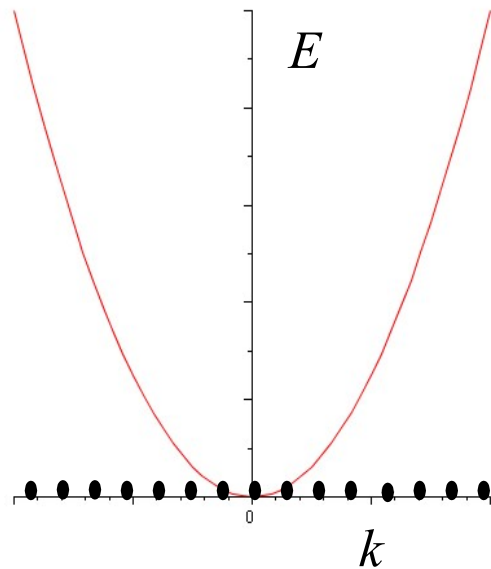
# Free particles in 1-d

$$D(k) = \frac{2}{\pi} \leftarrow \text{Spin}$$

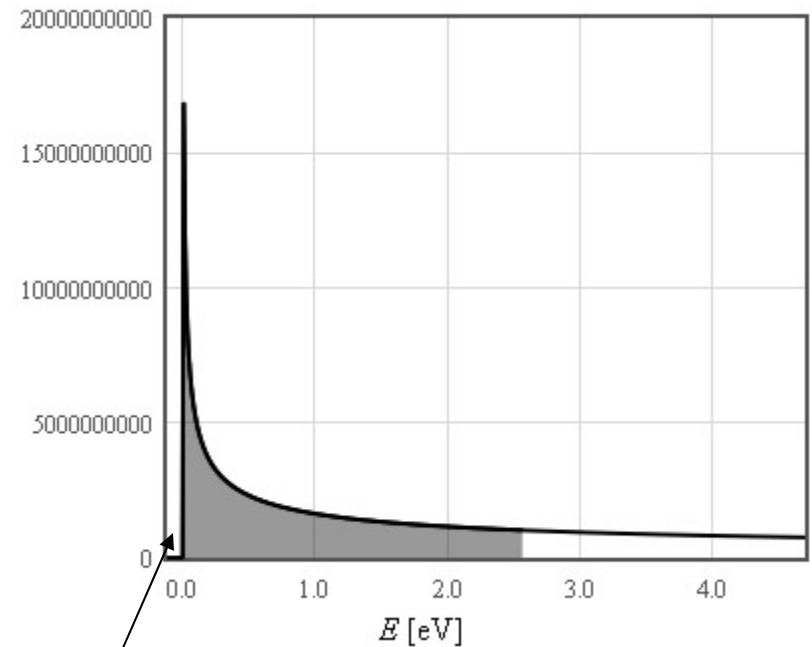
$$k = \frac{\sqrt{2mE}}{\hbar}$$

$$\frac{dk}{dE} = \frac{1}{2\hbar} \sqrt{\frac{2m}{E}}$$

$$D(E) = \frac{1}{\pi\hbar} \sqrt{\frac{2m}{E}}$$



## Density of states



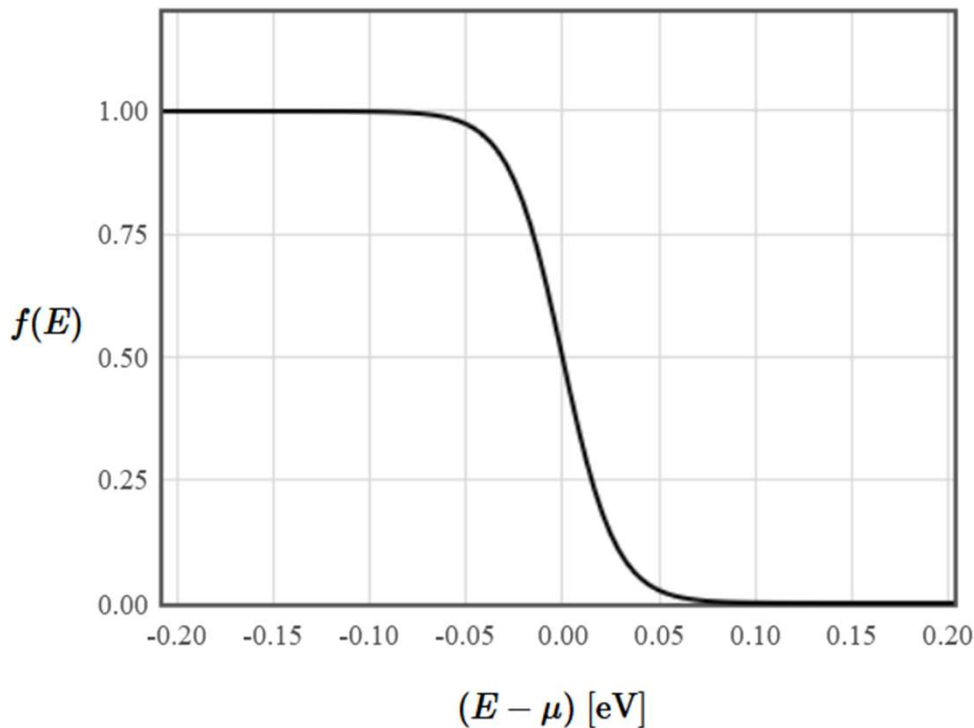
Van Hove singularity

$E$

# Fermi function

$f(E)$  is the probability that a state at energy  $E$  is occupied.

$$f(E) = \frac{1}{1 + \exp\left(\frac{E - \mu}{k_B T}\right)}$$



$\mu$  = chemical potential

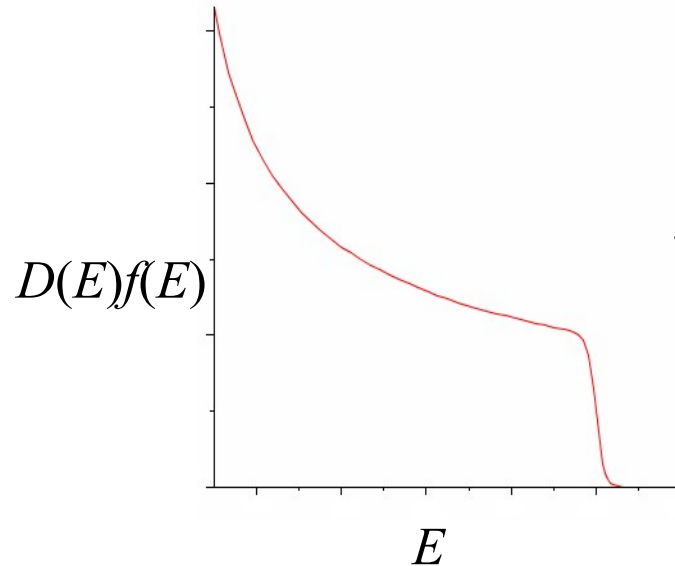


T=160 K

# Chemical potential

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$$D(E) = \frac{1}{\pi\hbar} \sqrt{\frac{2m}{E}}$$



