

# Phase transitions

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# Anharmonic terms

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Expand the energy in terms of the normal modes of the linearized problem  $u_k$

$$U = U_0 + \frac{\partial U}{\partial u_k} u_k + \frac{1}{2} \frac{\partial^2 U}{\partial u_j \partial u_k} u_j u_k + \frac{1}{6} \frac{\partial^3 U}{\partial u_i \partial u_j \partial u_k} u_i u_j u_k + \frac{1}{24} \frac{\partial^4 U}{\partial u_h \partial u_i \partial u_j \partial u_k} u_h u_i u_j u_k + \dots$$

Thermal expansion

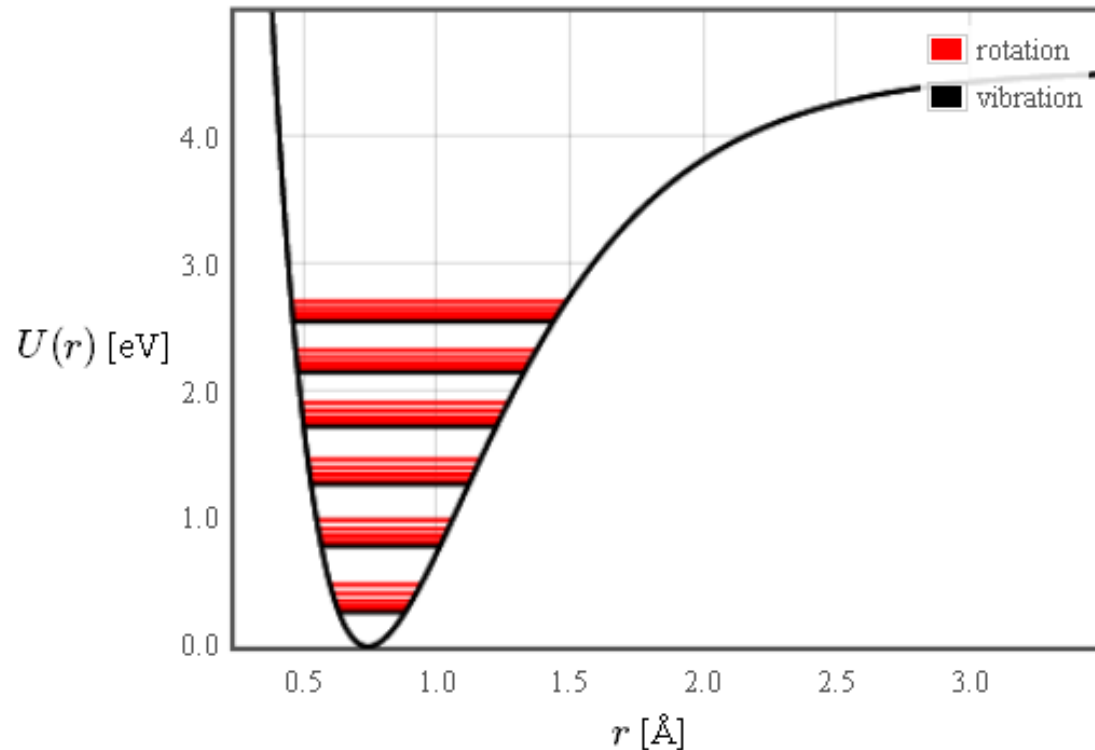
Thermal conductivity limited by Umklapp scattering

High temperature limit of specific heat does not approach the

Dulong-Petit law

# Nonlinear effects

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Thermal expansion

Thermal conductivity limited by Umklapp scattering

High temperature limit of specific heat does not approach the Dulong-Petit law

# Structural phase transitions

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Some materials make a transition from one crystal structure to another.

Two allotropes of tin: gray tin ( $\alpha$ -Sn) is stable at temperatures below 13.2°C and white tin ( $\beta$ -Sn) is stable above.

