

Landau theory, susceptibility

Add a magnetic field

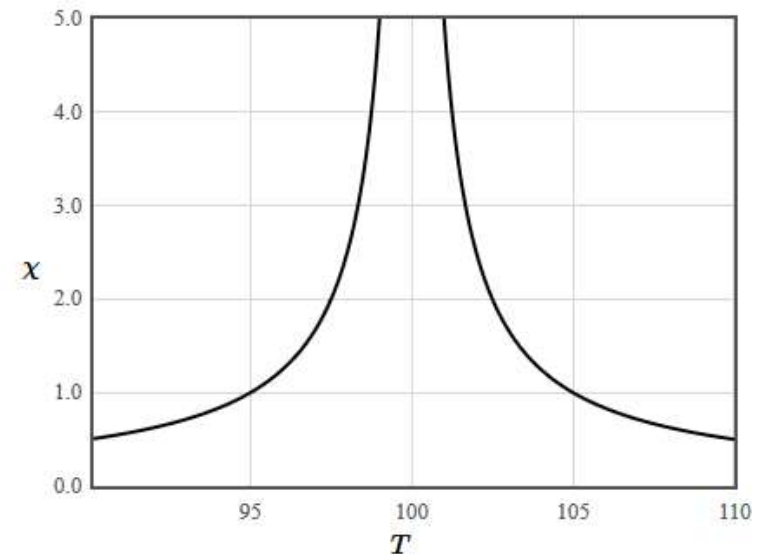
$$f = f_0 + \alpha_0 (T - T_c) m^2 + \frac{1}{2} \beta m^4 - mB$$

$$\frac{df}{dm} = 2\alpha_0 (T - T_c) m + 2\beta m^3 - B = 0$$

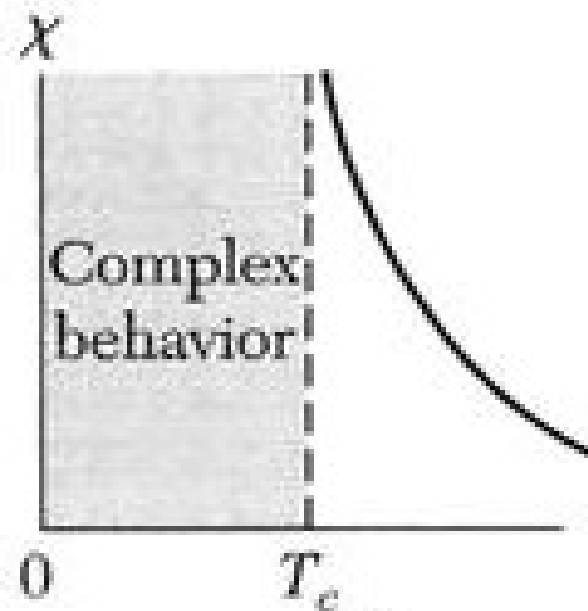
Above T_c , m is finite for finite B . For small m ,

