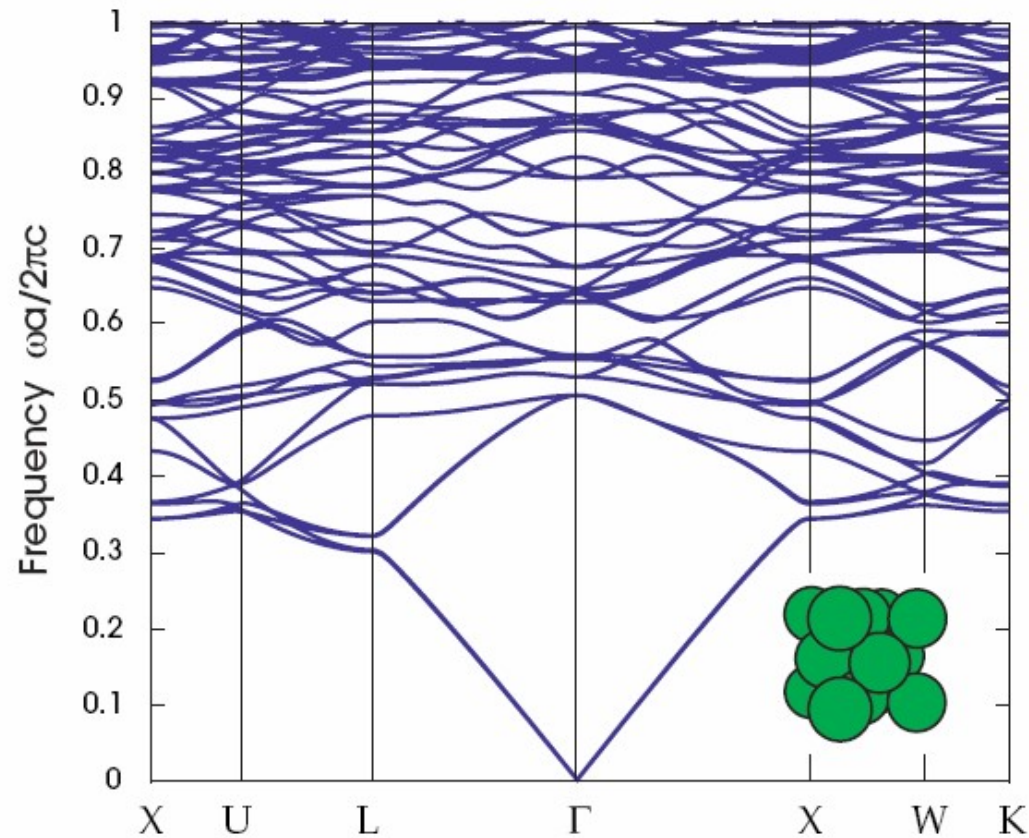


# Photonic Crystals \ Electrons

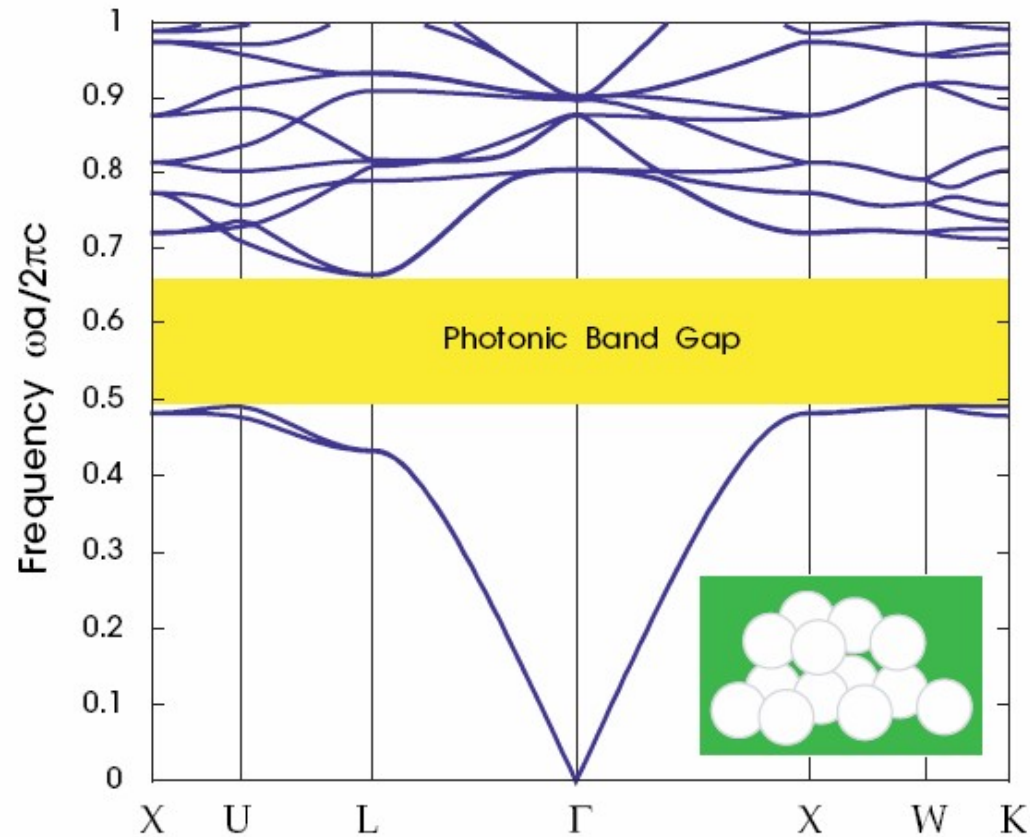
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fcc



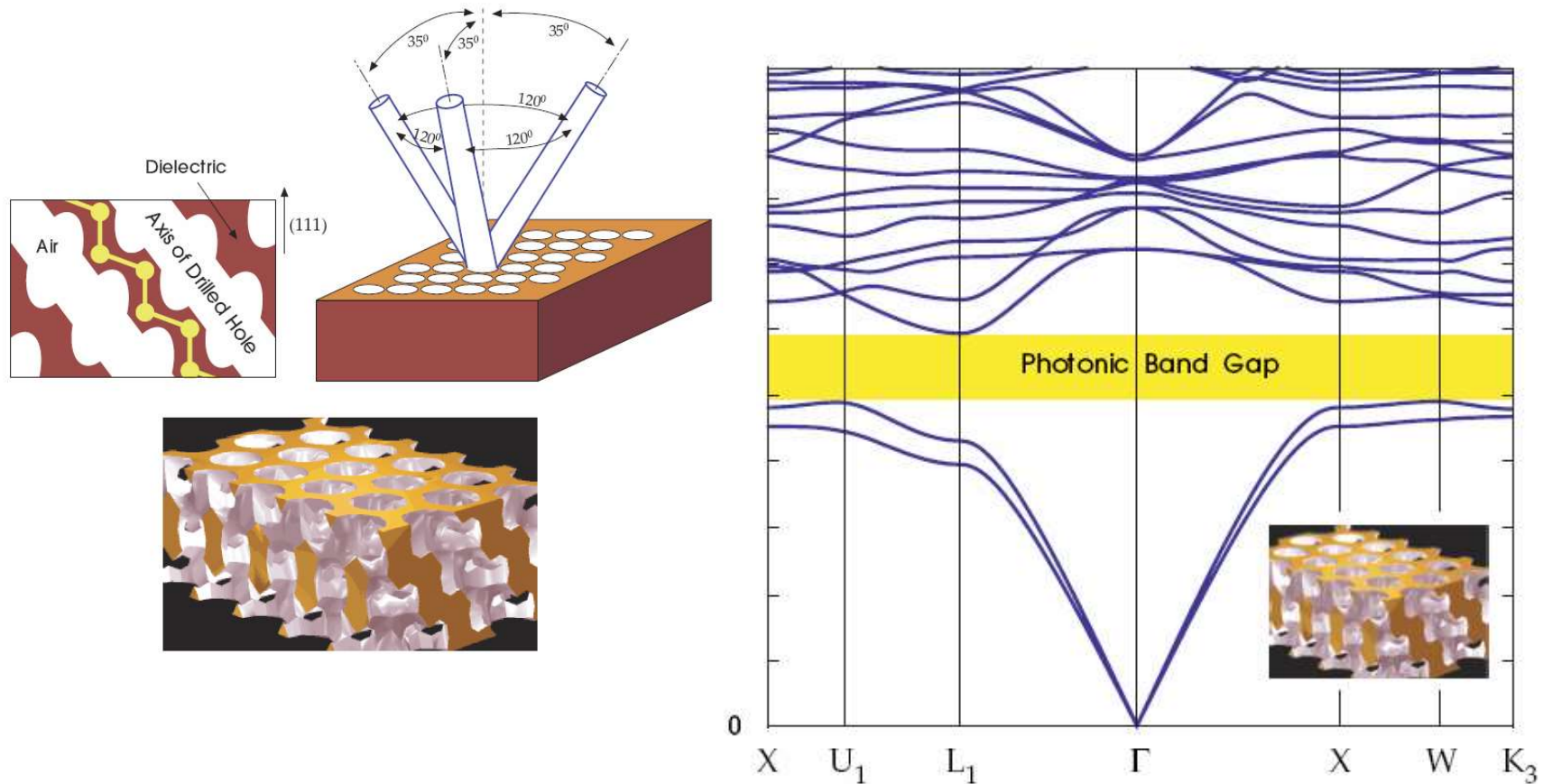
**Figure 2:** The photonic band structure for the lowest-frequency electromagnetic modes of a face-centered cubic (fcc) lattice of close-packed dielectric spheres ( $\epsilon = 13$ ) in air (inset). Note the *absence* of a complete photonic band gap. The wave vector varies across the irreducible Brillouin zone between the labelled high-symmetry points; see appendix B for a discussion of the Brillouin zone of an fcc lattice.