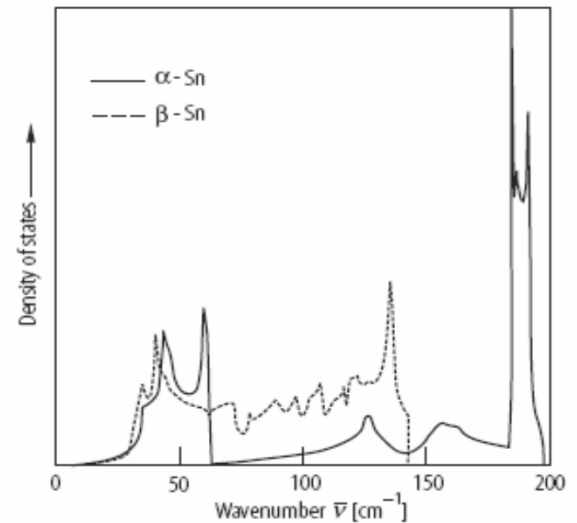
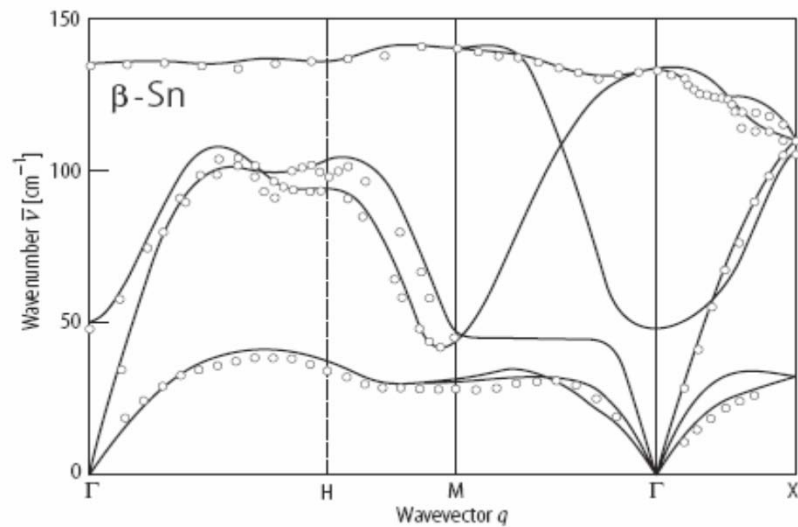
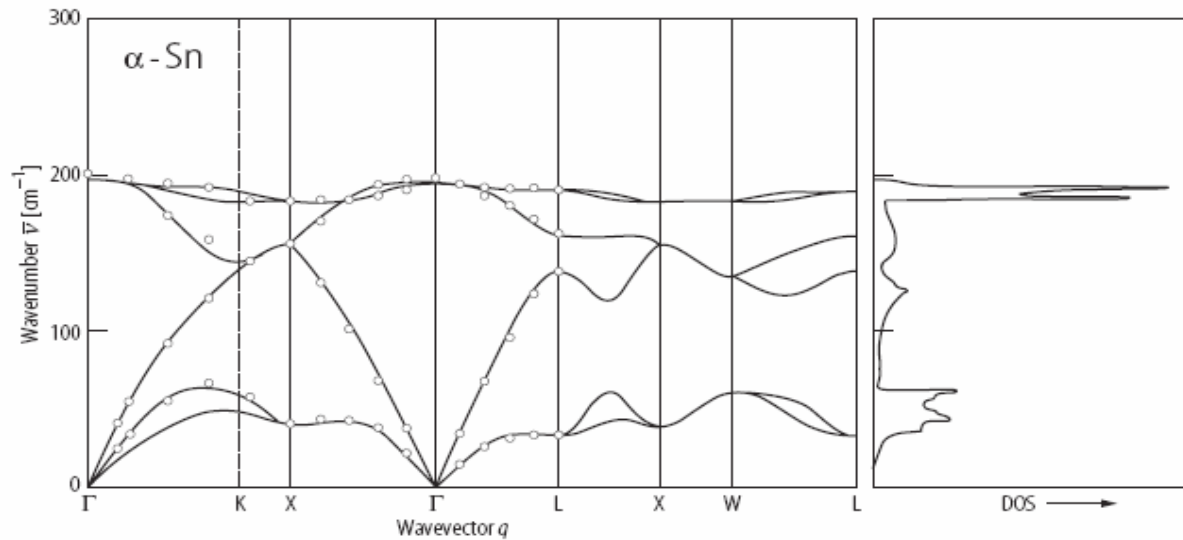
The phase with the lowest free energy prevails. (White tin can be stabilized below 13.2 C by adding impurities.)