$$f(E) = \frac{1}{1 + \exp\left(\frac{E - \mu}{k_B T}\right)}$$

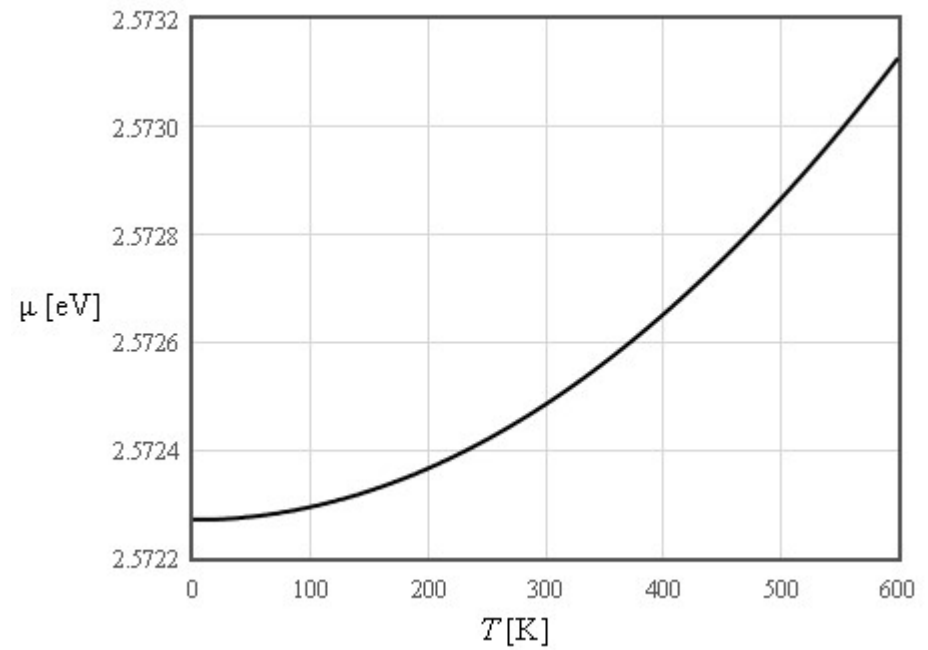
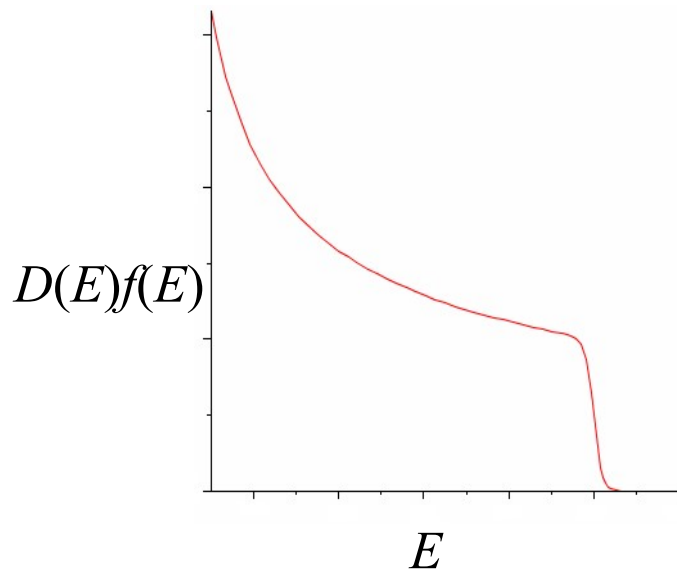
The chemical potential is implicitly defined as the energy that solves the following equation.

$$n = \frac{N}{L} = \int_{-\infty}^{\infty} D(E) f(E) dE = \int_{-\infty}^{\infty} \frac{D(E) dE}{1 + \exp\left(\frac{E - \mu}{k_B T}\right)}$$

Here  $N$  is the total number of electrons.

# Chemical potential

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$\mu$  is temperature dependent

# Fermi energy

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In solid state physics books,

$$E_F = \mu(T=0).$$

In semiconductor books,  $E_F(T) = \mu(T)$ .

$$\text{At } T = 0 \quad n = \int_{-\infty}^{E_F} D(E) dE$$

In one dimension,

$$n = \int_0^{E_F} \frac{1}{\pi \hbar} \sqrt{\frac{2m}{E}} dE = \frac{2}{\pi \hbar} \sqrt{2mE_F}$$

$$E_F = \frac{\pi^2 \hbar^2 n^2}{8m}$$

# Free particles in 1-d

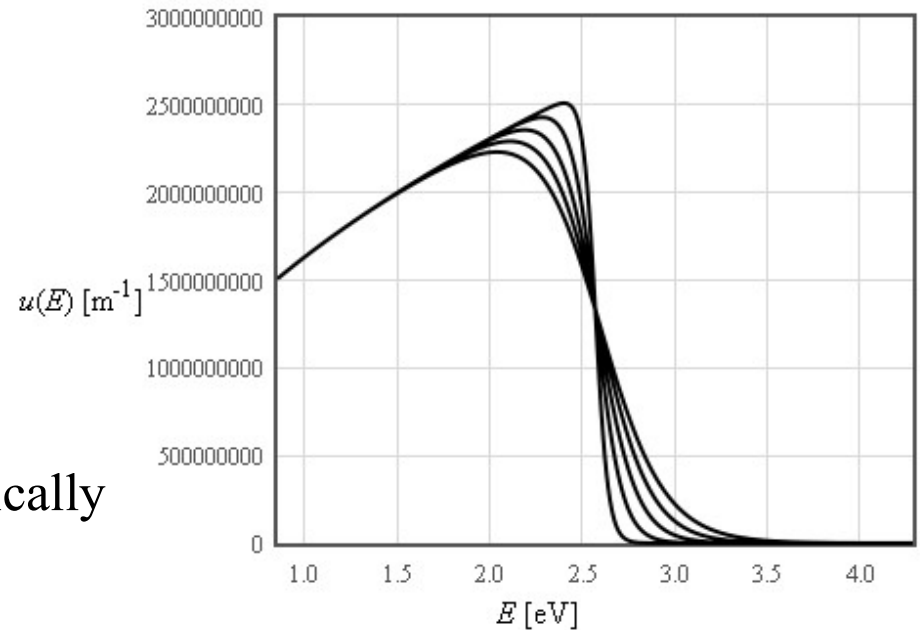
internal energy spectral density

$$u(E) = ED(E)f(E) = \frac{\sqrt{2mE}}{\pi\hbar} \frac{1}{\exp\left(\frac{E-\mu}{k_B T}\right) + 1}$$