# diamond



**Figure 3:** The photonic band structure for the lowest bands of a diamond lattice of air spheres in a high dielectric ( $\epsilon = 13$ ) material (inset). A complete photonic band gap is shown in yellow. The wave vector varies across the irreducible Brillouin zone between the labelled high-symmetry points; see appendix B for a discussion of the Brillouin zone of an fcc lattice.

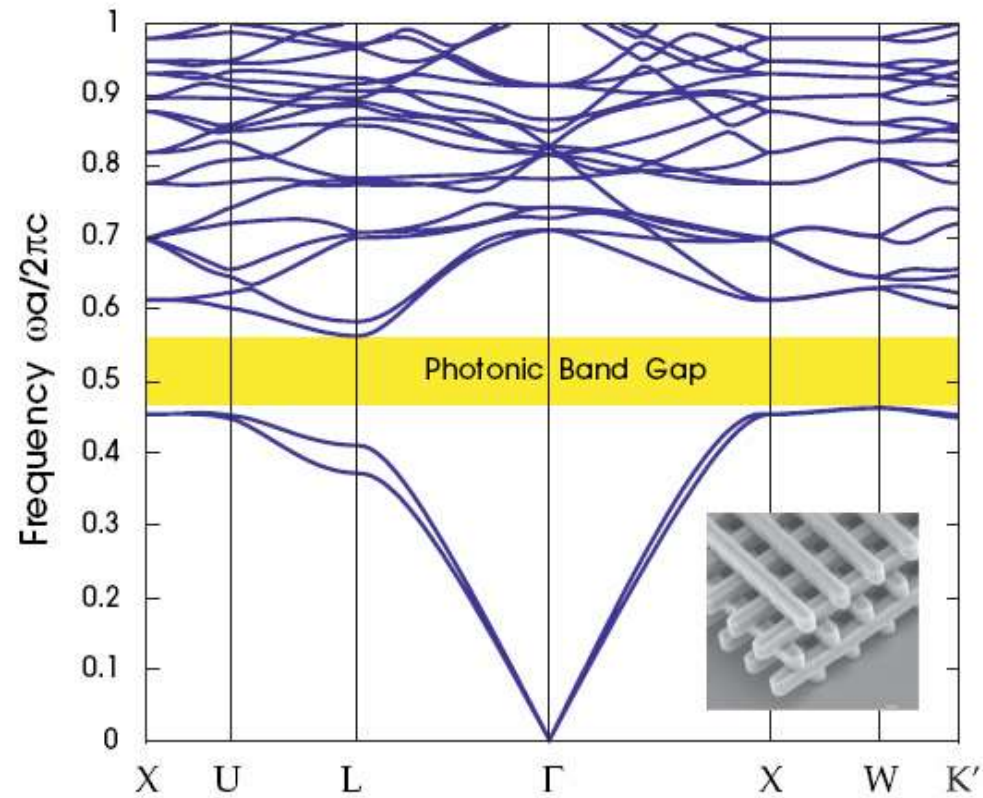
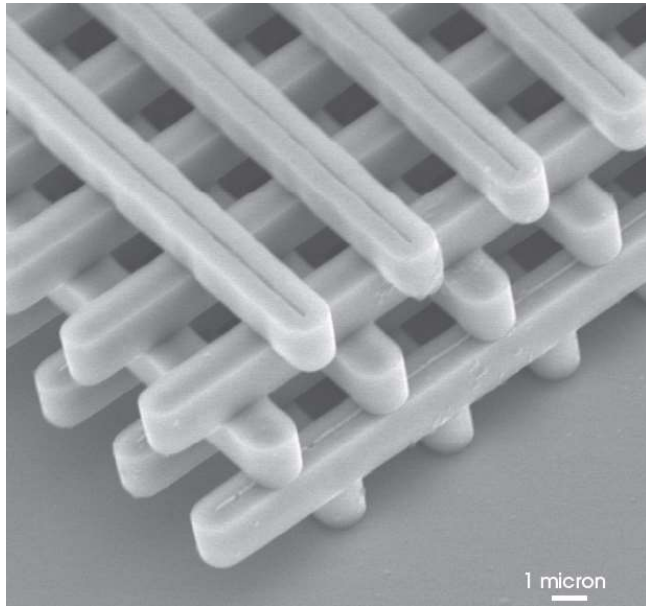
# Yablonovite