$$m = \frac{B}{2\alpha_0 (T - T_c)} \quad T > T_c$$

$$\chi = \frac{1}{2\alpha_0 (T - T_c)} \quad \text{Curie-Weiss}$$



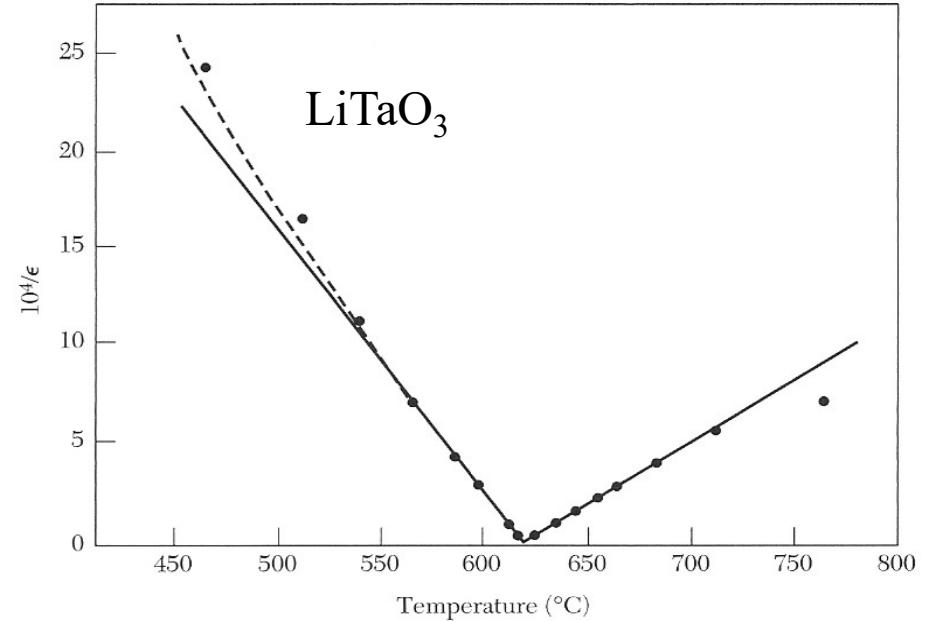
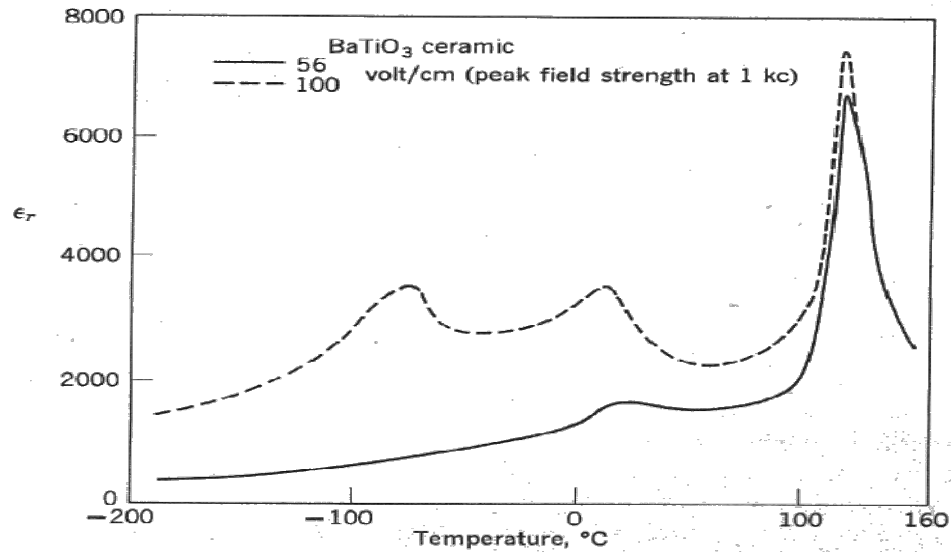
Ferromagnetism



$$\chi = \frac{C}{T - T_c}$$

Curie-Weiss law
($T > T_c$)