$$u = \int_{-\infty}^{\infty} u(E)dE$$

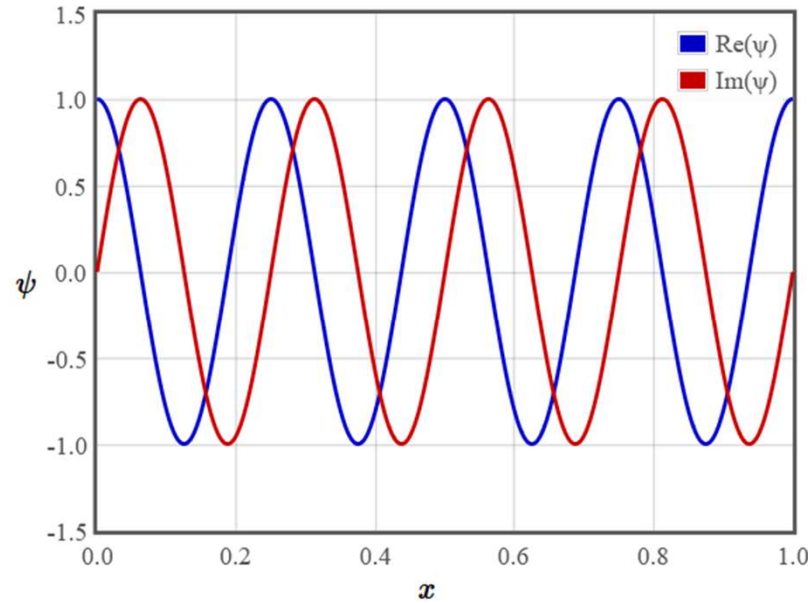
$$c_v = \frac{du}{dT}$$

Not possible to do this integral analytically



analog to the Planck curve for electrons in 1-d

### Free-electron model



$n_x = 4$

Single particle solutions to Schrödinger equation	$\psi(\vec{r}) = \frac{1}{\sqrt{L}} e^{ikx} \frac{1}{\sqrt{m}}$	$\psi(\vec{r}) = \frac{1}{\sqrt{L^2}} e^{i\vec{k}\cdot\vec{r}} \frac{1}{m}$	$\psi(\vec{r}) = \frac{1}{\sqrt{L^3}} e^{i\vec{k}\cdot\vec{r}} \frac{1}{m^{3/2}}$
Allowed k values	$k_x = \frac{2\pi n_x}{L} \frac{1}{m}$ $n_x = \dots - 2, -1, 0, 1, 2, \dots$	$\vec{k} = \left[ \frac{2\pi n_x}{L}, \frac{2\pi n_y}{L} \right] \frac{1}{m}$ $n_x, n_y = \dots - 2, -1, 0, 1, 2, \dots$	$\vec{k} = \left[ \frac{2\pi n_x}{L}, \frac{2\pi n_y}{L}, \frac{2\pi n_z}{L} \right] \frac{1}{m}$ $n_x, n_y, n_z = \dots - 2, -1, 0, 1, 2, \dots$
Density of electron states in reciprocal space	$D(\vec{k}) = \frac{2}{2\pi}$	$D(\vec{k}) = \frac{2}{(2\pi)^2}$	$D(\vec{k}) = \frac{2}{(2\pi)^3}$
Density of electron states in reciprocal space	$D(k) = \frac{2}{\pi}$	$D(k) = \frac{k}{\pi} \text{ m}^{-1}$	$D(k) = \frac{k^2}{\pi} \text{ m}^{-2}$