$$F = U - TS$$

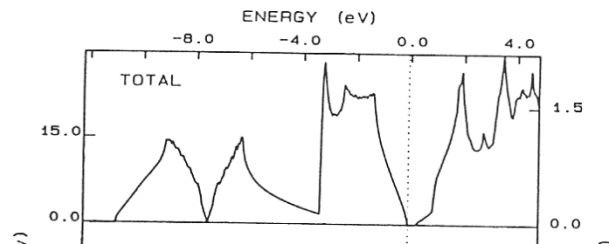
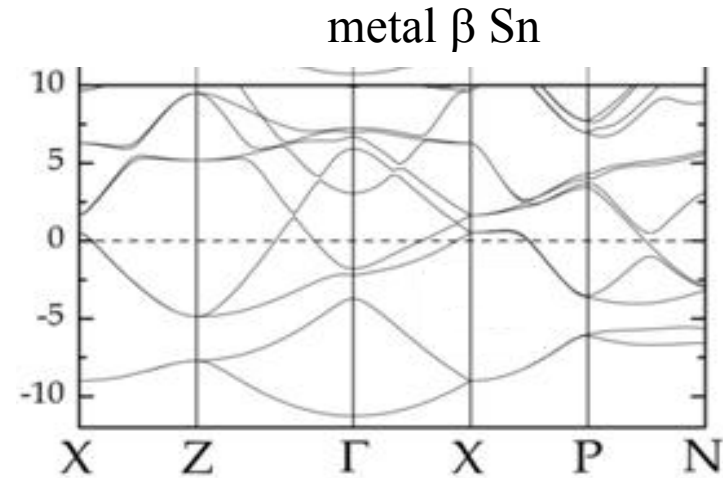
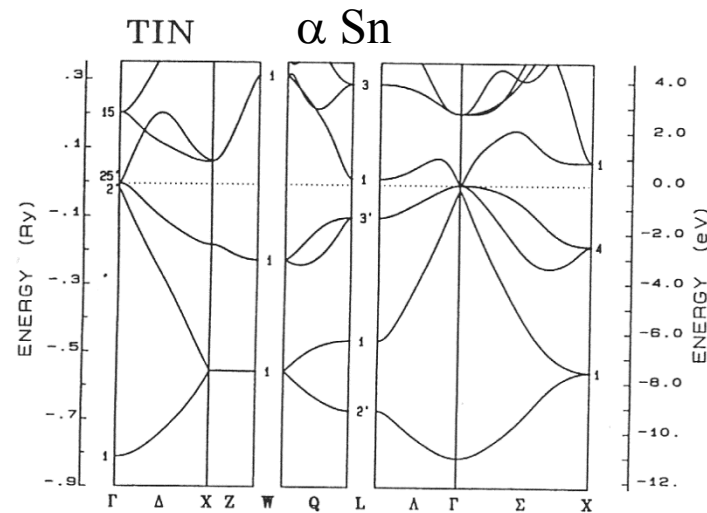
# Structural phase transition in Sn