**Figure 5:** The photonic band structure for the lowest bands of Yablonovite (inset, from figure 4). Wave vectors are shown for a portion of the irreducible Brillouin zone that includes the edges of the complete gap (yellow). A detailed discussion of this band structure can be found in Yablonovitch et al. (1991a).

<http://ab-initio.mit.edu/book/>

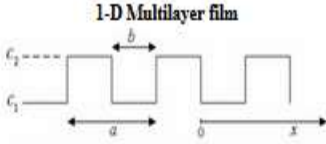
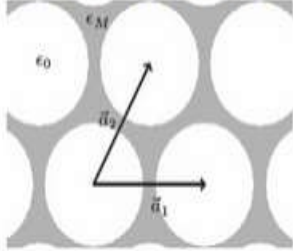

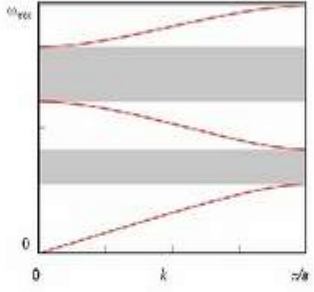
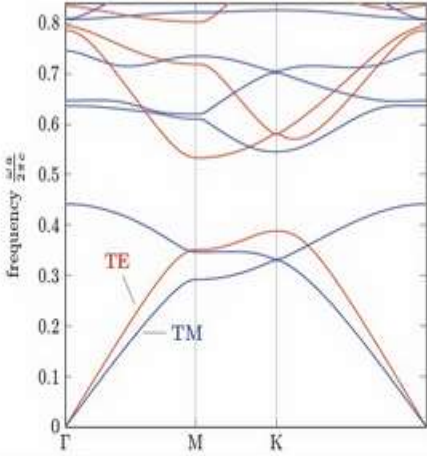
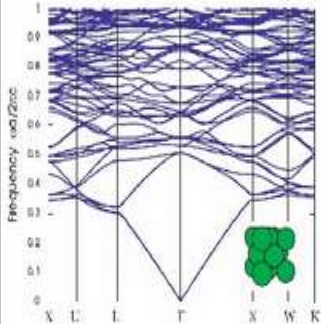
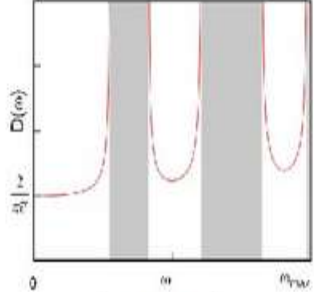
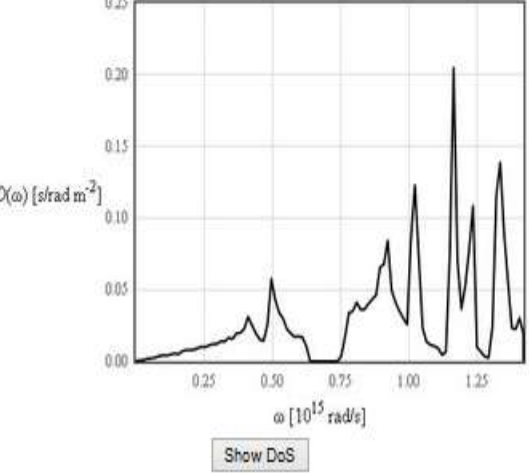
# Woodpile