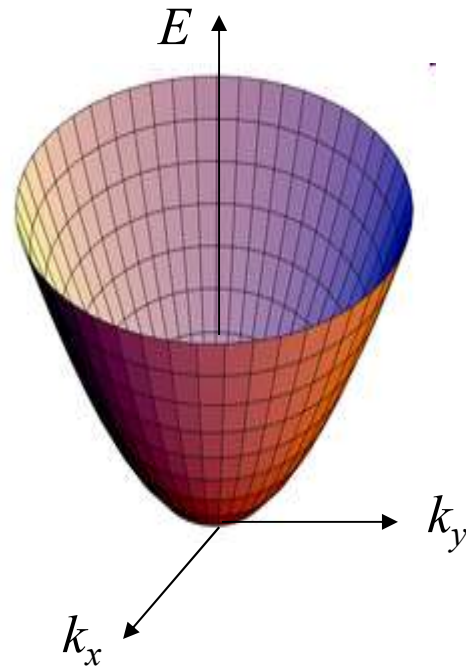
# Free particles in 2-d

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Density of states

$$E = \frac{\hbar^2 k^2}{2m}$$

$$k = \frac{\sqrt{2mE}}{\hbar}$$



$$\frac{dk}{dE} = \frac{1}{2\hbar} \sqrt{\frac{2m}{E}}$$

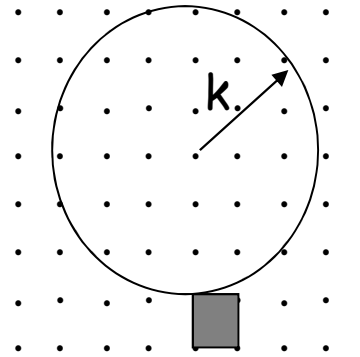


# Free particles in 2-d

$$E = \frac{\hbar^2 k^2}{2m}$$

$$k = \frac{\sqrt{2mE}}{\hbar}$$

$$\frac{dk}{dE} = \frac{1}{2\hbar} \sqrt{\frac{2m}{E}}$$



$$\left(\frac{2\pi}{L}\right)^2$$

$$D(E) = D(|k|) \frac{dk}{dE}$$

$$L^2 D(|k|) dk = \frac{2 \cdot 2\pi k dk}{\left(\frac{2\pi}{L}\right)^2}$$

$$D(|k|) = \frac{k}{\pi} \text{ m}^{-1}$$

$$D(E) = \frac{k}{\pi} \frac{dk}{dE} = \frac{\sqrt{2mE}}{\hbar\pi} \frac{1}{2\hbar} \sqrt{\frac{2m}{E}}$$

$$D(E) = \frac{m}{\pi\hbar^2}$$

# Free particles in 2-d

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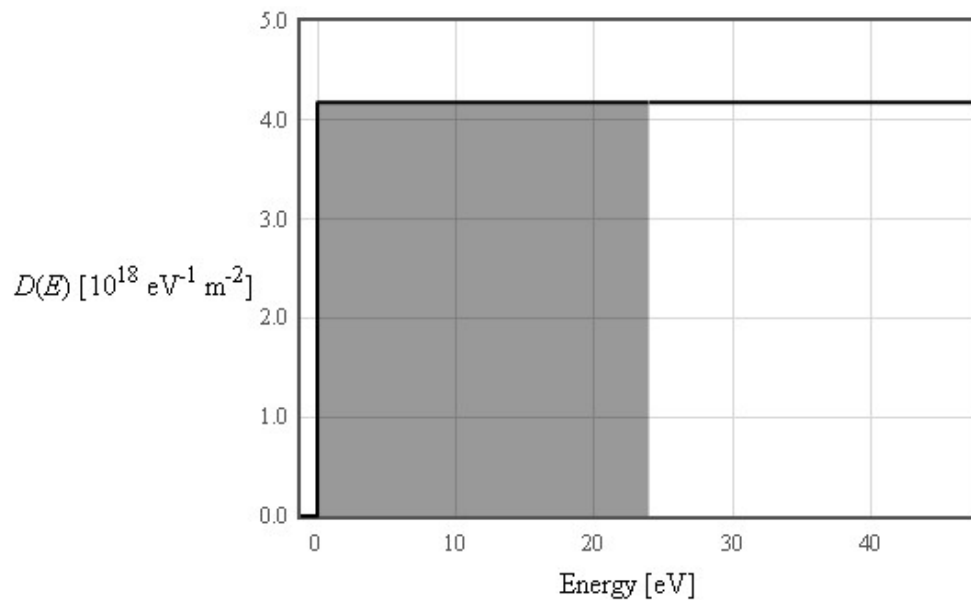
$$D(E) = \frac{m}{\pi \hbar^2}$$