Landau theory of phase transitions

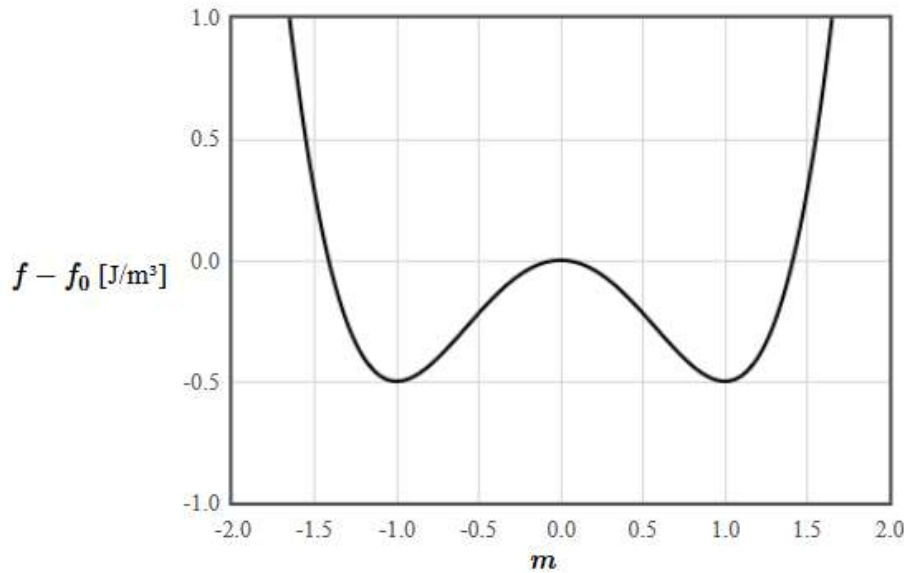


$$\epsilon_r = 1 + \chi$$

$$\chi = \frac{1}{2\alpha_0 (T - T_c)}$$

Curie-Weiss law

Fitting the α_0 and β parameters

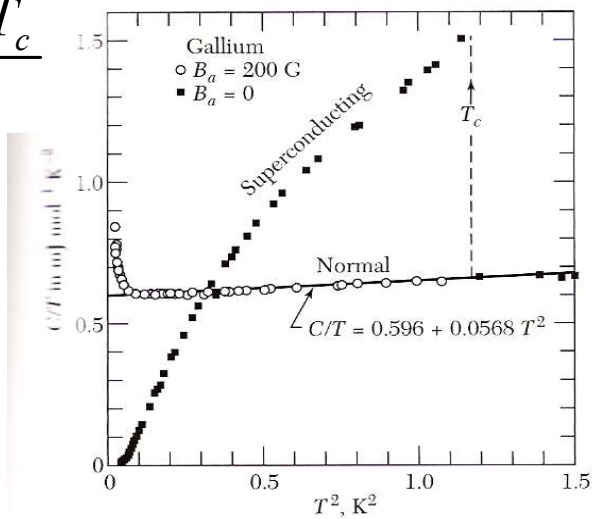


$\alpha_0 =$
 $\beta =$
 $T =$
 $T_c =$
 $f_0(T) =$

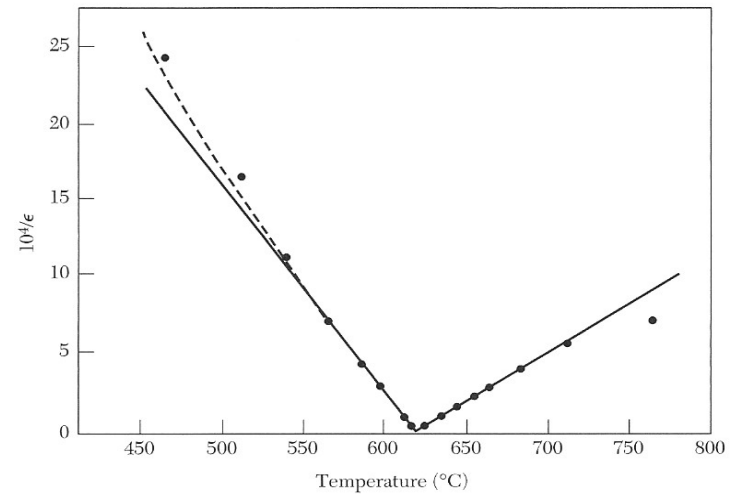
Superconductivity

Ferromagnetism

$$\Delta c_v = \frac{2\alpha_0^2 T_c}{\beta}$$

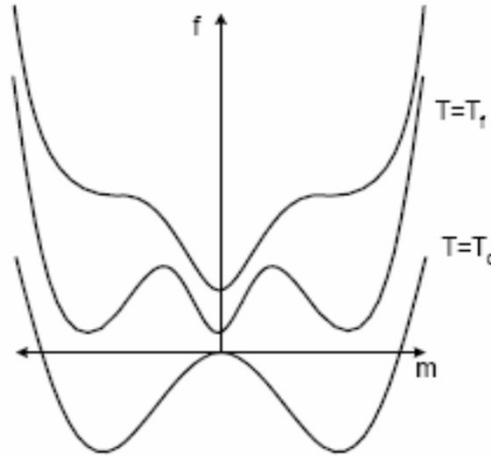


$$\chi = \frac{1}{2\alpha_0(T - T_c)}$$