semiconductor  
diamond crystal  
structure

metal  
tetragonal  
crystal structure



# Structural phase transition in Sn



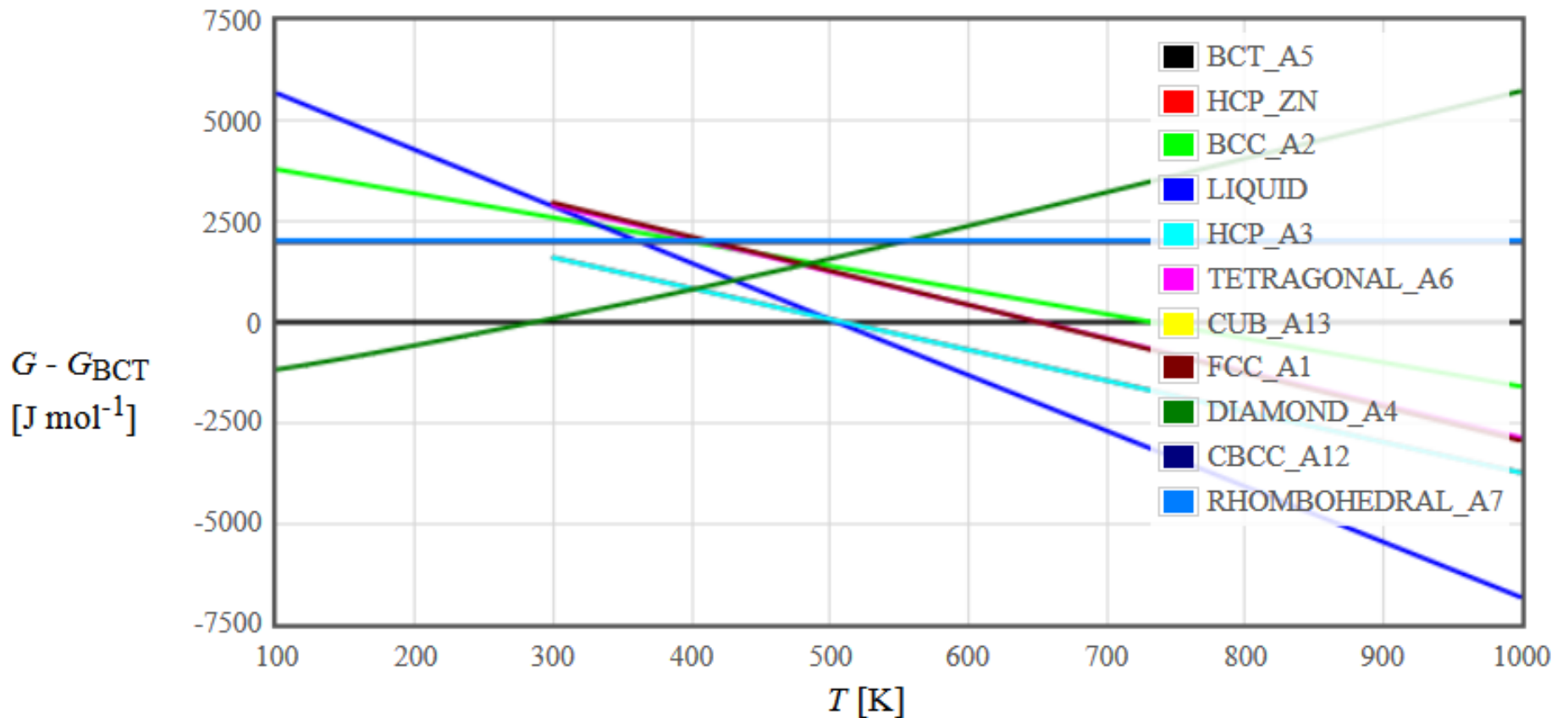
semiconductor: electrons make a negligible contribution to the entropy

$$s = \frac{\sqrt{2\pi}}{2\pi^2 \hbar^3} (m_e^* m_h^*)^{3/4} \exp\left(\frac{-E_g}{2k_B T}\right) (k_B T)^{3/2} \left(5k_B + \frac{E_g}{T}\right),$$

$$s \approx \frac{\pi^2 D(E_F)}{3} k_B^2 T$$

# Structural phase transition in Sn

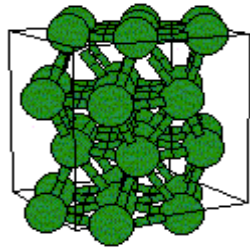
metal  $\beta$  Sn = A5



<http://lampx.tugraz.at/~hadley/ss1/materials/sgte/SGTE.html>

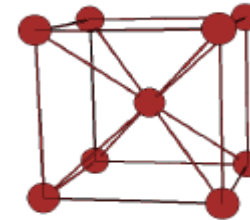
# Structural phase transitions

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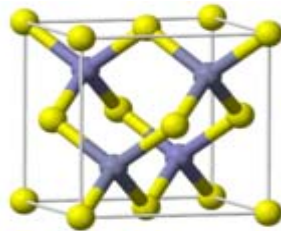