At  $T = 0$ :

$$n = \int_0^{E_F} D(E) dE$$

$$n = \frac{N}{L^2} = \frac{m}{\pi \hbar^2} \int_0^{E_F} dE = \frac{m}{\pi \hbar^2} E_F$$

$$E_F = \frac{\pi \hbar^2 n}{m}$$



# Fermi circle

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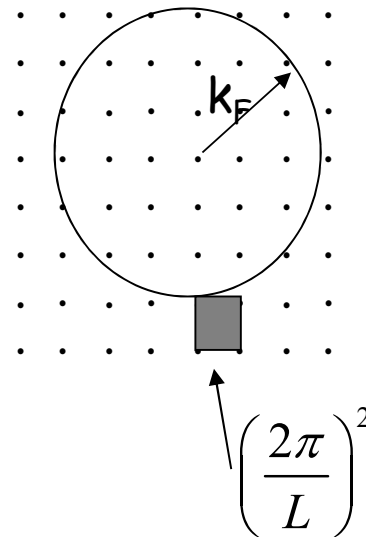
$$N = \frac{2\pi k_F^2}{\left(\frac{2\pi}{L}\right)^2}$$

spin ↘

$n = N/L^2 =$  electron density

$$k_F = \sqrt{2\pi n}$$

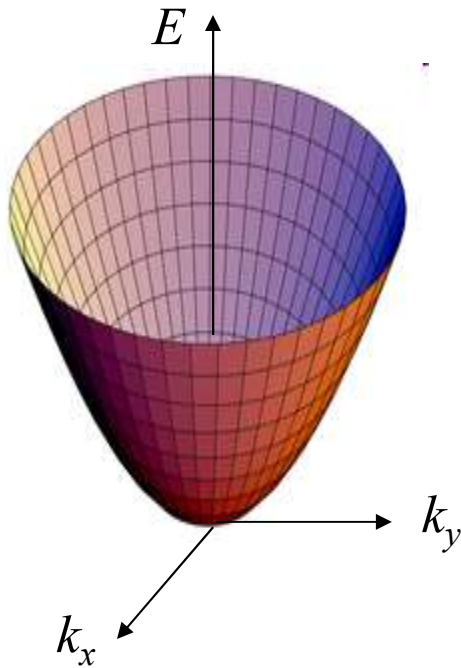
$$E_F = \frac{\hbar^2 k_F^2}{2m} = \frac{\pi \hbar^2 n}{m}$$



At  $T = 0$ , all states inside the Fermi circle are occupied and those outside are empty.