First order transitions

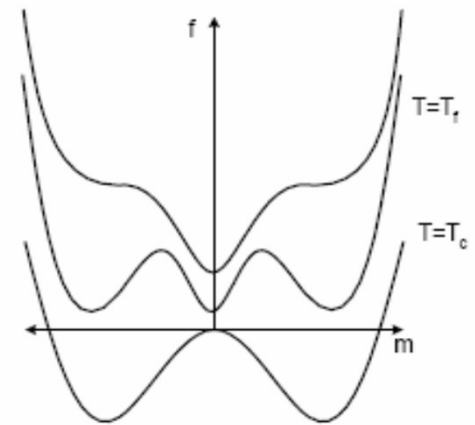
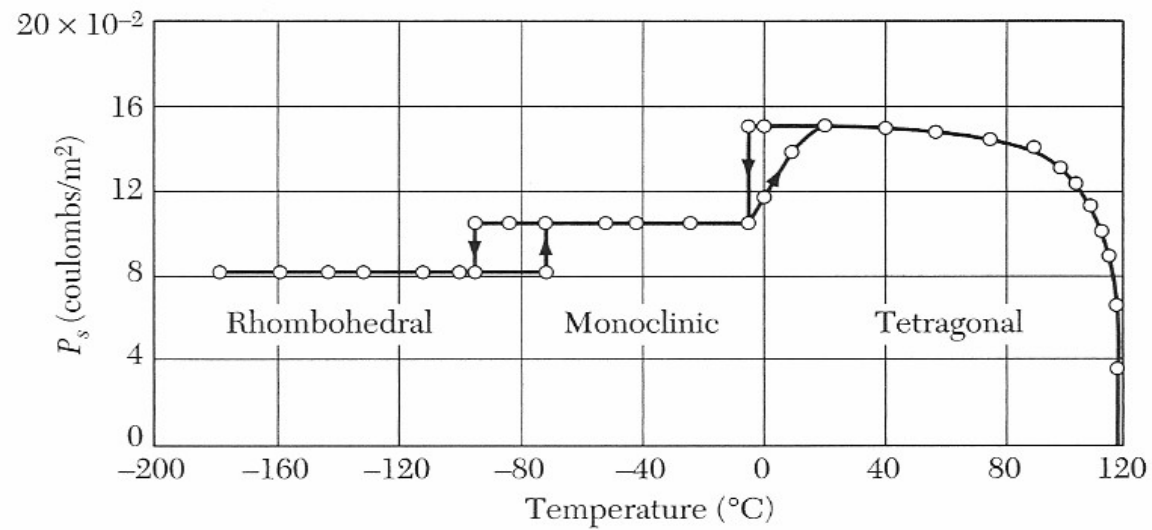
$$f = f_0 + \alpha_0 (T - T_c) m^2 + \frac{1}{2} \beta m^4 + \frac{1}{3} \gamma m^6 + \dots \quad \alpha_0 > 0, \quad \beta < 0, \quad \gamma > 0$$



There is a jump in the order parameter at the phase transition.

First order transitions

BaTiO₃



First order transitions

$$f = f_0 + \alpha_0 (T - T_c) m^2 + \frac{1}{2} \beta m^4 + \frac{1}{3} \gamma m^6 + \dots \quad \beta < 0$$