**Figure 7:** The photonic band structure for the lowest bands of the woodpile structure (inset, from figure 6) with  $\epsilon = 13$  logs in air. The irreducible Brillouin zone is larger than that of the fcc lattice described in appendix B, because of reduced symmetry—only a portion is shown, including the edges of the complete photonic band gap (yellow).

<http://ab-initio.mit.edu/book/>

# Photonic crystals

<p><b>1-D Multilayer film</b></p>  <p><u>1-D Multilayer film</u></p>	<p><b>2-D triangular array of air holes</b></p> 	<p><b>fcc</b></p>  <p><a href="http://ab-initio.mit.edu/book">http://ab-initio.mit.edu/book</a></p>	
<p><b>Dispersion relation</b></p>	 <p>Calculate <math>\omega(k)</math></p>		 <p><a href="http://ab-initio.mit.edu/book">http://ab-initio.mit.edu/book</a></p>
<p><b>Density of states</b></p>	 <p>Calculate DoS</p>	 <p>Show DoS</p>	

# Student projects

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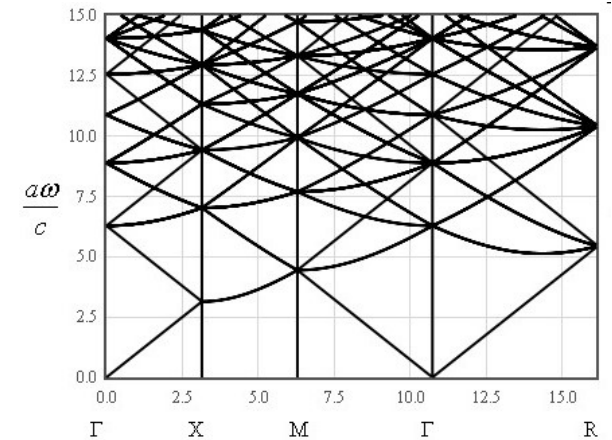
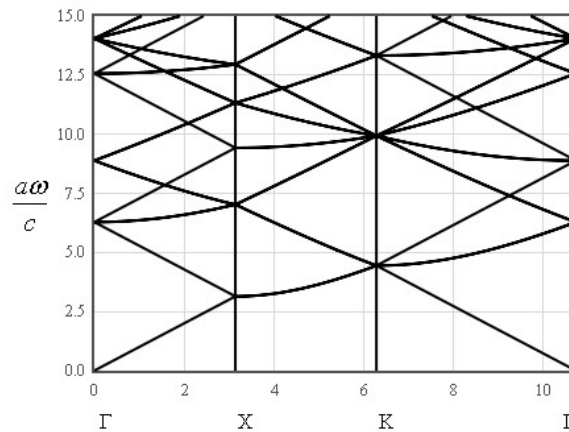
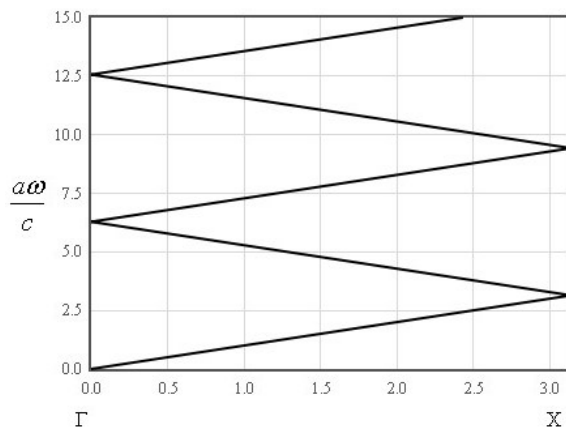
Describe the plane wave method for the web page.