$\alpha$ -Ti, hcp

transition at 890 C

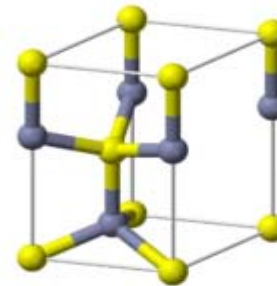


$\beta$ -Ti, bcc



Zincblende, ZnS

transition at 1020 C



Wurtzite, ZnS

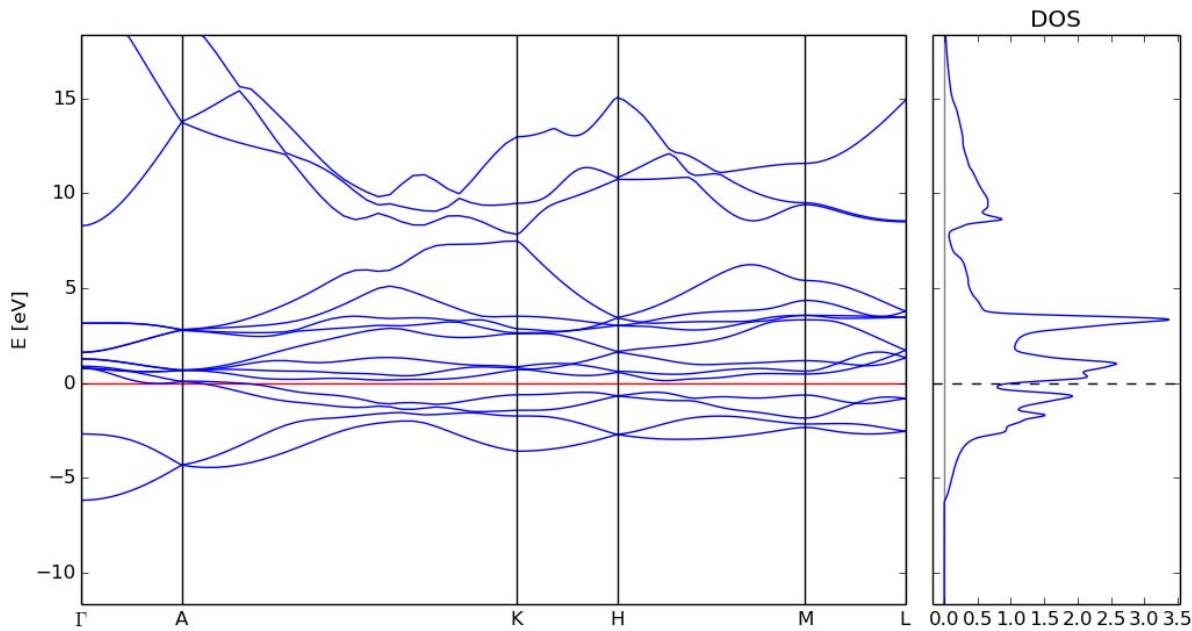
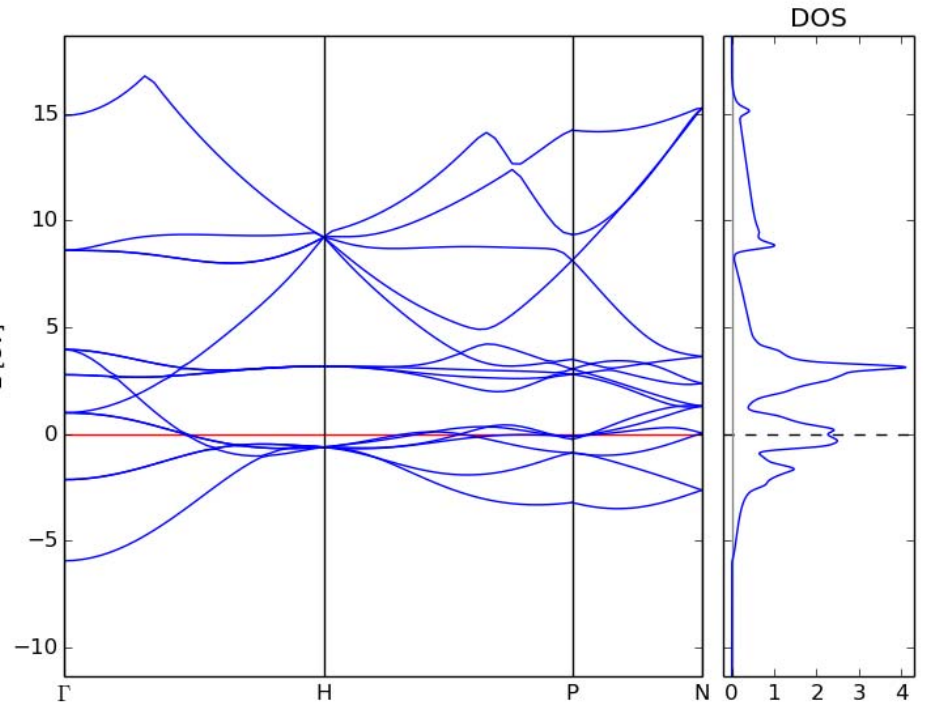
The crystal structure with the lowest free energy will be observed.  
Softer phonons >> lower Debye frequency >> more modes occupied >> higher entropy