$$\frac{df}{dm} = 2\alpha_0 (T - T_c) m + 2\beta m^3 + 2\gamma m^5 = 0$$

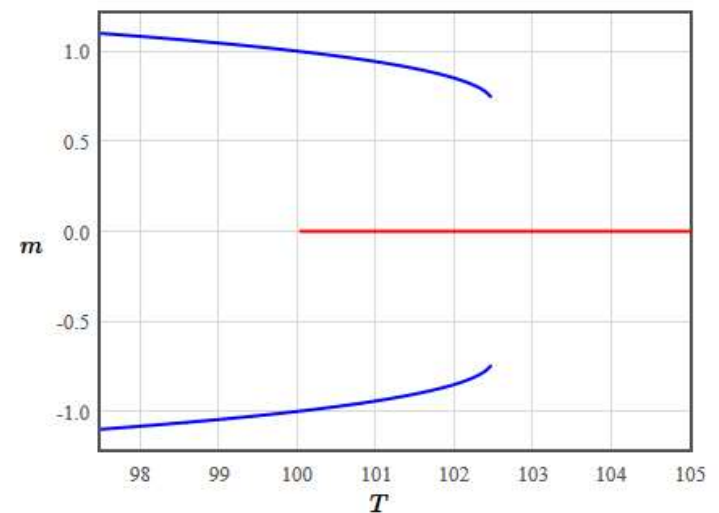
One solution for $m = 0$.

$$\alpha_0 (T - T_c) + \beta m^2 + \gamma m^4 = 0$$

$$m^2 = 0, \frac{-\beta \pm \sqrt{\beta^2 - 4\alpha_0 (T - T_c) \gamma}}{2\gamma}$$

There will be a minimum at finite m as long as m^2 is real

$$T_1 = \frac{\beta^2}{4\alpha_0 \gamma} + T_c$$



Jump in the order parameter

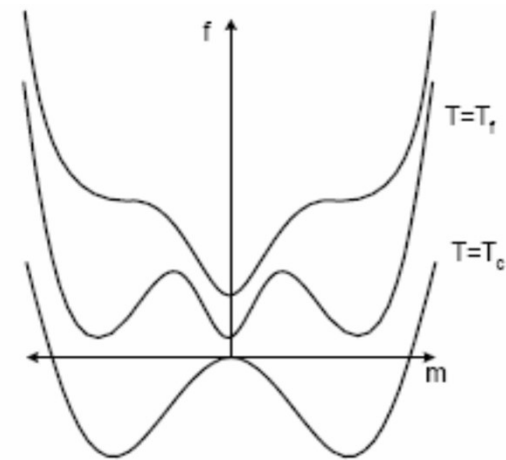
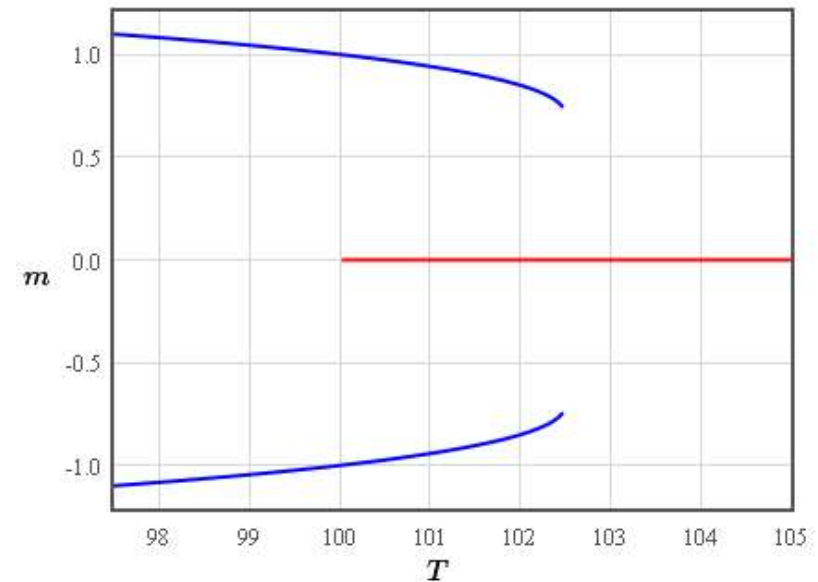
$$m^2 = \frac{-\beta \pm \sqrt{\beta^2 - 4\alpha_0(T - T_c)\gamma}}{2\gamma}$$