Calculate the band structure and density of states for a photonic crystal.

Help complete the table of the empty lattice approximation

Write a program that calculates the photonic band structure of any 1-D crystal.

Plot the thermodynamic properties of some photonic crystal (you need the density of states)



# Electrons

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Thermodynamic properties

Free electrons

Fermi surfaces

Band structure calculations

- Empty lattice approximation

- Plane wave method

- Tight binding



# Thermodynamic properties of non-interacting fermions

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The grand canonical partition function is

$$Z_{gr} = \sum_q \exp\left(\frac{\mu}{k_B T}\right)^{N_q} \exp\left(-\frac{E_q}{k_B T}\right) = \sum_q \exp\left(-\frac{E_q - \mu N_q}{k_B T}\right)$$

Here  $q$  sums over the macro states. Only one fermion is allowed per microscopic quantum state.

$$N_q = \sum_i n_{qi} \quad E_q = \sum_i n_{qi} \varepsilon_i \quad n_{qi} \in 0,1$$

$n_{qi}$  are occupation numbers that specify if microstate  $i$  is occupied in macrostate  $q$

<http://lamp.tu-graz.ac.at/~hadley/ss2/fermigas/thermo/thermo.php>

# Thermodynamic properties of non-interacting fermions

---

$$Z_{gr} = \sum_q \exp\left(-\frac{E_q - \mu N_q}{k_B T}\right) = \sum_q \exp\left(-\frac{\sum_i n_{qi} (\varepsilon_i - \mu)}{k_B T}\right) = \sum_q \prod_i \exp\left(-\frac{n_{qi} (\varepsilon_i - \mu)}{k_B T}\right)$$

The sum over all possible macrostates can also be written as the sum over all possible microstates.

$$Z_{gr} = \sum_{n_1=0}^1 \sum_{n_2=0}^1 \cdots \sum_{n_j=0}^1 \cdots \prod_i \exp\left(-\frac{n_i (\varepsilon_i - \mu)}{k_B T}\right) = \prod_i \left( \exp\left(-\frac{(\varepsilon_i - \mu)}{k_B T}\right) + 1 \right)$$

Pull the  $n_i$  factors through the other sums then write out the sum over 0 and 1.

# Thermodynamic properties of non-interacting fermions

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Grand potential:  $\Phi = U - TS - \mu N = -k_B T \ln(Z_{gr})$

$$\Phi = -k_B T \sum_i \ln \left( \exp \left( -\frac{(\varepsilon_i - \mu)}{k_B T} \right) + 1 \right)$$

Approximate the sum by an integral over the density of states.

$$\phi = \frac{\Phi}{V} = -k_B T \int_{-\infty}^{\infty} D(E) \ln \left( \exp \left( -\frac{(E - \mu)}{k_B T} \right) + 1 \right) dE$$

Differentiate to find the number density.

$$n = -\frac{\partial \phi}{\partial \mu} = \int_{-\infty}^{\infty} \frac{D(E)}{1 + \exp \left( \frac{E - \mu}{k_B T} \right)} dE$$

Fermi function

# Thermodynamic properties

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Grand potential density:

$$\phi = -k_B T \int_{-\infty}^{\infty} D(E) \ln \left( \exp \left( -\frac{(E - \mu)}{k_B T} \right) + 1 \right) dE$$

Helmholtz free energy density:

$$f = \phi + \mu n = \int_{-\infty}^{\infty} D(E) \left( \frac{\mu}{1 + \exp \left( \frac{E - \mu}{k_B T} \right)} - k_B T \ln \left( \exp \left( -\frac{(E - \mu)}{k_B T} \right) + 1 \right) \right) dE$$