# Free particles in 3-d

$$E = \frac{\hbar^2 k^2}{2m}$$



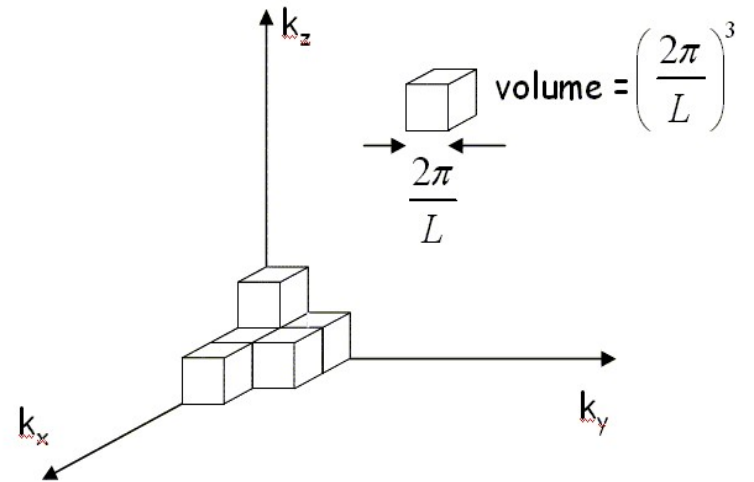
Density of states

$$D(|k|) = \frac{k^2}{\pi^2}$$

$$\frac{dk}{dE} = \frac{1}{2\hbar} \sqrt{\frac{2m}{E}}$$

$$k^2 = \frac{2mE}{\hbar^2}$$

$$D(E) = D(|k|) \frac{dk}{dE}$$

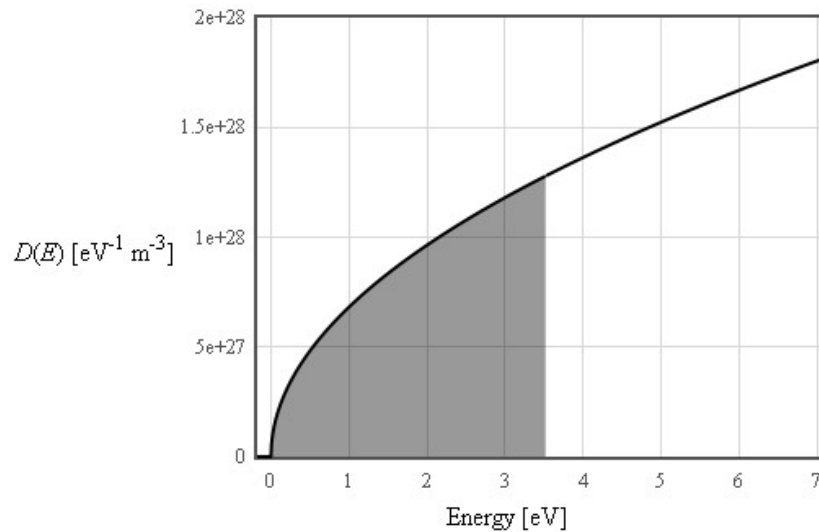