At T_c

$$\Delta m = \sqrt{\frac{-\beta}{\gamma}}$$

At T_1

$$\Delta m = \sqrt{\frac{-\beta}{2\gamma}}$$



First order transitions, entropy, c_v

$$f = f_0 + \alpha_0 (T - T_c) m^2 + \frac{1}{2} \beta m^4 + \frac{1}{3} \gamma m^6 + \dots \quad \beta < 0$$

$$m = 0, \pm \sqrt{\frac{-\beta \pm \sqrt{\beta^2 - 4\alpha_0 (T - T_c) \gamma}}{2\gamma}}$$

$$s = -\frac{\partial f}{\partial T} = \frac{\alpha_0}{2\gamma} \left(\beta - \sqrt{\beta^2 - 4\alpha_0 \gamma (T - T_c)} \right)$$

$$c_v = T \frac{\partial s}{\partial T} = \frac{\alpha_0^2 T}{\sqrt{\beta^2 - 4\alpha_0 \gamma (T - T_c)}}$$

branch where the order parameter is nonzero

First order transitions, susceptibility

$$f = f_0 + \alpha_0 (T - T_c) m^2 + \frac{1}{2} \beta m^4 + \frac{1}{3} \gamma m^6 - mB \quad \beta < 0$$

$$\frac{df}{dm} = 2\alpha_0 (T - T_c) m + 2\beta m^3 + 2\gamma m^5 - B = 0$$

At the minima $B = 2\alpha_0 (T - T_c) m + 2\beta m^3 + 2\gamma m^5$

For small m ,

$$\chi = \left. \frac{dm}{dB} \right|_{m=0} = \frac{1}{2\alpha_0 (T - T_c)} \quad \text{Curie - Weiss}$$

$$\chi = \left. \frac{dm}{dB} \right|_{m=\sqrt{\frac{-\beta}{2\gamma}}} = \frac{1}{2\alpha_0 (T - T_1)}$$

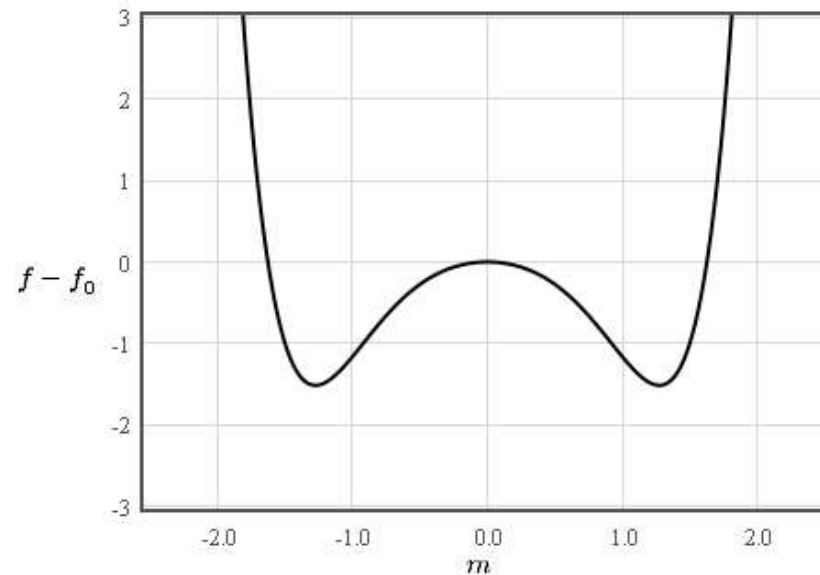
- Outline
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Landau theory of a first order phase transition

The free energy for a first order transition in Landau theory is,

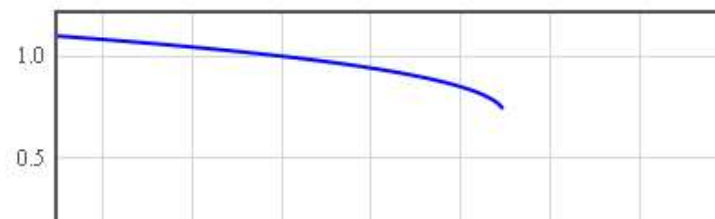
$$f(T) = f_0(T) + \alpha_0(T - T_c)m^2 + \frac{1}{2}\beta m^4 + \frac{1}{3}\gamma m^6 \quad \alpha_0 > 0, \quad \beta < 0, \quad \gamma > 0.$$

Here $f_0(T)$ describes the temperature dependence of the high temperature phase near the phase transition.