Entropy density:

$$s = -\frac{\partial \phi}{\partial T} = \frac{1}{T} \int_{-\infty}^{\infty} D(E) \left( \frac{(E - \mu)}{1 + \exp \left( \frac{(E - \mu)}{k_B T} \right)} - k_B T \ln \left( \exp \left( -\frac{(E - \mu)}{k_B T} \right) + 1 \right) \right) dE$$

# Thermodynamic properties

Chemical potential  
(implicitly defined by):

$$n = \int_{-\infty}^{\infty} \frac{D(E)}{1 + \exp\left(\frac{E - \mu}{k_B T}\right)} dE$$

DoS →  
mu

Internal energy density:

$$u = \phi + Ts + \mu n = \int_{-\infty}^{\infty} \frac{ED(E)}{1 + \exp\left(\frac{E - \mu}{k_B T}\right)} dE$$

DoS →  
u(T)

Energy spectral density:

$$u(E, T) = \frac{ED(E)}{1 + \exp\left(\frac{E - \mu}{k_B T}\right)}$$

DoS →  
u(E)

Specific heat:

$$c_v = \frac{\partial u}{\partial T} = \int_{-\infty}^{\infty} \frac{ED(E)(E - \mu) \exp\left(\frac{E - \mu}{k_B T}\right)}{k_B T^2 \left(1 + \exp\left(\frac{E - \mu}{k_B T}\right)\right)^2} dE$$

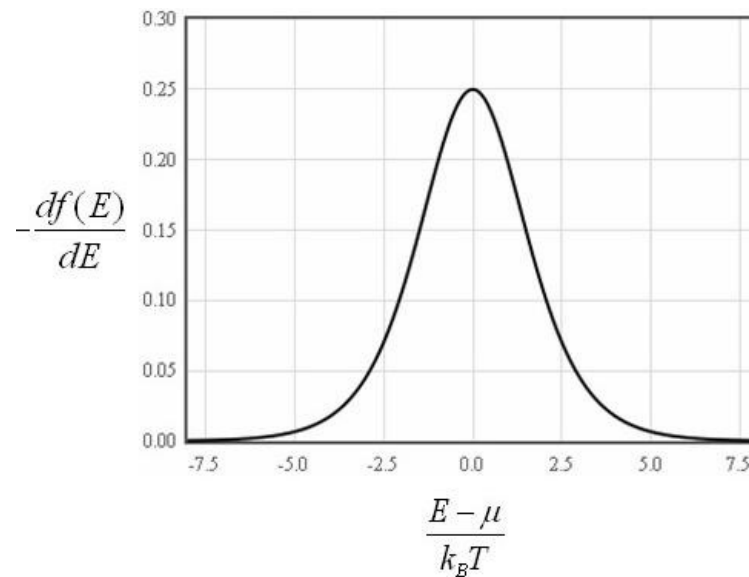
DoS →  
cv(T)

# Properties of metals depend mostly on the electron states at the Fermi surface

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$$n = \int_{-\infty}^{\infty} D(E) f(E) dE = \int_{-\infty}^{\infty} \frac{D(E) dE}{\exp\left(\frac{E - \mu}{k_B T}\right) + 1} \quad K(E) = \int_{-\infty}^E D(E') dE'$$

$$n = \int_{-\infty}^{\infty} D(E) f(E) dE = K(\infty) f(\infty) - K(-\infty) f(-\infty) - \int_{-\infty}^{\infty} K(E) \frac{df(E)}{dE} dE.$$



# Free electron Fermi gas 1-d

$$D(k) = \frac{2}{\pi}$$

$$E = \frac{\hbar^2 k^2}{2m} \quad k = \sqrt{\frac{2mE}{\hbar^2}}$$

$$dE = \frac{\hbar^2 k}{m} dk$$

$$D(k)dk = D(E)dE$$

$$D(E) = \sqrt{\frac{2m}{\hbar^2 \pi^2 E}} = \frac{n}{2\sqrt{E_F E}} \quad \text{J}^{-1} \text{m}^{-1}$$