# Titanium



bcc →

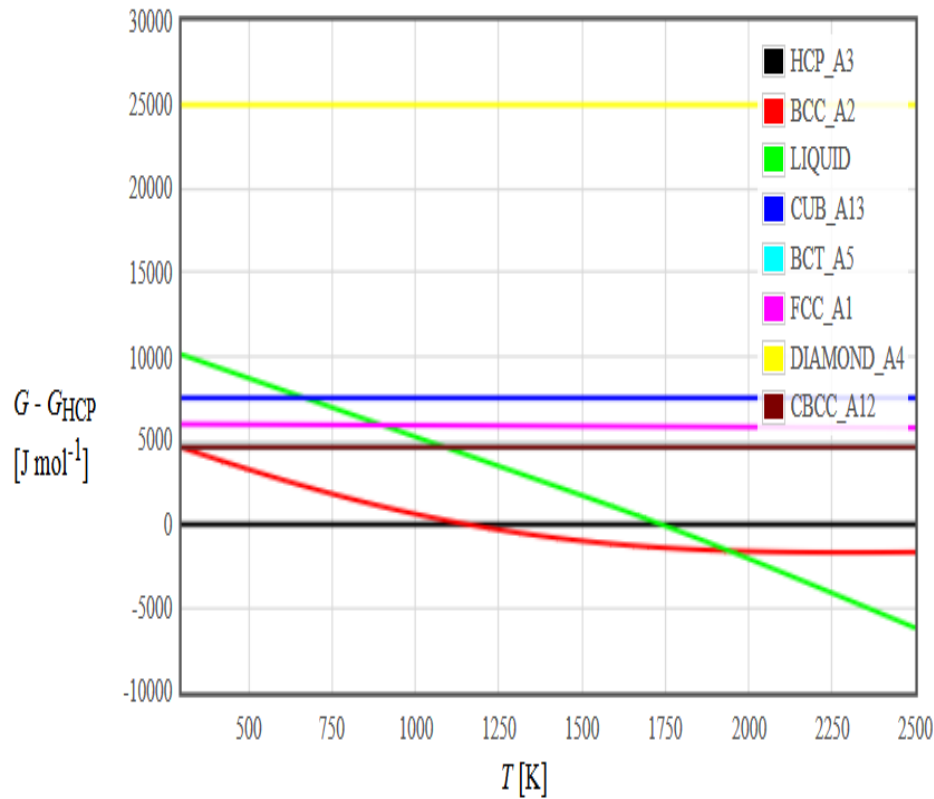


← hcp

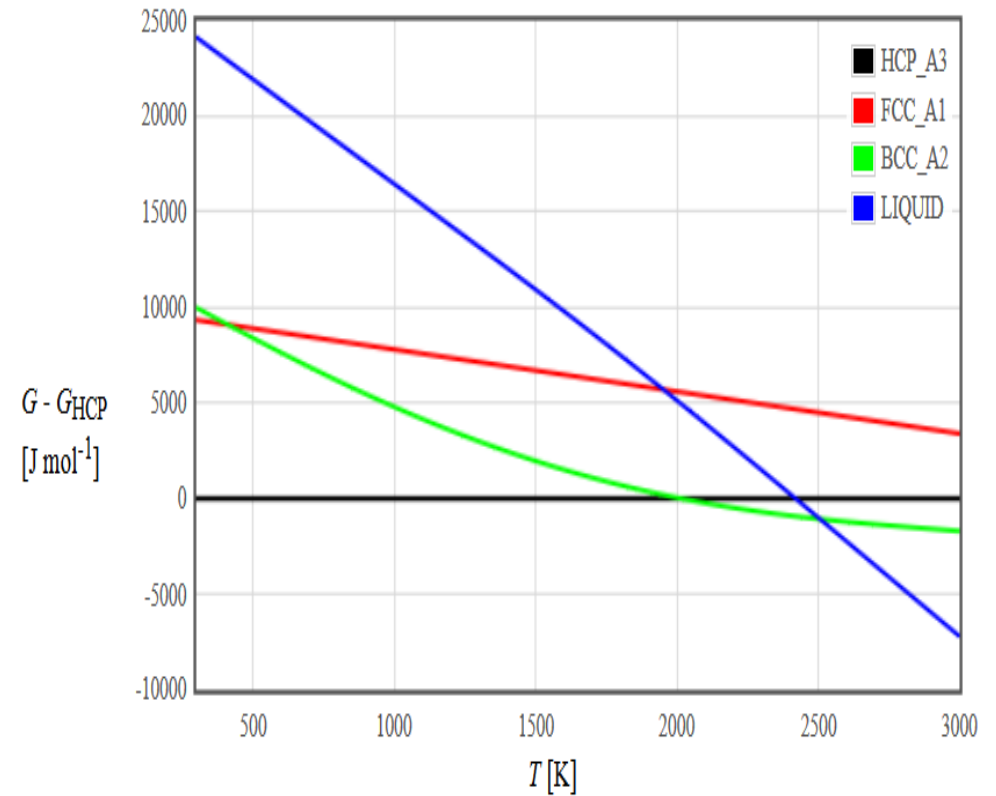
Calculated using FHI-aims by Lydia Nemec