$\alpha_0 =$
 $\beta =$
 $\gamma =$
 $T =$
 $T_c =$
 $f_0(T) =$

Order parameter



Superconductivity

Primary characteristic: zero resistance at dc

There is a critical temperature T_c above which superconductivity disappears

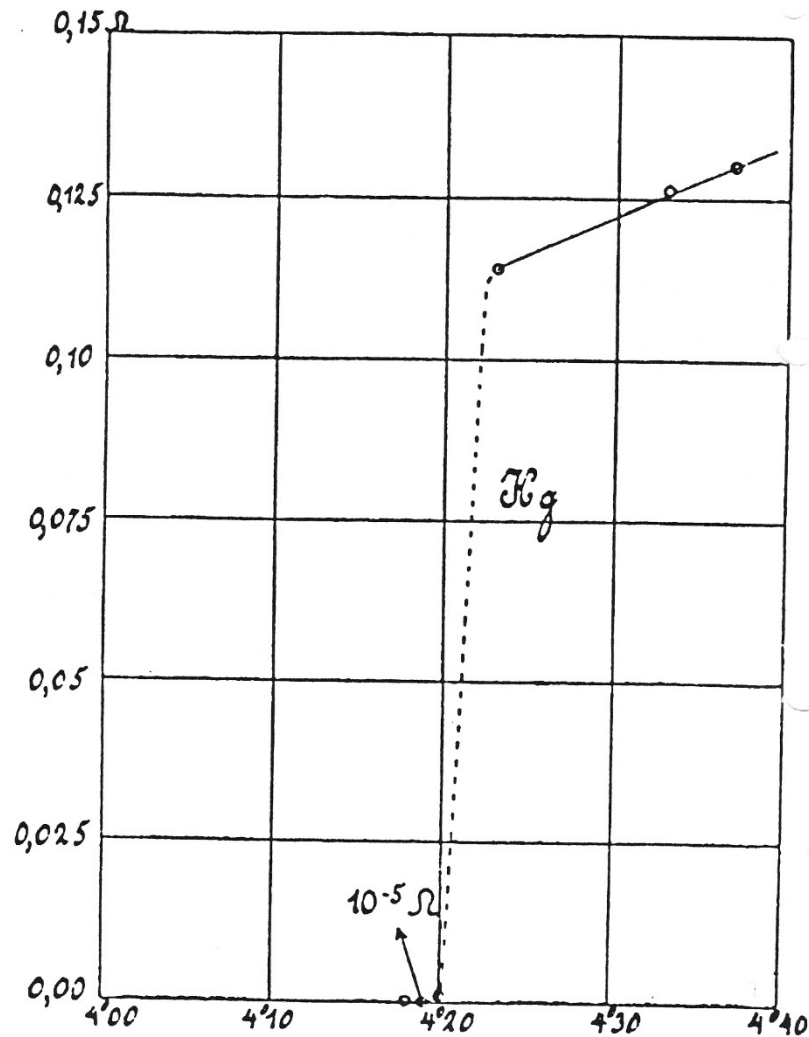
About 1/3 of all metals are superconductors