$$k_x, k_y, k_z = \dots, \frac{-4\pi}{L}, \frac{-2\pi}{L}, 0, \frac{2\pi}{L}, \frac{4\pi}{L}, \dots$$

$$D(E) = \frac{(2m)^{\frac{3}{2}}}{2\pi^2 \hbar^3} \sqrt{E}$$

# Free particles in 3-d

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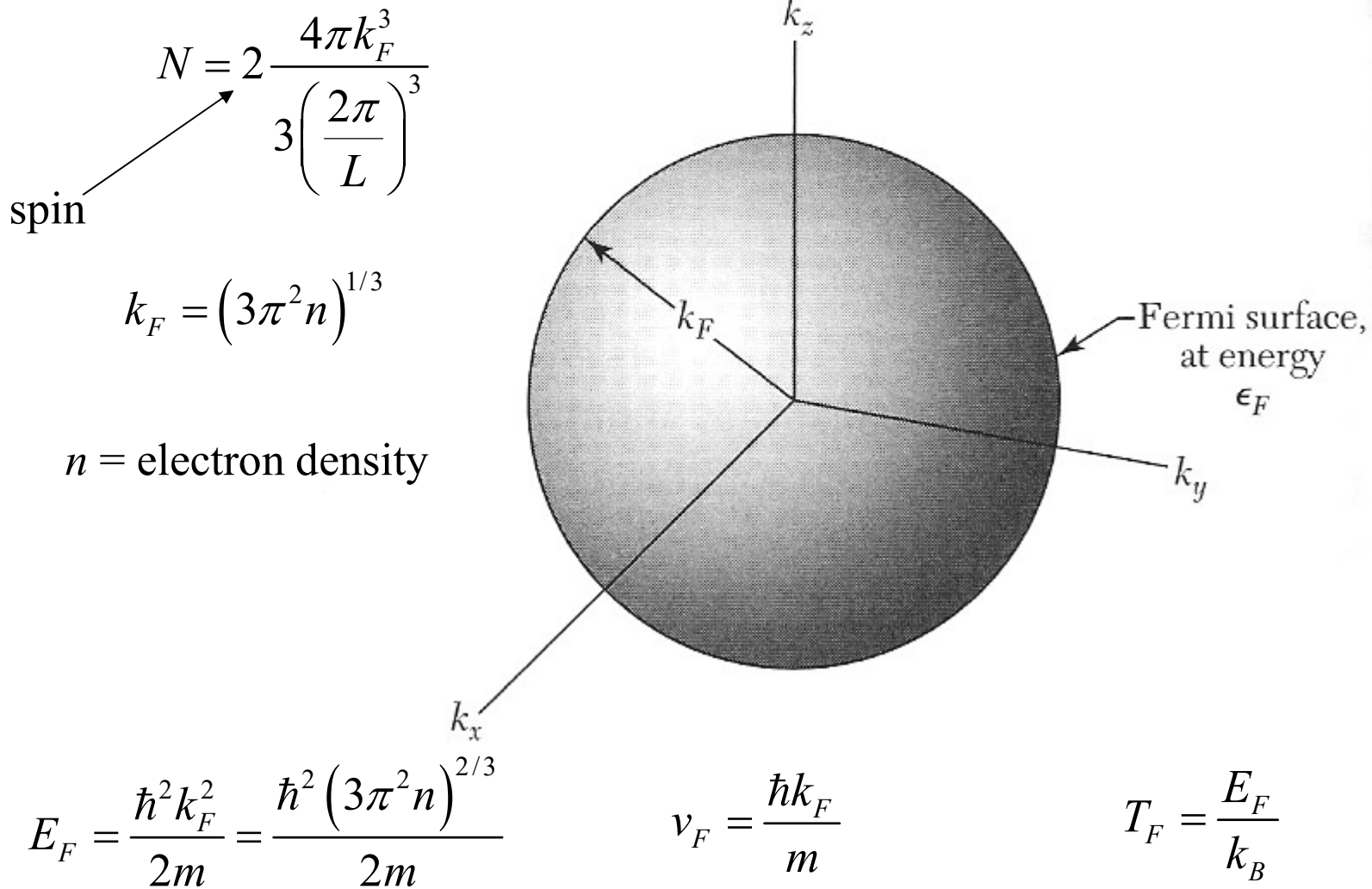
At  $T = 0$ :