# Close packed $\rightarrow$ bcc

Ti



Hf



Close packed  $\rightarrow$  bcc: Am, Be, Ca, Gd, Nd, Pr, Hf, Sc, Sm, Sr, Ti, Tb, Th, Tl, Y, Yb, Zr

<http://lampx.tugraz.at/~hadley/ss1/materials/sgte/SGTE.html>

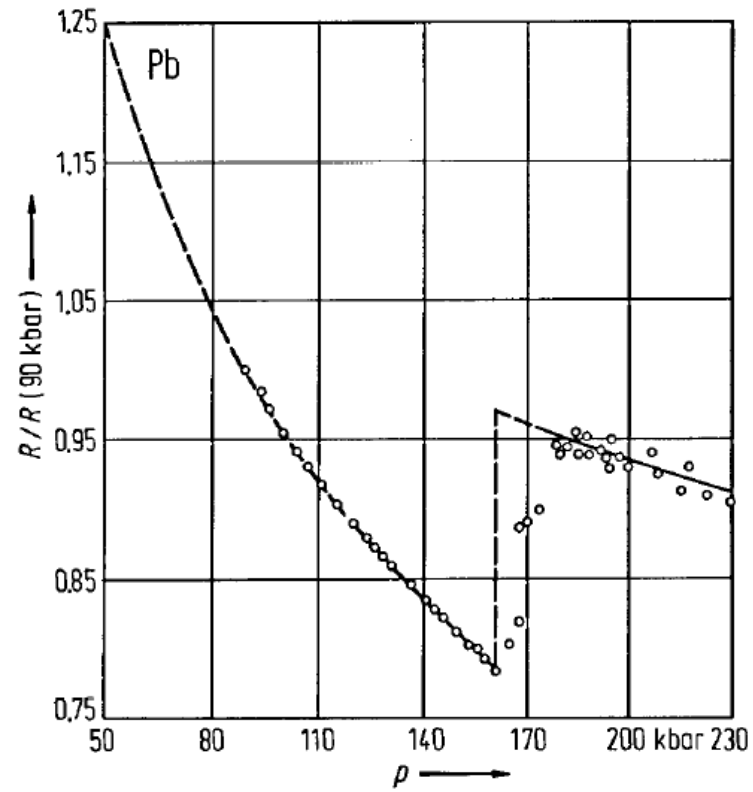
# Strain

Strain displaces the atoms and the band structure needs to be recalculated.

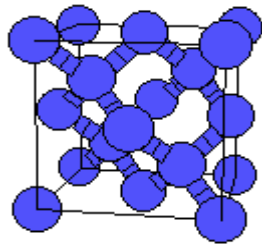
This changes the density of states and the thermodynamic properties.

Make Legendre transformations from the internal energy to the enthalpy that has temperature and pressure as independent variables. The crystal structure with lowest enthalpy will be observed.

Enthalpy is calculated from the microscopic states of electrons and phonons.

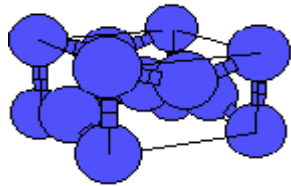


# Structural phase transitions



$\alpha$ -Sn, gray tin, diamond

transition at 13 C



$\beta$ -Sn, white tin, tetragonal

silicon makes a diamond to  $\beta$ -Sn transition under pressure