Metals are usually superconductors OR magnetic, not both

Good conductors are bad superconductors

Kittel chapter 10

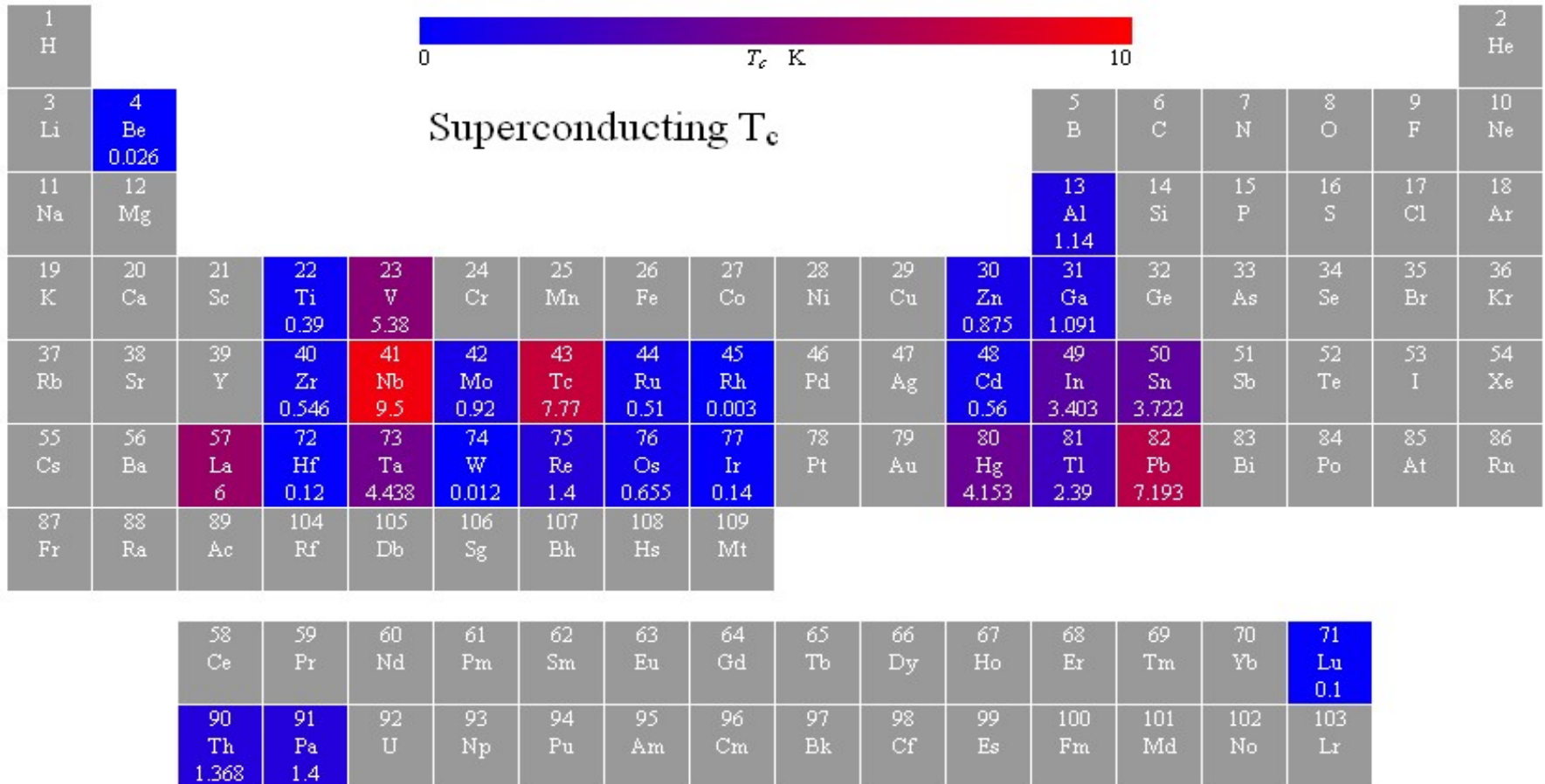
Superconductivity



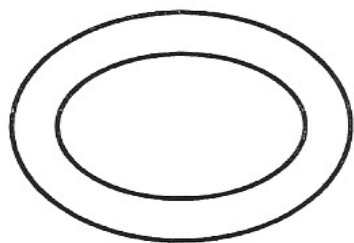
Heike Kamerling-Onnes

Superconductivity was discovered in 1911

Critical temperature



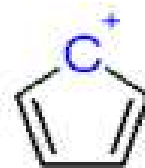
Superconductivity



Superconducting ring



A



B



C

Molecule with magnetic moment

Antiaromatic molecules are unstable and highly reactive

No measurable decay in current after 2.5 years. $\rho < 10^{-25} \Omega\text{m}$.