$$n = \int_0^{E_F} D(E) dE$$

$$n = \frac{N}{L^3} = \frac{\sqrt{2} m^{3/2}}{\pi^2 \hbar^3} \int_0^{E_F} \sqrt{E} dE = \frac{(2m)^{3/2}}{3\pi^2 \hbar^3} E_F^{3/2}$$

$$E_F = \frac{\hbar^2}{2m} (3\pi^2 n)^{2/3}$$

# Fermi sphere



The thermal and electronic properties depend on the states at the Fermi surface.

**Table 1** Calculated free electron Fermi surface parameters for metals at room temperature

(Except for Na, K, Rb, Cs at 5 K and Li at 78 K)

Valency	Metal	Electron concentration, in $\text{cm}^{-3}$	Radius <sup>a</sup> parameter $r_n$	Fermi wavevector, in $\text{cm}^{-1}$	Fermi velocity, in $\text{cm s}^{-1}$	Fermi energy, in eV	Fermi temperature $T_F \equiv \epsilon_F/k_B$ , in deg K
1	Li	$4.70 \times 10^{22}$	3.25	$1.11 \times 10^8$	$1.29 \times 10^8$	4.72	$5.48 \times 10^4$
	Na	2.65	3.93	0.92	1.07	3.23	3.75
	K	1.40	4.86	0.75	0.86	2.12	2.46
	Rb	1.15	5.20	0.70	0.81	1.85	2.15
	Cs	0.91	5.63	0.64	0.75	1.58	1.83
	Cu	8.45	2.67	1.36	1.57	7.00	8.12
	Ag	5.85	3.02	1.20	1.39	5.48	6.36
	Au	5.90	3.01	1.20	1.39	5.51	6.39
2	Be	24.2	1.88	1.93	2.23	14.14	16.41
	Mg	8.60	2.65	1.37	1.58	7.13	8.27
	Ca	4.60	3.27	1.11	1.28	4.68	5.43
	Sr	3.56	3.56	1.02	1.18	3.95	4.58
	Ba	3.20	3.69	0.98	1.13	3.65	4.24
	Zn	13.10	2.31	1.57	1.82	9.39	10.90
	Cd	9.28	2.59	1.40	1.62	7.46	8.66
3	Al	18.06	2.07	1.75	2.02	11.63	13.49
	Ga	15.30	2.19	1.65	1.91	10.35	12.01
	In	11.49	2.41	1.50	1.74	8.60	9.98
4	Pb	13.20	2.30	1.57	1.82	9.37	10.87
	Sn( <i>w</i> )	14.48	2.23	1.62	1.88	10.03	11.64

<sup>a</sup>The dimensionless radius parameter is defined as  $r_n = r_0/a_H$ , where  $a_H$  is the first Bohr radius and  $r_0$  is the radius of a sphere that contains one electron.

$$k_F = (3\pi^2 n)^{1/3}$$

$$E_F \gg k_B T$$

# Free electron Fermi gas

$$1 - d \quad D(E) = \sqrt{\frac{2m}{\hbar^2 \pi^2 E}} = \frac{n}{2\sqrt{E_F E}} \quad \text{J}^{-1} \text{m}^{-1}$$

$$2 - d \quad D(E) = \frac{m}{\hbar^2 \pi} = \frac{n}{E_F} \quad \text{J}^{-1} \text{m}^{-2}$$

$$3 - d \quad D(E) = \frac{\pi}{2} \left( \frac{2m}{\hbar^2 \pi^2} \right)^{3/2} \sqrt{E} = \frac{3n}{2E_F^{3/2}} \sqrt{E} \quad \text{J}^{-1} \text{m}^{-3}$$

